

Resonant atomic gases



Leo Radzihovsky

for details see: *Gurarie, L.R., Annals of Physics, 322, 2-119 (2007)*

Sheehy, L.R., Annals of Physics, 322, 1790 (2007)

Giorgini, et al., RMP, 80, 885 (2008)

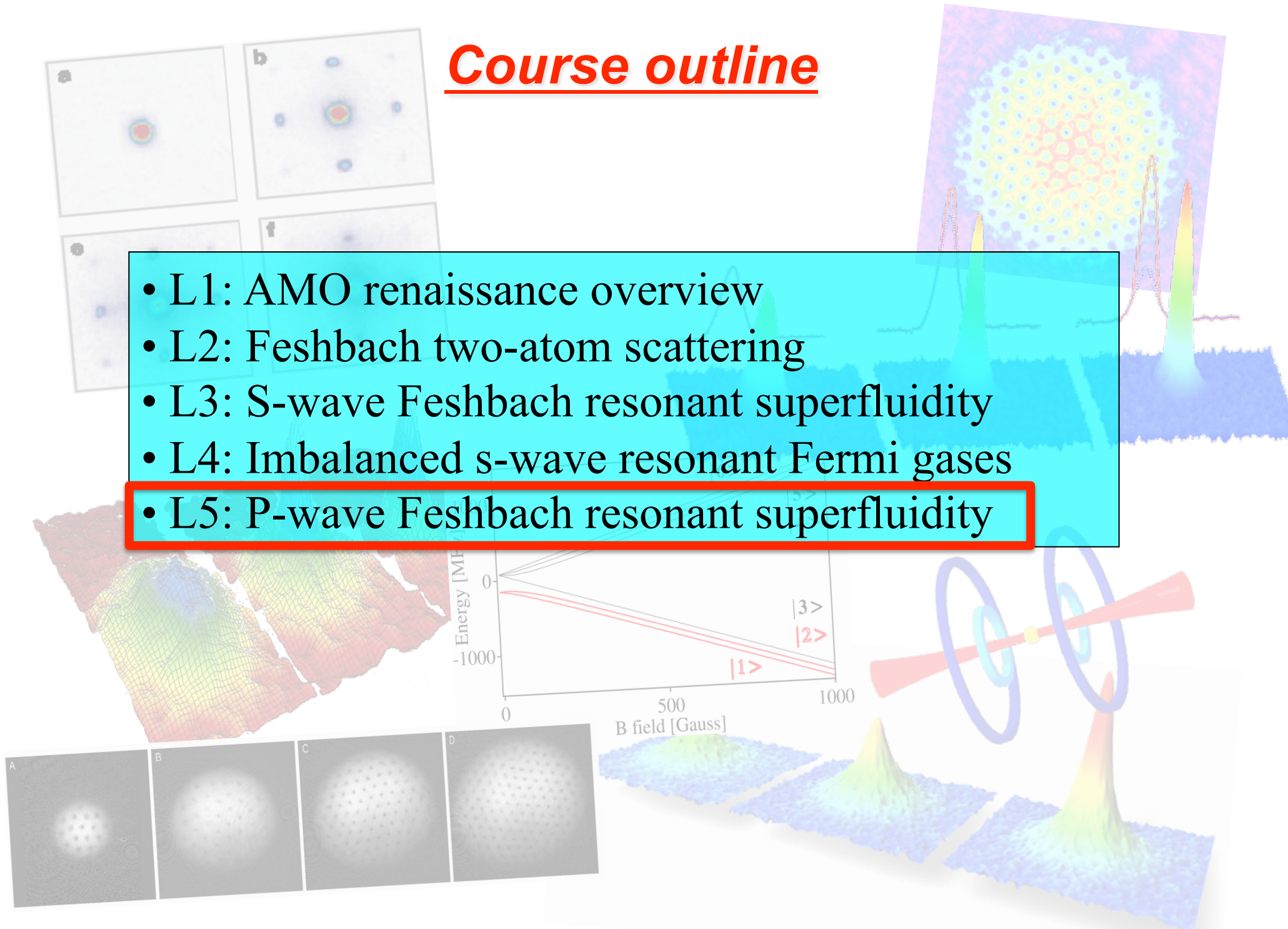
Ketterle and Zwierlein, Varenna lectures (2006)

\$: NSF

Mysore, India, Dec 2010

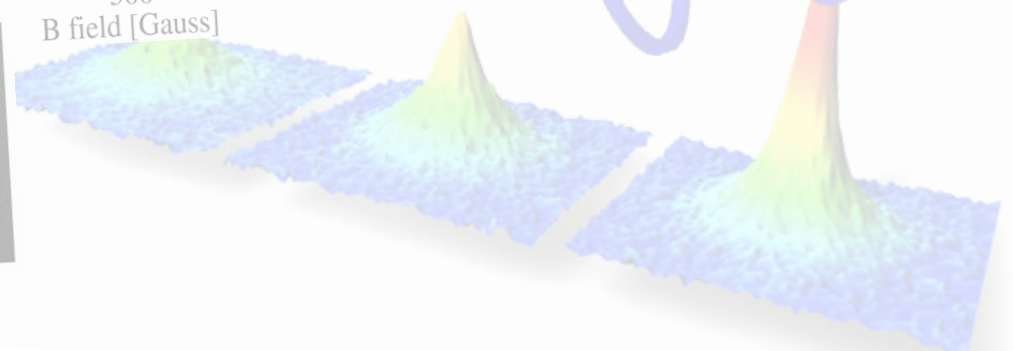
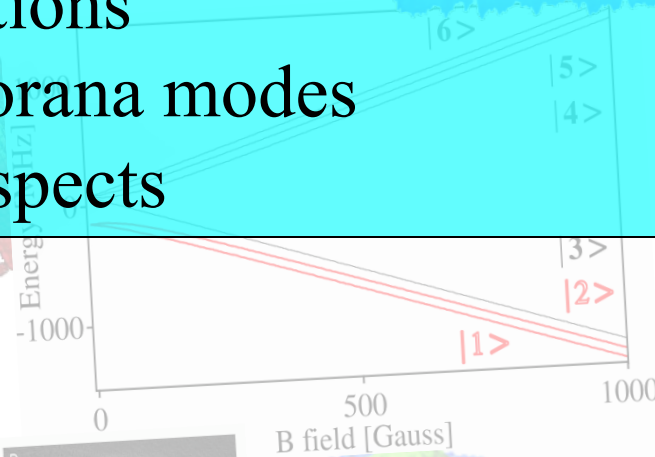
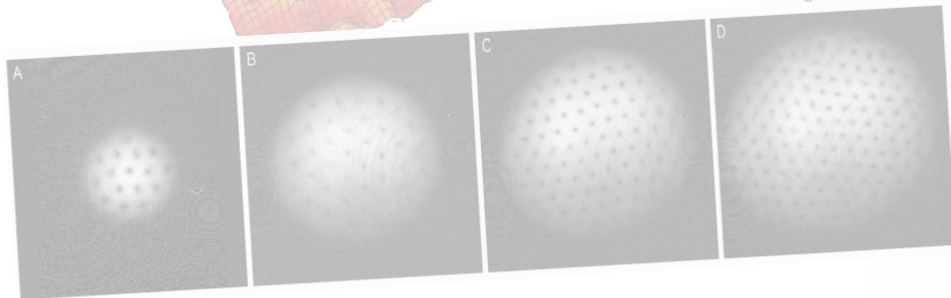
Course outline

- L1: AMO renaissance overview
- L2: Feshbach two-atom scattering
- L3: S-wave Feshbach resonant superfluidity
- L4: Imbalanced s-wave resonant Fermi gases
- L5: P-wave Feshbach resonant superfluidity



Lecture 5: *P-wave Feshbach resonance* superfluidity

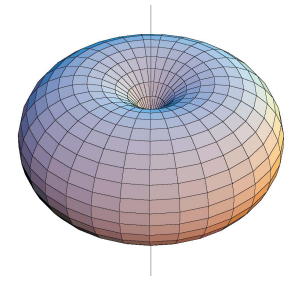
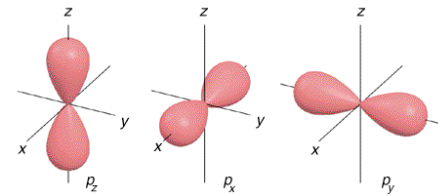
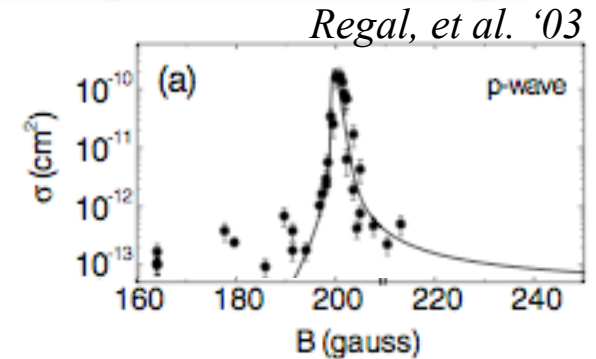
- motivation and experiments
- review of p-wave scattering theory
- two-channel model
- phases and transitions
- vortices and Majorana modes
- experimental prospects



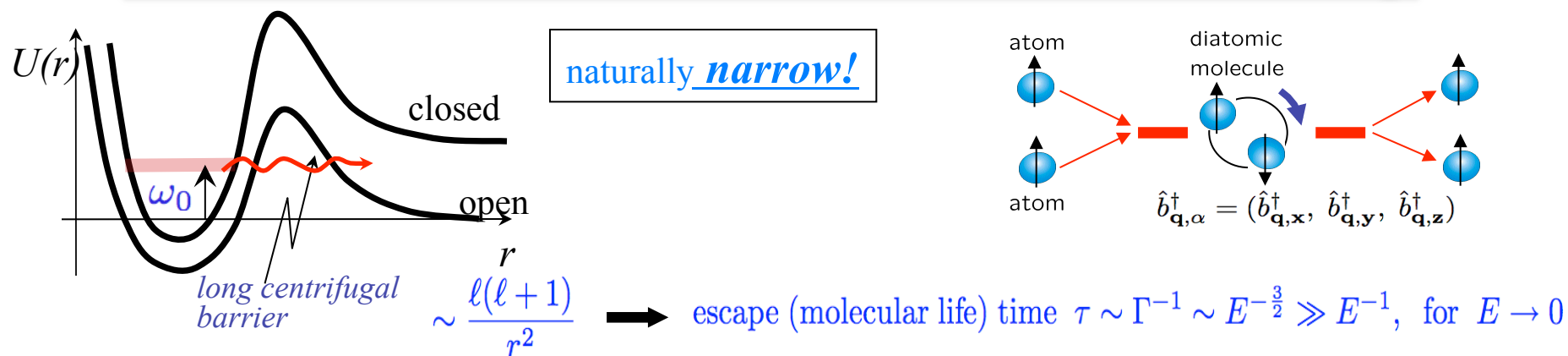
Finite angular momentum superfluidity

Motivation:

- *p-wave Feshbach resonances exist*
- *examples of ^3He and high- T_c superconductors*
- *multiple superfluids phases*
- *anisotropic gap with gapless excitations*
- *conventional (thermal and quantum) and topological phase transitions with detuning*
- *non-Abelian vortex excitations $\Rightarrow$ topological QC?*



P-wave Feshbach resonant scattering

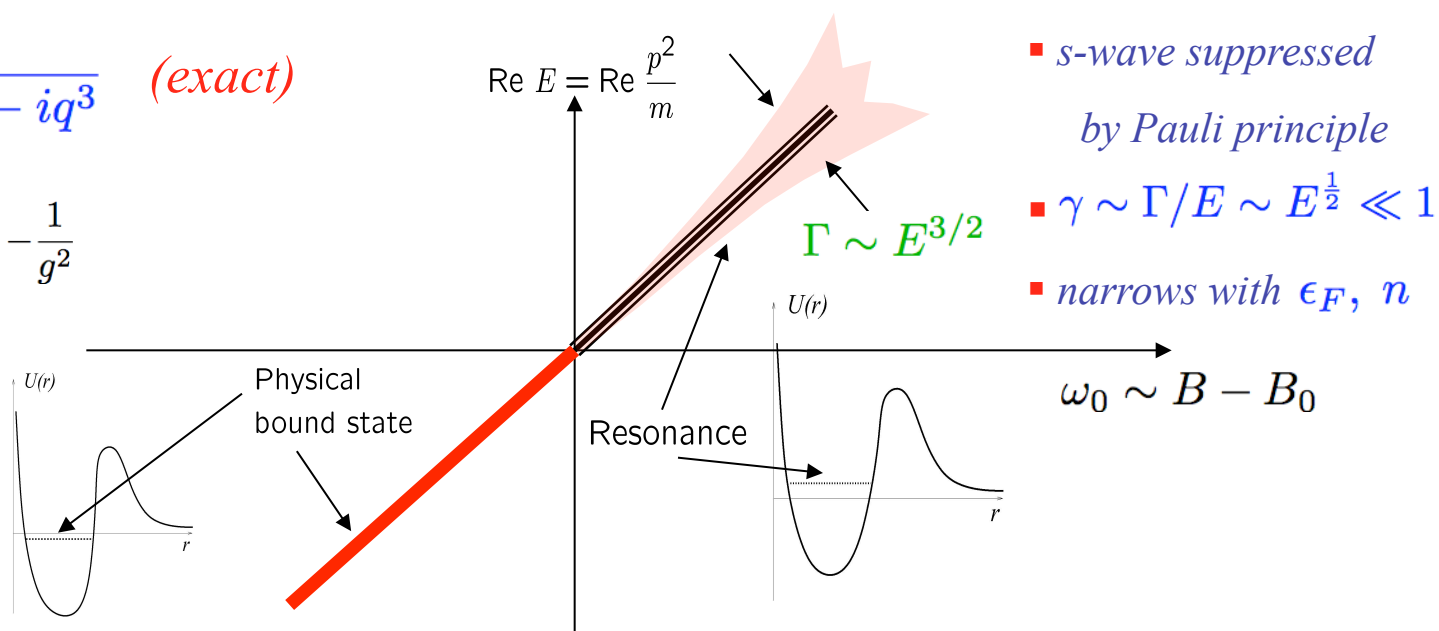


$$\mathcal{H}_{2ch} = \psi^\dagger \frac{\hat{p}^2}{2m} \psi + \vec{\phi}^\dagger \left(\frac{\hat{p}^2}{4m} + \epsilon_0 \right) \vec{\phi} - ig \vec{\phi} \cdot \psi^\dagger \nabla \psi$$

$$f_p = \frac{q^2}{-v^{-1} + \frac{q_0}{2} q^2 - iq^3} \quad (\text{exact})$$

$$\text{with } v^{-1} \sim -\frac{g^2}{\omega_0}, \quad q_0 \sim -\frac{1}{g^2}$$

$$f_p(q) = \frac{q^2}{F(q^2) - iq^3}$$

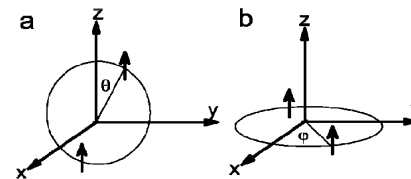
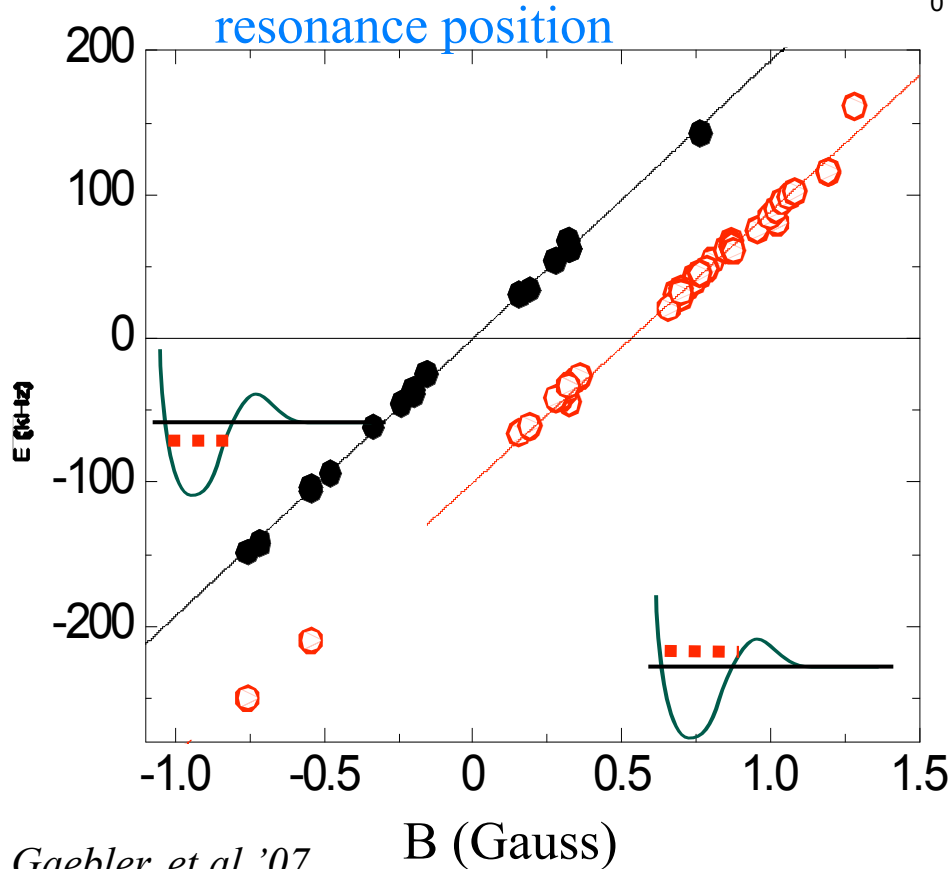
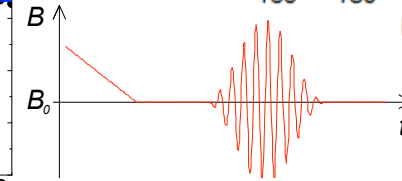
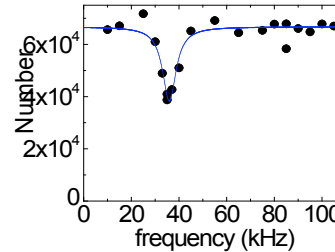
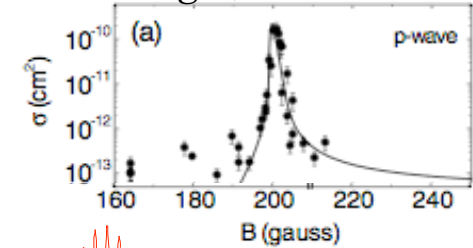


Experimental hopes for p-wave superfluidity

- p-wave Feshbach resonance in ^{40}K , ^6Li
- making p-wave molecules:

*resonant disappearance of atoms
with oscillating $B(t)$*

Regal, et al. '03



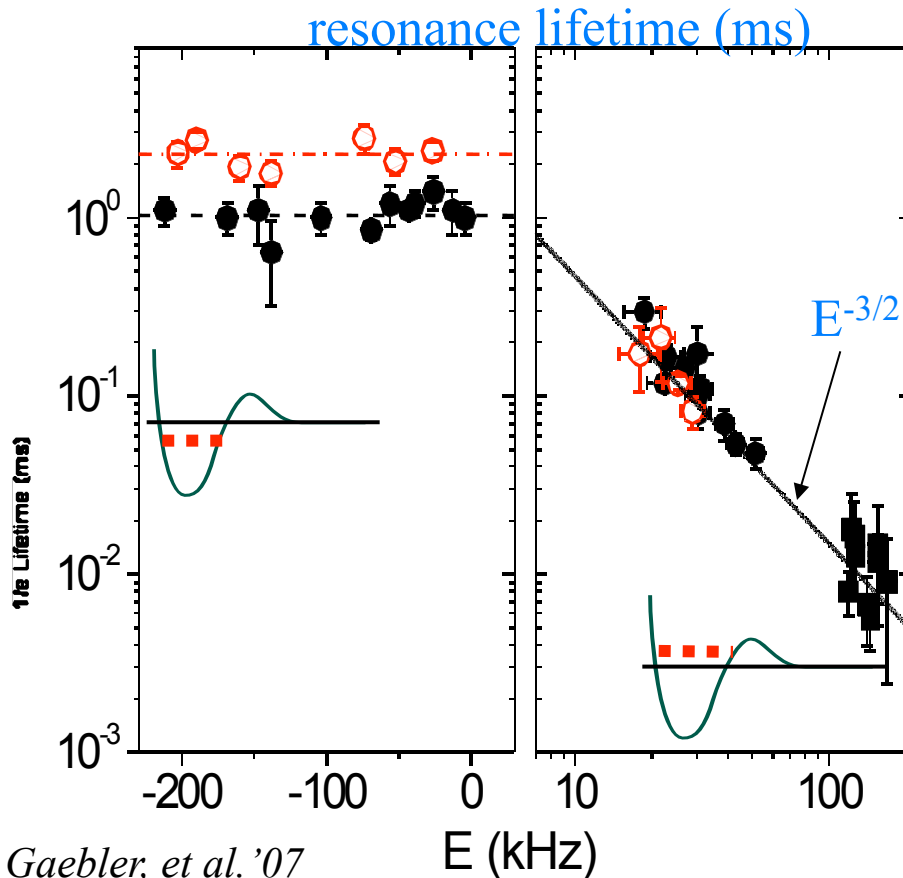
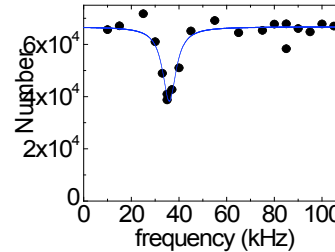
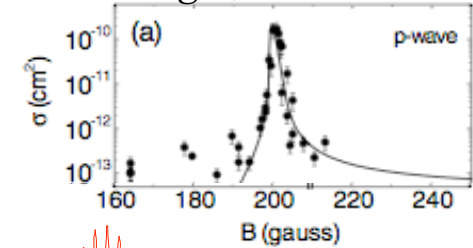
Gaebler, et al. '07

Experimental hopes for p-wave superfluidity

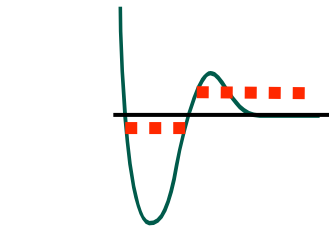
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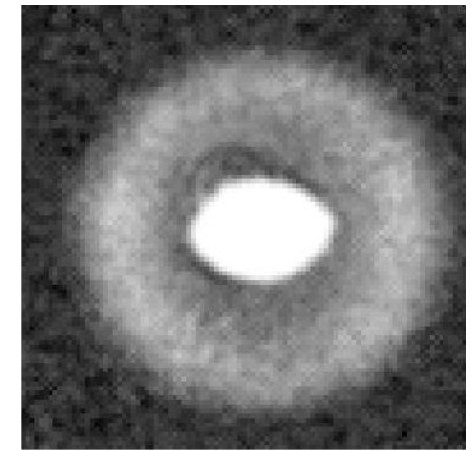
Regal, et al. '03



Gaebler, et al. '07



to see molecules:



look for energetic atoms

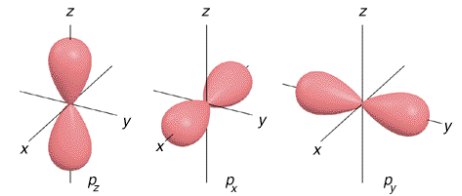
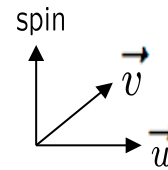
P-wave resonant superfluidity

$$\mathcal{H}_{2ch} = \psi^\dagger \frac{\hat{p}^2}{2m} \psi + \vec{\phi}^\dagger \left(\frac{\hat{p}^2}{4m} + \epsilon_0 \right) \vec{\phi} - ig \vec{\phi} \cdot \psi^\dagger \nabla \psi$$

dimensionless coupling: $\gamma \sim \left(\frac{g\sqrt{n}k_F}{\epsilon_F} \right)^2 \sim g^2 \epsilon_F^{1/2} \sim \frac{n^{1/3}}{q_0}$

- **narrow** resonance $\gamma \ll 1 \rightarrow$ MFT : $\vec{\phi}(x) = \vec{B}$

- **complex vector** order parameter:



$$\vec{B} = \vec{u} + i \vec{v} \iff \psi_0 = B_z, \psi_{\pm} = \pm(B_x \pm iB_y)$$

- sample states:

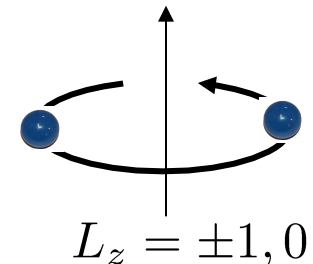
$$v = 0 \iff |m = 0\rangle \text{ along } \vec{u}$$

(k_x β - state in ^3He)

$$u = v \iff |m = 1\rangle \text{ along } \vec{u} \times \vec{v}$$

($k_x + ik_y$ "axial" Anderson - Morel state in ^3He)

$$\vec{B} \cdot \vec{k} = \sum_{m=0,\pm k} \psi_m Y_{1,m}(\hat{k}) k$$



Mean-field theory ($\gamma \sim g^2 \epsilon_F^{1/2} \ll 1$)

$$H_p = \sum_{\mathbf{k}} \frac{k^2}{2m} a_{\mathbf{k}}^\dagger a_{\mathbf{k}} + \sum_{\mathbf{q}, \alpha} \left(\frac{q^2}{4m} + \epsilon_{0\alpha} \right) b_{\mathbf{q}, \alpha}^\dagger b_{\mathbf{q}, \alpha} + \sum_{\mathbf{k}} [\Delta_{\mathbf{k}}(\vec{B}) a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger + h.c.]$$

- *superfluid ground state:*

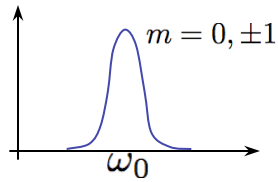
$$\begin{array}{cc} \text{molecular BEC } |\vec{B}\rangle & + \text{ Cooper pairing } |\text{BCS}_{\vec{B}}\rangle = \prod_{\mathbf{k}} (u_{\mathbf{k}} + v_{\mathbf{k}} a_{-\mathbf{k}}^\dagger a_{\mathbf{k}}^\dagger) |0\rangle \\ \text{(closed)} & \text{(open)} \end{array}$$

- *excitation spectrum:* $H_{ex} = \sum_{\mathbf{k}} E_{\mathbf{k}}^{(a)} \alpha_{\mathbf{k}}^\dagger \alpha_{\mathbf{k}} + \sum_{\mathbf{k}, \alpha} E_{\mathbf{k}, \alpha}^{(m)} \beta_{\mathbf{k}, \alpha}^\dagger \beta_{\mathbf{k}, \alpha}$

$$E_{\mathbf{k}}^{(a)} = \sqrt{\epsilon_{\mathbf{k}}^2 + \Delta_{\mathbf{k}}^2}, \quad E_{\mathbf{k}, \alpha}^{(m)} = \sqrt{\epsilon_{\mathbf{k}}^2 + \mu_{\alpha} \epsilon_{\mathbf{k}}} \quad \text{with gap : } \Delta_{\mathbf{k}} = 2g|\vec{B} \cdot \vec{k}|$$

- $\vec{B}$, n_b , n_a , μ determined by :

- energy minimization (gap equation) $\rightarrow \frac{\partial E(\vec{B})}{\partial B_{\alpha}} = 0$
- atom number equation $\rightarrow 2n_b + n_a = n$



Isotropic resonance at $T=0$ ($\omega_\alpha = \omega_0$)

$$E = (u^2 + v^2) [\omega_0 - 2\mu + \underbrace{a_1 \ln \{a_0 (u + v)\}}_{\text{BCS}} + \underbrace{a_1 \frac{u^3 + v^3}{u + v} + a_2 \left[(u^2 + v^2)^2 + \frac{1}{2} (u^2 - v^2)^2 \right]}_{\text{BEC: } (\vec{B}^* \cdot \vec{B})^2 + \frac{1}{2} |\vec{B} \cdot \vec{B}|^2}]$$

- BCS ($\omega_0 \gg 2\epsilon_F$) :

➤ $\mu \approx \epsilon_F + \mathcal{O}(\gamma)$

➤ $\frac{E_{k_x + ik_y}}{E_{k_x}} = \frac{1}{2}e > 1$

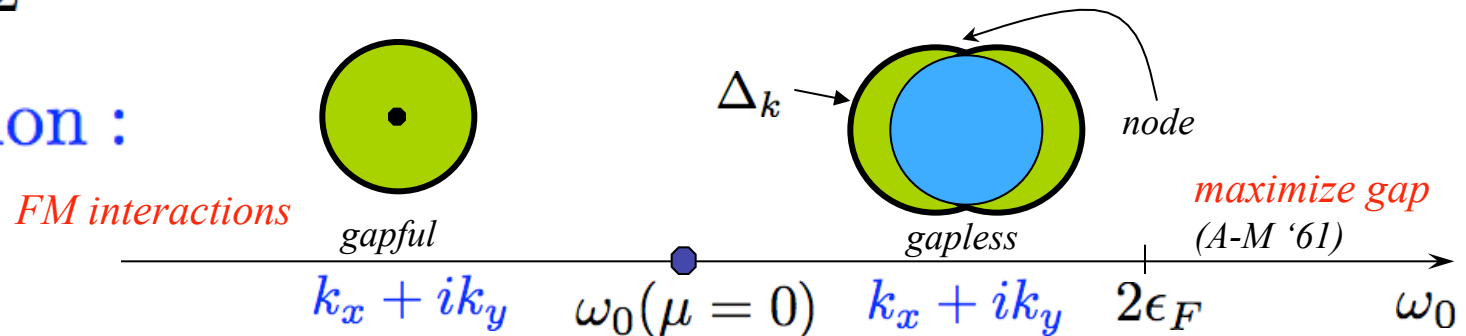
⇒ (Anderson – Morel A_1 phase)

$u = v \sim e^{-(\omega_0 - 2\epsilon_F)/\gamma\epsilon_F} \Rightarrow k_x + ik_y (m = 1)$

- BEC ($\omega \ll 2\epsilon_F$) :

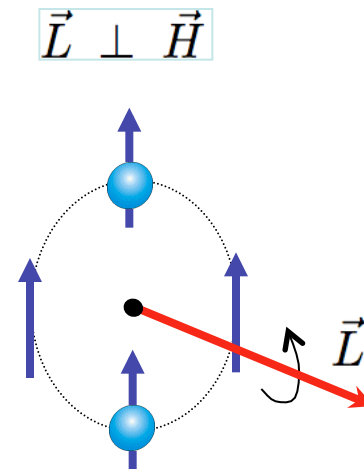
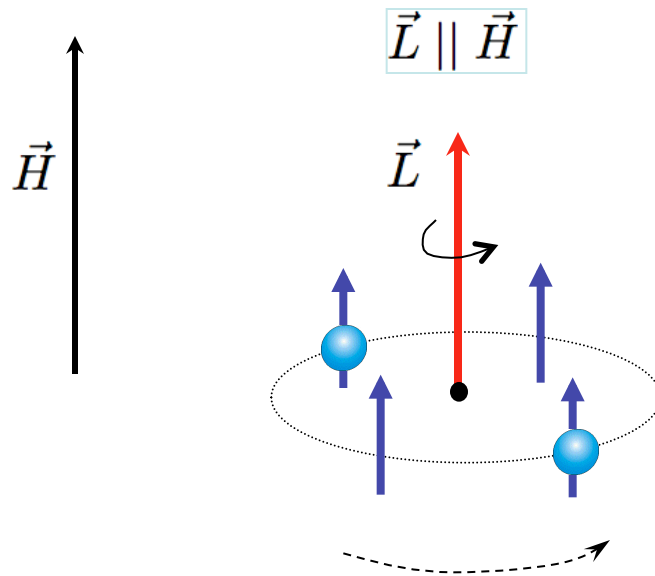
➤ $\mu \approx \frac{1}{2}\omega_0 + \mathcal{O}(\gamma) \Rightarrow u = v \approx \sqrt{n - n(\omega_0/2)} \Rightarrow k_x + ik_y (m = 1)$

- transition :

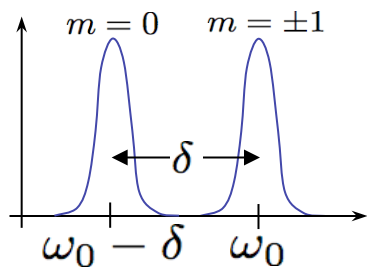


Dipolar-interaction FR splitting

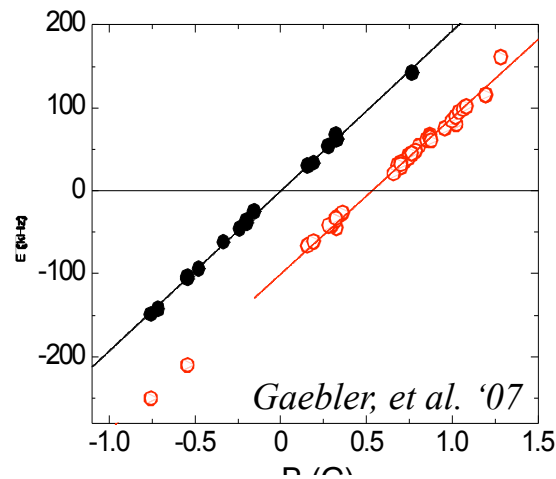
Ticknor, et al '03



$$H_{\text{molecule}} = \omega_{\parallel} b_{\parallel}^{\dagger} b_{\parallel} + \omega_{\perp} \vec{b}_{\perp}^{\dagger} \cdot \vec{b}_{\perp} \quad \omega_{\parallel} < \omega_{\perp}$$



Ticknor, Regal, et al. '03

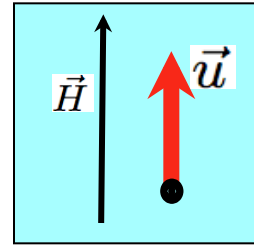


$$E = \sum_{\alpha} (u_{\alpha}^2 + v_{\alpha}^2) [\omega_{\alpha} - 2\mu + \dots]$$

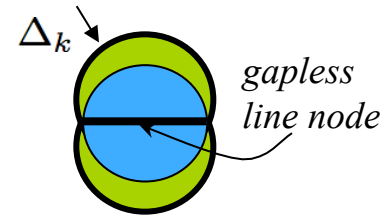
P-wave superfluid phases

$$E = \sum_{\alpha} (u_{\alpha}^2 + v_{\alpha}^2) [\omega_{0\alpha} - 2\mu + a_1 \ln \{a_0 (u + v)\}] + a_1 \frac{u^3 + v^3}{u + v} + a_2 \left[(u^2 + v^2)^2 + \frac{1}{2} (u^2 - v^2)^2 \right]$$

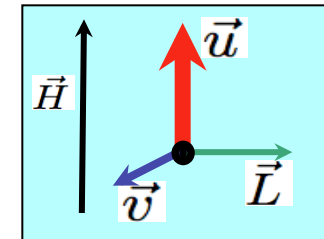
- k_x - state: β -phase of ^3He ($m_{||}=0$)



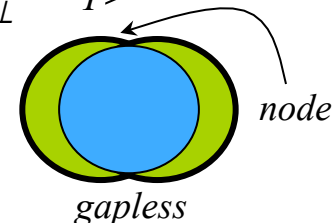
- $u \approx u_0 e^{\delta}$, $v = 0$
- equatorial node line for $\mu > 0 \Rightarrow C \sim T^{\alpha}$, ...
- fully gapped for $\mu < 0 \Rightarrow C \sim e^{-|B-B_0|^2/T}$, ...
- spontaneously broken symmetries: $U(1)$



- $k_x + i \sigma k_y$ - state: “deformed” A_1 -phase of ^3He ($m_{\perp}=1$)

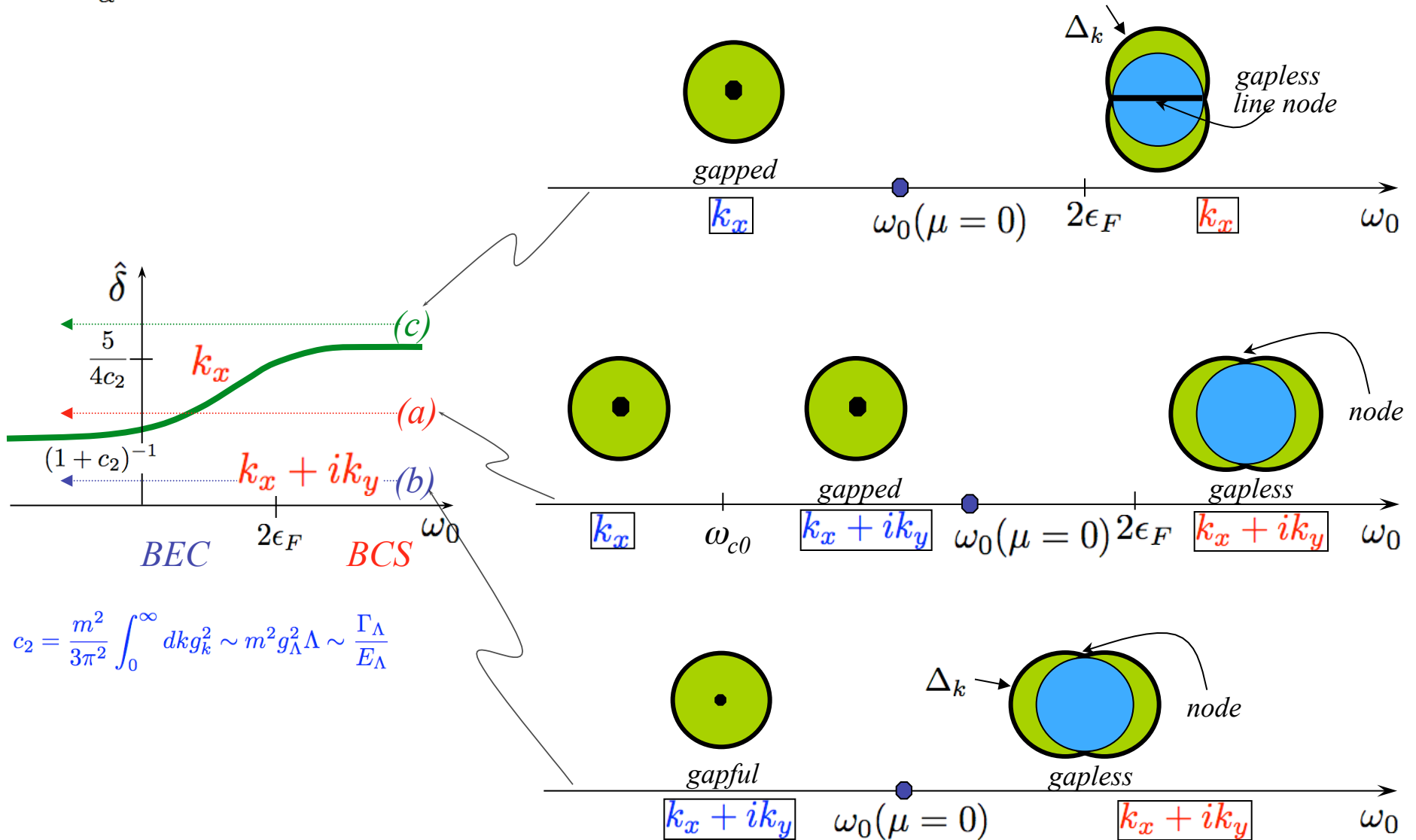


- $u \approx u_0 (1 + \delta) e^{\delta/2}$, $v \approx u_0 (1 - \delta) e^{\delta/2} \Rightarrow |m_{\perp}=1\rangle + \delta |m_{\perp}=-1\rangle$
- polar point nodes for $\mu > 0 \Rightarrow C \sim T^{\alpha}$, ...
- fully gapped for $\mu < 0 \Rightarrow C \sim e^{-|B-B_0|^2/T}$, ...
- spontaneously broken symmetries: $U(1)$, $O(2)$, T



$T=0$ phase diagrams also see Cheng, Yip '05

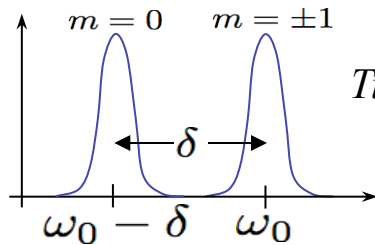
$$E = \sum_{\alpha} (u_{\alpha}^2 + v_{\alpha}^2) [\omega_{0\alpha} - 2\mu + a_1 \ln \{a_0 (u + v)\}] + a_1 \frac{u^3 + v^3}{u + v} + a_2 \left[(u^2 + v^2)^2 + \frac{1}{2} (u^2 - v^2)^2 \right]$$



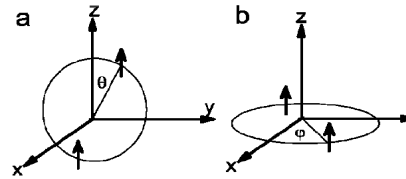
Anisotropic p -wave superfluidity

Gurarie, L.R., Andreev '05
Cheng and Yip '05

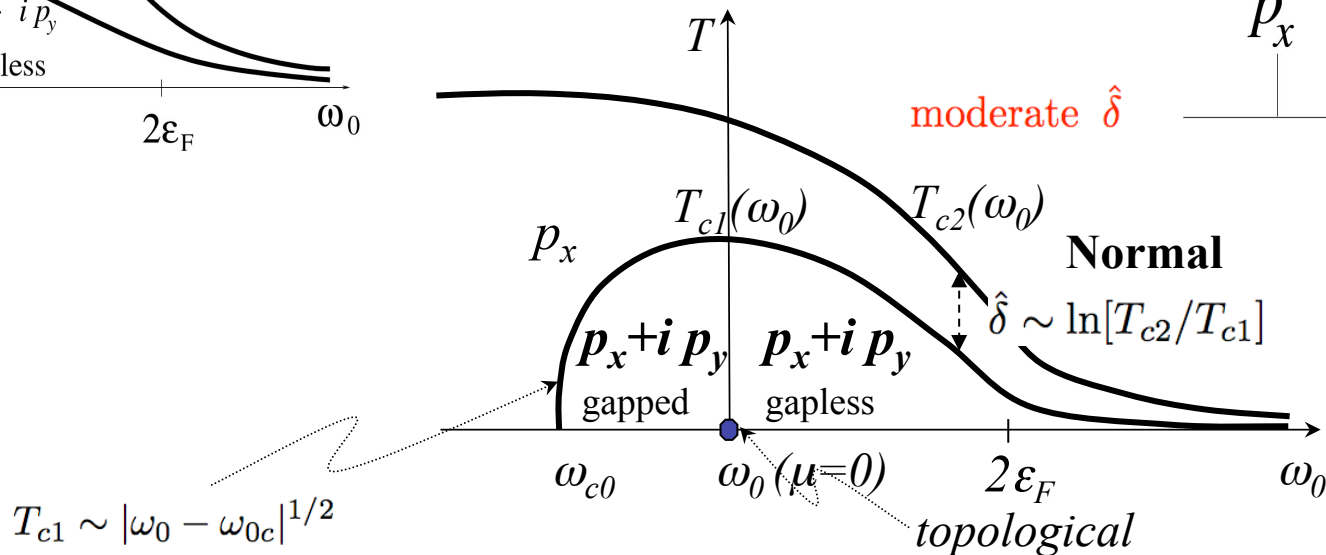
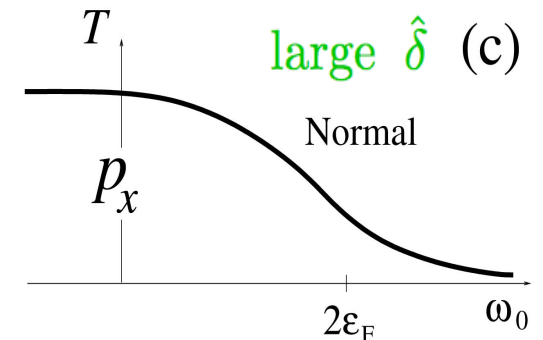
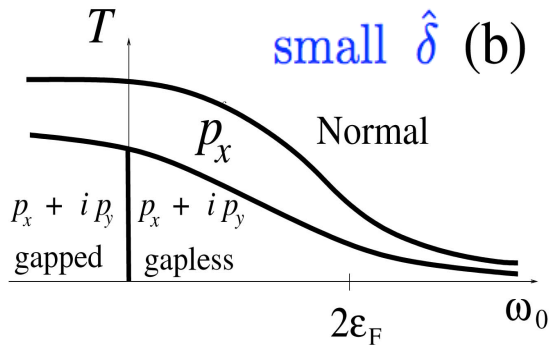
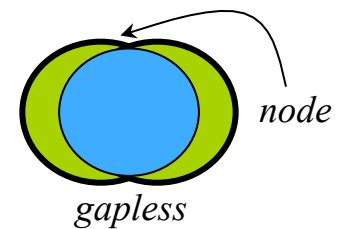
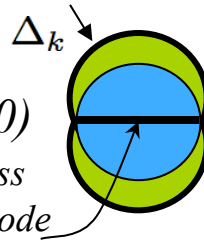
- resonance splits:



Ticknor, Regal, et al.

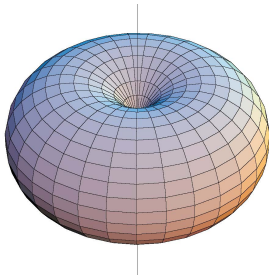


- two competing states: p_x ($m_z=0$) and $p_x + i p_y$ ($m_z=\pm 1$)

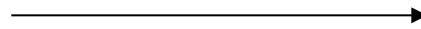


Gapless $\longleftrightarrow$ gapped superfluid transitions

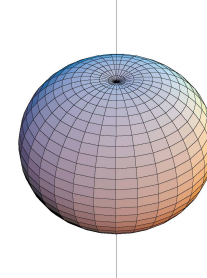
$p_x + i p_y$



μ changes sign



$p_x + i p_y$



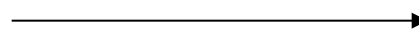
$$E_p = \sqrt{\left(\frac{p^2}{2m} - \mu\right)^2 + g^2 (p_x^2 + p_y^2) |B|^2}$$

$$E_p = \sqrt{\left(\frac{p^2}{2m} + |\mu|\right)^2 + g^2 (p_x^2 + p_y^2) |B|^2}$$

G. E. Volovik, JETP Lett. 80, 343 (2004)

p_x

μ changes sign



p_x

$$E_p = \sqrt{\left(\frac{p^2}{2m} - \mu\right)^2 + g^2 p_x^2 |B|^2}$$

$$E_p = \sqrt{\left(\frac{p^2}{2m} + |\mu|\right)^2 + g^2 p_x^2 |B|^2}$$

$p_x + i p_y$ superfluid in 2D

- Pfaffian (Moore-Read) state from FQH $|p_x + i p_y_{BCS}\rangle = \prod_p [u_p + v_p a_{-p}^\dagger a_p^\dagger] |0\rangle$

$$\Psi(z_1, z_2, \dots, z_{2N}) = \sum_P (-1)^P \frac{1}{z_{P_1} - z_{P_2}} \frac{1}{z_{P_3} - z_{P_4}} \dots \frac{1}{z_{P_{N-1}} - z_{P_N}}$$

- topological classification in terms of u_p and v_p

Anderson's pseudospin

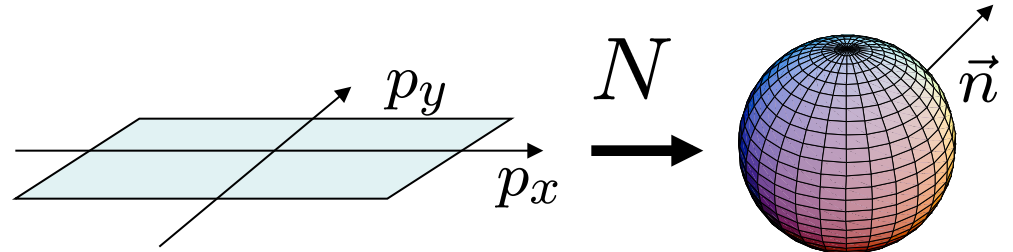
$$\begin{cases} n_x + i n_y = 2v^* u \\ n_z = |v|^2 - |u|^2 \end{cases}$$

$$\vec{n} = \frac{1}{\sqrt{\left(\frac{p^2}{2m_a} - \mu\right)^2 + 4B^2(p_x^2 + p_y^2)}} \begin{pmatrix} 2gBp_x \\ -2gBp_y \\ \frac{p^2}{2m_a} - \mu \end{pmatrix}$$

Explicit calculations show that

$$N=0 \text{ if } \mu < 0$$

$$N=1 \text{ if } \mu > 0$$



$$N = \frac{1}{8\pi} \int d^2p \left[\vec{n} \cdot \partial_\alpha \vec{n} \times \partial_\beta \vec{n} \epsilon_{\alpha\beta} \right]$$

topological invariant

$p_x + i p_y$ superfluid in 2D

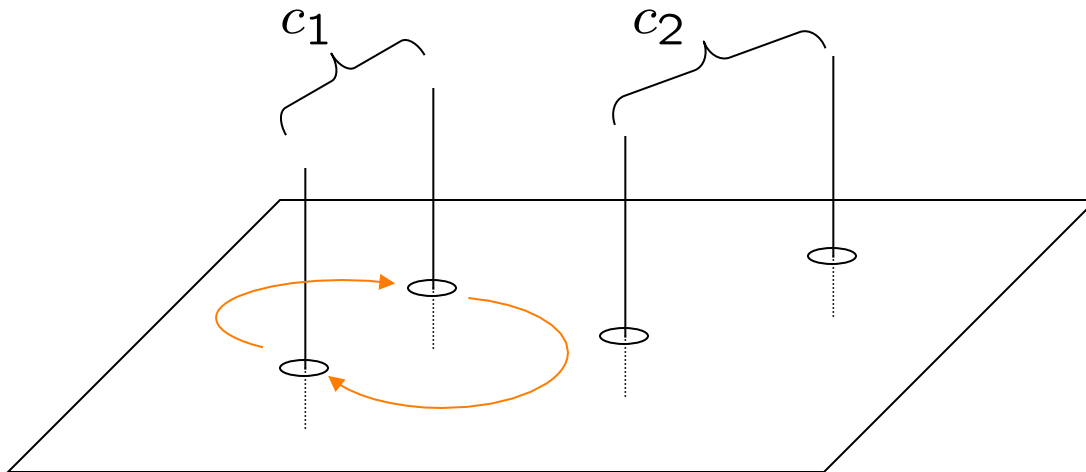
- Pfaffian (Moore-Read) state from FQH $|p_x + i p_y_{BCS}\rangle = \prod_p [u_p + v_p a_{-p}^\dagger a_p^\dagger] |0\rangle$

$$\Psi(z_1, z_2, \dots, z_{2N}) = \sum_P (-1)^P \frac{1}{z_{P_1} - z_{P_2}} \frac{1}{z_{P_3} - z_{P_4}} \dots \frac{1}{z_{P_{N-1}} - z_{P_N}}$$

- topological classification in terms of u_p and v_p
- gapped (N=1, BCS) $\Rightarrow$ gapped (N=0, BEC) superfluid transition at $\mu=0$

Read and Green, PRB 61, 10267 (2000)

- vortex excitations with non-Abelian statistics *Ivanov, PRL (2001)*



one fermion (2 states – either empty or occupied fermion) per two vortices

$2^{\frac{n}{2}}$ states per n vortices

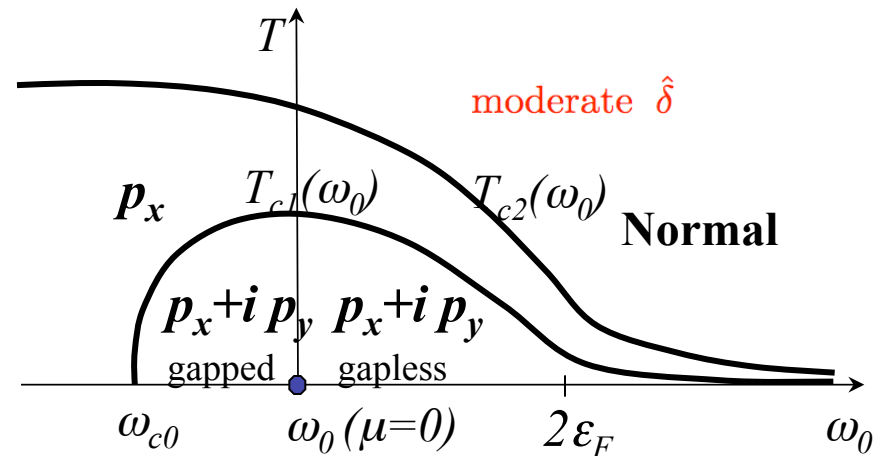
- suggested to be used as qubits for quantum computers

Kitaev, Ann. Phys. 303, 2 (2003)

Summary of p-wave superfluidity

- mapped out T , $\omega_0 \propto B$, δ phase diagram for p-wave Feshbach resonant Fermi gas

- p_x and $p_x + i p_y$ superfluids
 - thermal, quantum and topological $SF \Rightarrow SF$ transitions



- quantitatively accurate description for small $\gamma = \Gamma/\epsilon_F$ (low n)
- realization of topological states, majorana zero modes, and non-Abelian statistics of Pfaffian (Moore-Read) state
- p-wave Feshbach molecules observed in K^{40}
- ...BUT
 - ❖ short (msec) molecular lifetime (see Levinson, et al, PRL 2007)
 - ❖ what about Li^6
 - ❖ need better quantitative understanding of stability