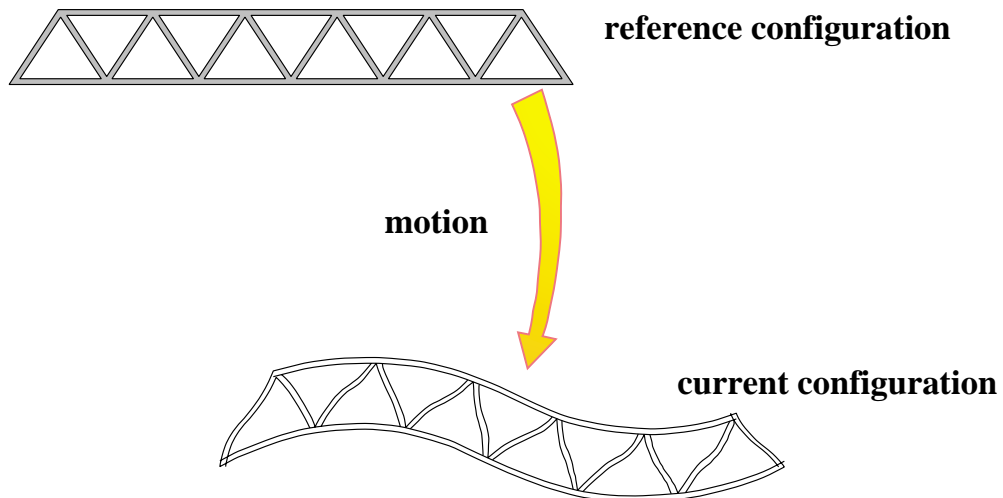


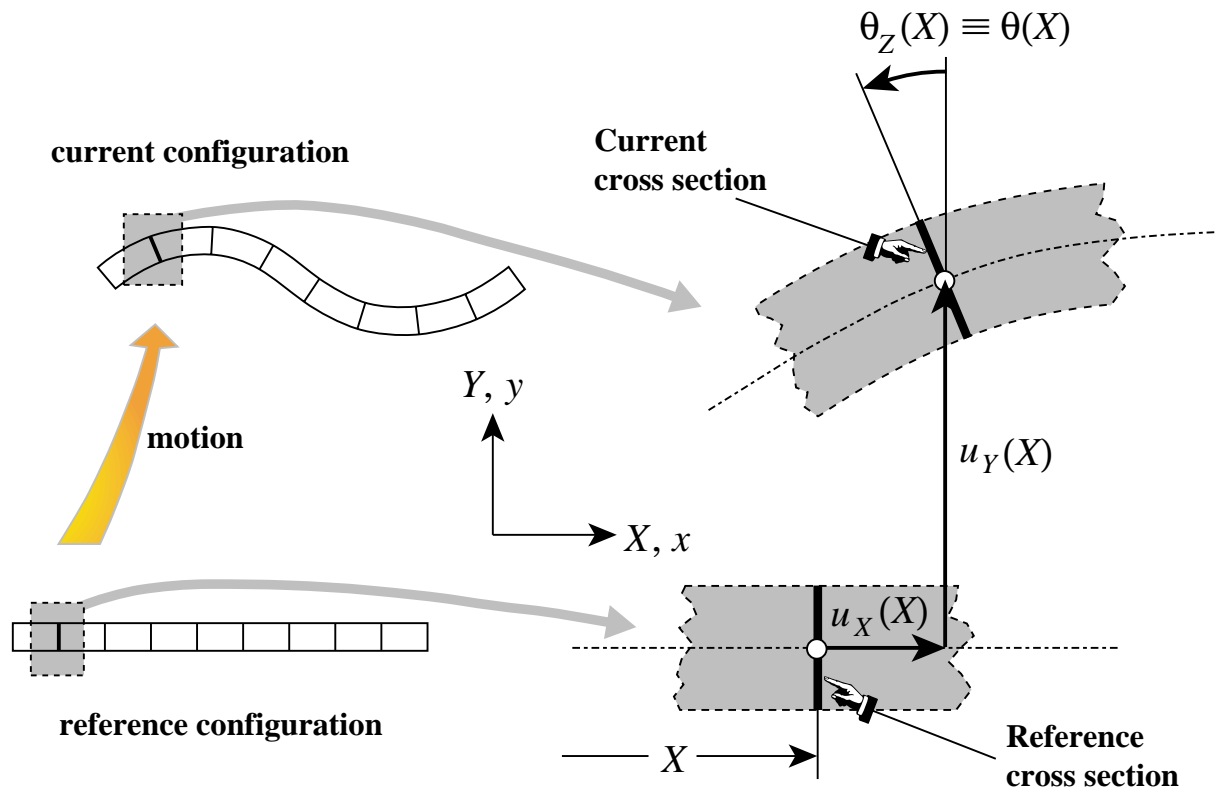
# 11

## The TL Timoshenko Beam Element: Formulation

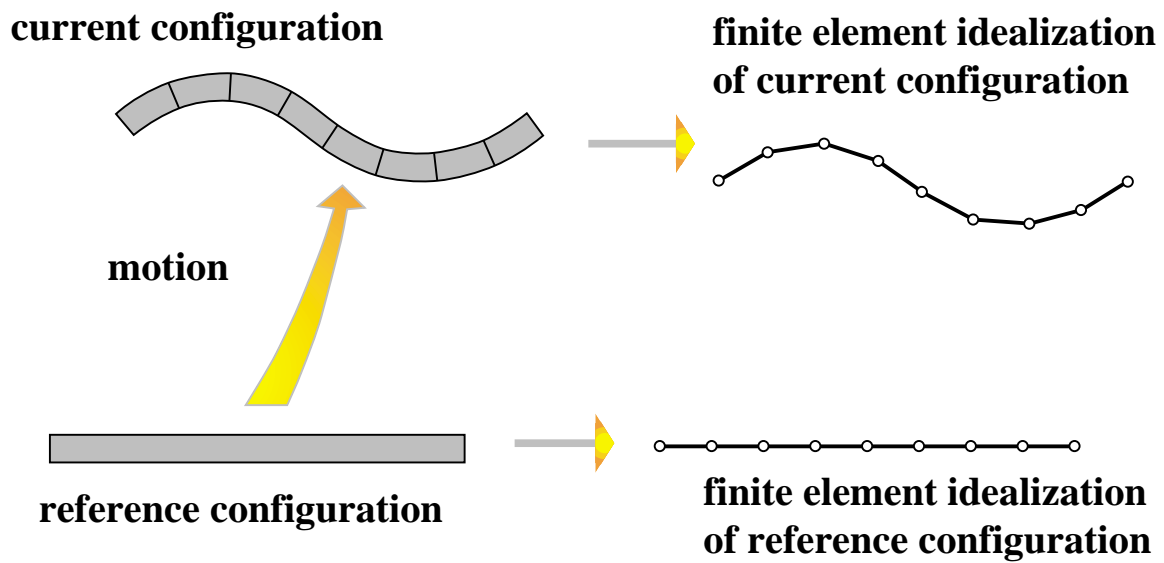
# TL Beam Element Application



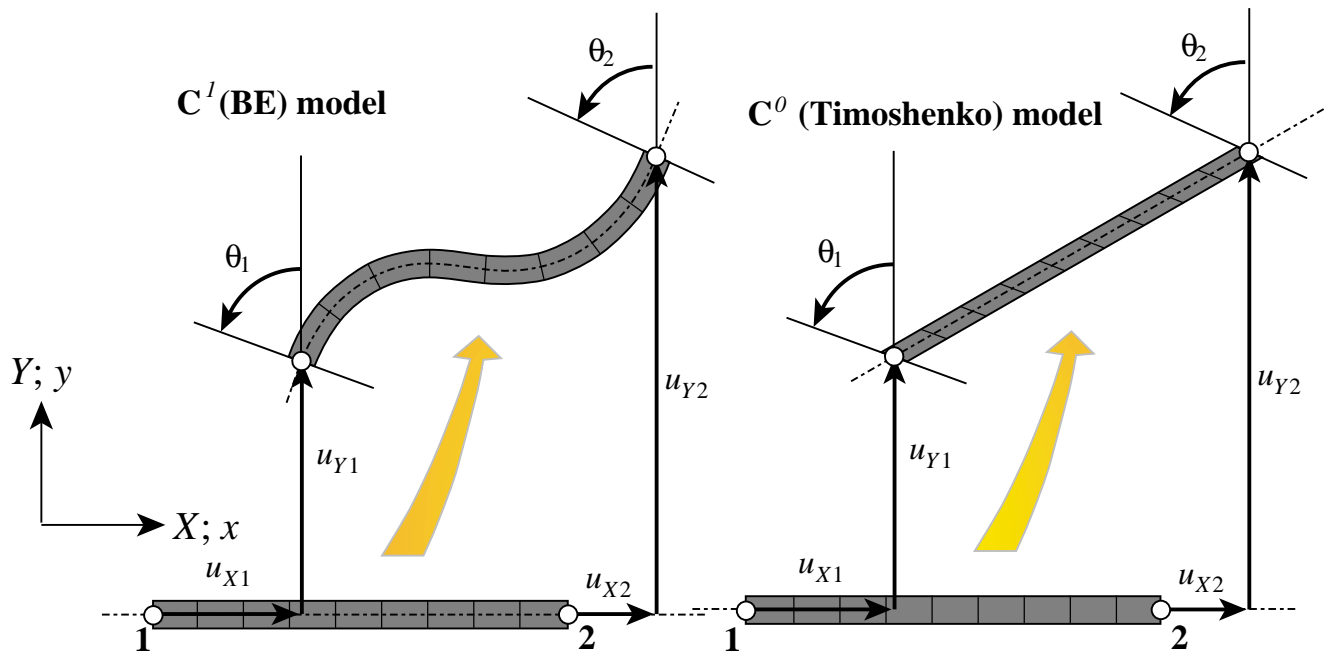
# Beam Terminology



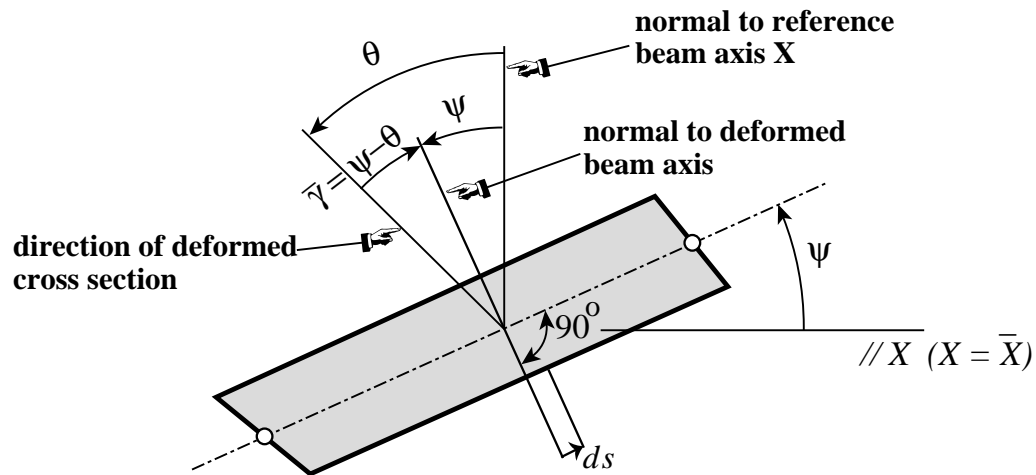
## Reduction To One Dimensional Model



# Beam Mathematical Models



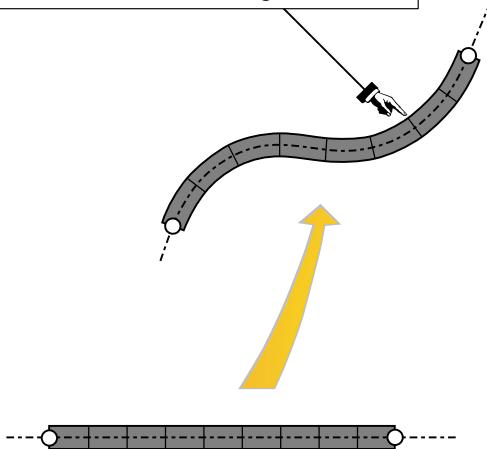
# Timoshenko Beam Kinematics



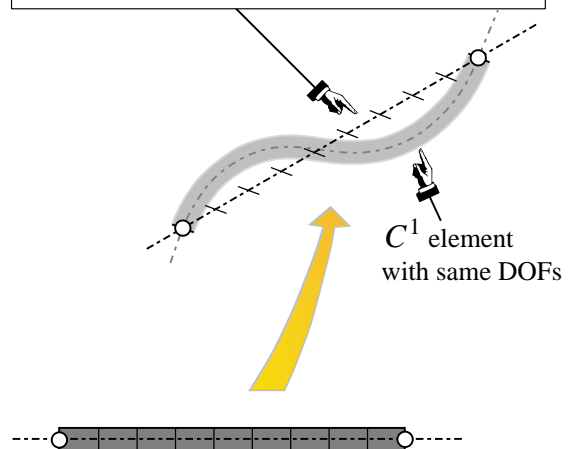
**Note: in practice  $\bar{\gamma} \ll \theta$ ; typically 0.1% or less. Magnitude of  $\bar{\gamma}$  is grossly exaggerated in the figure for visualization convenience.**

## Comparing The Two Models Over Individual Finite Element

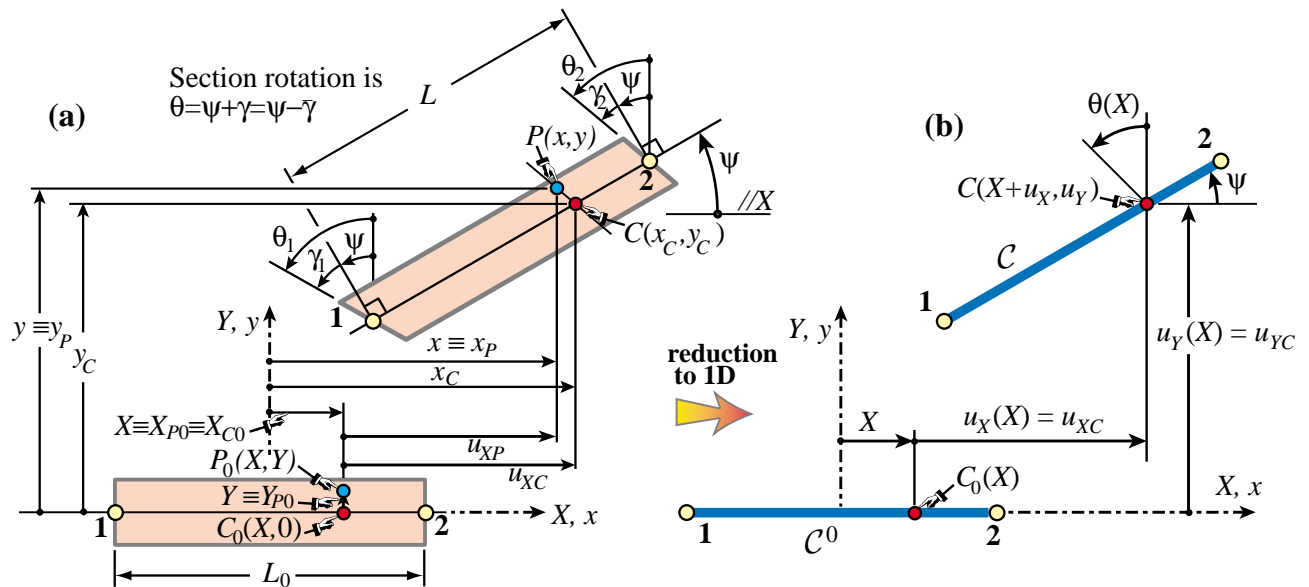
2-node  $C^1$  (cubic) element  
for Euler-Bernoulli beam model:  
plane sections remain plane and  
*normal* to deformed longitudinal axis



2-node  $C^0$  linear-displacement-and-rotations  
element for Timoshenko beam model:  
plane sections remain plane but *not*  
*normal* to deformed longitudinal axis



## Element Kinematics: Reduction from 2D to 1D

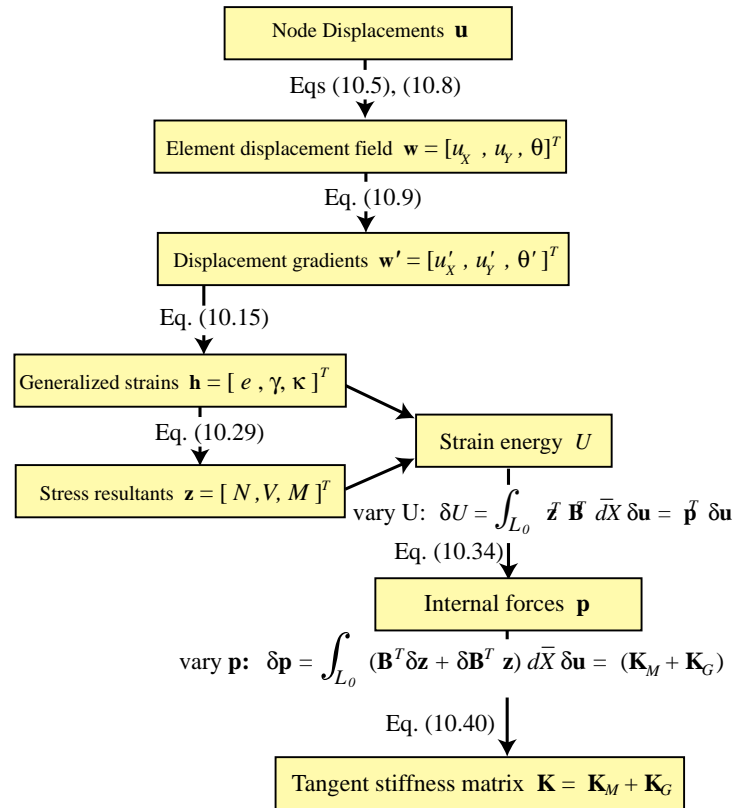




## Degrees of Freedom of Beam Element (they are model independent)

$$\mathbf{u} = \begin{bmatrix} u_{X1} \\ u_{Y1} \\ \theta_1 \\ u_{X2} \\ u_{Y2} \\ \theta_2 \end{bmatrix} \quad \mathbf{f} = \begin{bmatrix} f_{X1} \\ f_{Y1} \\ f_{\theta1} \\ f_{X2} \\ f_{Y2} \\ f_{\theta2} \end{bmatrix}$$

# Element Formulation Path



# The Longest Trip Begins With The First Step

## Element motion in Lagrangian description

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} X + u_X - Y \sin \theta \\ u_Y + Y \cos \theta \end{bmatrix}$$

## Extended displacement vector

$$\mathbf{w} = \begin{bmatrix} u_X(X) \\ u_Y(X) \\ \theta(X) \end{bmatrix}, \quad \mathbf{w}' = \frac{d\mathbf{w}}{dX} = \begin{bmatrix} du_X/dX \\ du_Y/dX \\ d\theta/dX \end{bmatrix} = \begin{bmatrix} u'_X \\ u'_Y \\ \theta' \end{bmatrix}$$

## Isoparametric Interpolation Of Extended Displacements

$$\mathbf{w} = \begin{bmatrix} u_X(X) \\ u_Y(X) \\ \theta(X) \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 - \xi & 0 & 0 & 1 + \xi & 0 & 0 \\ 0 & 1 - \xi & 0 & 0 & 1 + \xi & 0 \\ 0 & 0 & 1 - \xi & 0 & 0 & 1 + \xi \end{bmatrix} \begin{bmatrix} u_{X1} \\ u_{Y1} \\ \theta_1 \\ u_{X2} \\ u_{Y2} \\ \theta_2 \end{bmatrix} = \mathbf{N} \mathbf{u},$$

$$\mathbf{w}' = \begin{bmatrix} u'_X \\ u'_Y \\ \theta' \end{bmatrix} = \frac{1}{L_0} \begin{bmatrix} -1 & 0 & 0 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 & 1 & 0 \\ 0 & 0 & -1 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} u_{X1} \\ u_{Y1} \\ \theta_1 \\ u_{X2} \\ u_{Y2} \\ \theta_2 \end{bmatrix} = \mathbf{N}' \mathbf{u} = \mathbf{G} \mathbf{u}.$$

# Deformation Gradient and Displacement Gradient Tensors

$$\mathbf{F} = \begin{bmatrix} \frac{\partial x}{\partial X} & \frac{\partial x}{\partial Y} \\ \frac{\partial y}{\partial X} & \frac{\partial y}{\partial Y} \end{bmatrix} = \begin{bmatrix} 1 + u'_X - Y \kappa \cos \theta & -\sin \theta \\ u'_Y - Y \kappa \sin \theta & \cos \theta \end{bmatrix}$$

$$\mathbf{G}_F = \mathbf{F} - \mathbf{I} = \begin{bmatrix} u'_X - Y \kappa \cos \theta & -\sin \theta \\ u'_Y - Y \kappa \sin \theta & \cos \theta - 1 \end{bmatrix}$$

# Green-Lagrange Strains

**Tensor form:**

$$\underline{\mathbf{e}} = \begin{bmatrix} e_{XX} & e_{XY} \\ e_{YX} & e_{YY} \end{bmatrix} = \frac{1}{2}(\mathbf{F}^T \mathbf{F} - \mathbf{I}) = \frac{1}{2}(\mathbf{G}_F + \mathbf{G}_F^T) + \frac{1}{2}\mathbf{G}_F^T \mathbf{G}_F$$

**Tensor component form:**

$$\underline{\mathbf{e}} = \frac{1}{2} \begin{bmatrix} 2(u'_X - Y\kappa \cos \theta) + (u'_X - Y\kappa \cos \theta)^2 + (u'_Y - Y\kappa \sin \theta)^2 & -(1 + u'_X) \sin \theta + u'_Y \cos \theta \\ -(1 + u'_X) \sin \theta + u'_Y \cos \theta & 0 \end{bmatrix}$$

**These are geometrically exact, but unwieldy.**

**Simplifications are necessary - those are detailed in Chapter 10.**

## Simplified Strain Field

**Axial strain** is linearized with respect to displacements using a technique called **polar decomposition**. The result is

$$\mathbf{e} = \begin{bmatrix} e_1 \\ e_2 \end{bmatrix} = \begin{bmatrix} e_{XX} \\ 2e_{XY} \end{bmatrix} = \begin{bmatrix} (1 + u'_X) \cos \theta + u'_Y \sin \theta - Y\theta' - 1 \\ -(1 + u'_X) \sin \theta + u'_Y \cos \theta \end{bmatrix} = \begin{bmatrix} e - Y\kappa \\ \gamma \end{bmatrix}$$

We can characterize axial strains, shear strains and bending strains

$$e = (1 + u'_X) \cos \theta + u'_Y \sin \theta - 1 \quad \gamma = -(1 + u'_X) \sin \theta + u'_Y \cos \theta \quad \kappa = \theta'$$

Collected in the generalized strain vector

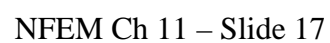
$$\mathbf{h} = \begin{bmatrix} e \\ \gamma \\ \kappa \end{bmatrix}$$

## Geometrically Invariant Form of Axial and Shear Strains

$$1 + e = s' \cos(\theta - \psi) = \frac{L \cos \bar{\gamma}}{L_0}$$
$$\gamma = -s' \sin(\theta - \psi) = \frac{L \sin \bar{\gamma}}{L_0}$$

Using this invariant form we can pass to a beam element **arbitrarily oriented** in the reference configuration





## What Remains To Be Done

**The remaining steps include:**

**PK2 stresses and stress resultants**

**Strain energy**

**Internal force**

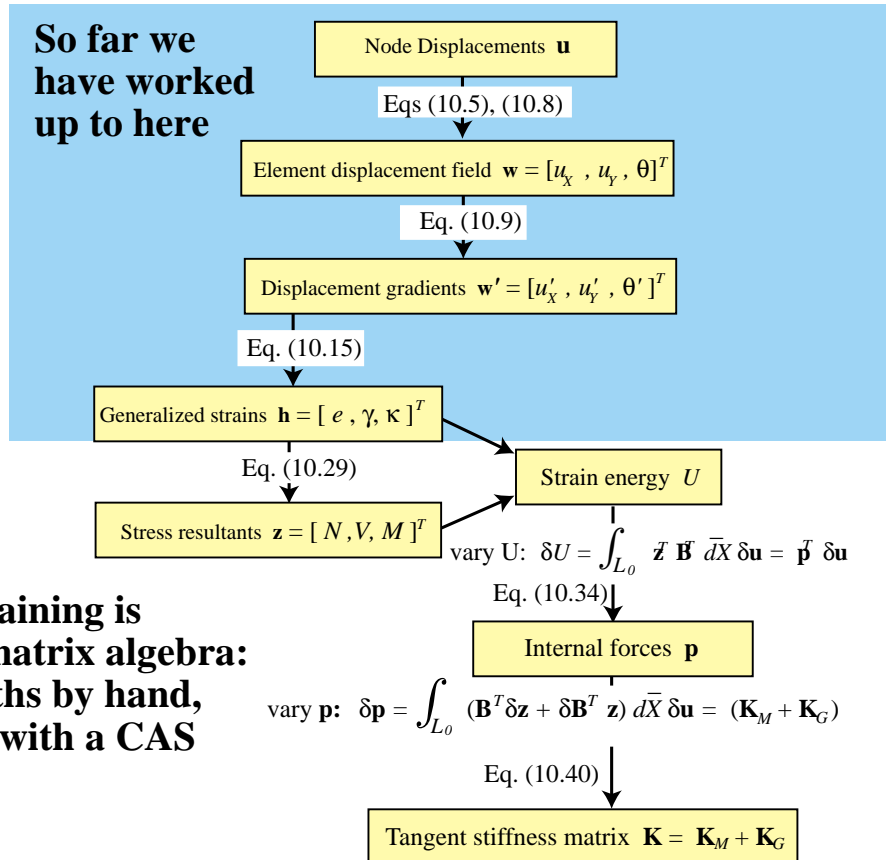
**Tangent stiffness: material + geometric**

**Improving element performance**

**There is a lot of algebra left, even for this relatively simple element. A computer algebra system (CAS) can really help, especially to avoid algebra errors**

# Flowcharted Steps (again)

So far we  
have worked  
up to here



The remaining is  
a lot of matrix algebra:  
1-2 months by hand,  
~1 week with a CAS

# Energy and Internal Forces

Nonlinear FEM

Strain energy of beam element

$$U = \int_{L_0} (N^0 e + \frac{1}{2} E A_0 e^2) d\bar{X} + \int_{L_0} (V^0 \gamma + \frac{1}{2} G A_0 \gamma^2) d\bar{X} + \int_{L_0} (M^0 \kappa + \frac{1}{2} E I_0 \kappa^2) d\bar{X}$$

Vary to get internal force

$$\delta U = \int_{L_0} (N \delta e + V \delta \gamma + M \delta \kappa) d\bar{X} = \int_{L_0} \mathbf{z}^T \delta \mathbf{h} d\bar{X} = \int_{L_0} \mathbf{z}^T \mathbf{B} d\bar{X} \delta \mathbf{u}.$$

$$\mathbf{p} = \int_{L_0} \mathbf{B}^T \mathbf{z} d\bar{X}.$$

Evaluate by one point Gauss quadrature:

$$\mathbf{B}_m = \mathbf{B}|_{\xi=0} = \frac{1}{L_0} \begin{bmatrix} -c_m & -s_m & -\frac{1}{2}L_0\gamma_m & c_m & s_m & -\frac{1}{2}L_0\gamma_m \\ s_m & -c_m & \frac{1}{2}L_0(1+e_m) & s_m & -c_m & \frac{1}{2}L_0(1+e_m) \\ 0 & 0 & -1 & 0 & 0 & 1 \end{bmatrix}$$

$$\mathbf{p} = L_0 \mathbf{B}_m^T \mathbf{z} = \begin{bmatrix} -c_m & -s_m & \frac{1}{2}L_0\gamma_m & c_m & s_m & -\frac{1}{2}L_0\gamma_m \\ s_m & -c_m & -\frac{1}{2}L_0(1+e_m) & -s_m & c_m & \frac{1}{2}L_0(1+e_m) \\ 0 & 0 & -1 & 0 & 0 & 1 \end{bmatrix}^T \begin{bmatrix} N \\ V \\ M \end{bmatrix}$$

# Tangent Stiffness: Material Component

Vary internal force to get tangent stiffness

$$\delta \mathbf{p} = \int_{L_0} (\mathbf{B}^T \delta \mathbf{z} + \delta \mathbf{B}^T \mathbf{z}) d\bar{X} = (\mathbf{K}_M + \mathbf{K}_G) \delta \mathbf{u} = \mathbf{K} \delta \mathbf{u}.$$

Material stiffness matrix

$$\mathbf{K}_M = \int_{L_0} \mathbf{B}^T \mathbf{S} \mathbf{B} d\bar{X} \quad \delta \mathbf{z} = \begin{bmatrix} \delta N \\ \delta V \\ \delta M \end{bmatrix} = \begin{bmatrix} E A_0 & 0 & 0 \\ 0 & G A_0 & 0 \\ 0 & 0 & E I_0 \end{bmatrix} \begin{bmatrix} \delta e \\ \delta \gamma \\ \delta \kappa \end{bmatrix} = \mathbf{S} \delta \mathbf{h}$$

Evaluate by one point Gauss quadrature:

$$\mathbf{K}_M = \int_{L_0} \mathbf{B}_m^T \mathbf{S} \mathbf{B}_m d\bar{X} = \mathbf{K}_M^a + \mathbf{K}_M^b + \mathbf{K}_M^s$$

# Material Stiffness Subcomponents

Nonlinear FEM

Material stiffness subcomponents due to axial, bending and shear:

$$\begin{aligned}
 \mathbf{K}_M^a &= \frac{E A_0}{L_0} \begin{bmatrix} c_m^2 & c_m s_m & -c_m \gamma_m L_0/2 & -c_m^2 & -c_m s_m & -c_m \gamma_m L_0/2 \\ c_m s_m & s_m^2 & -\gamma_m L_0 s_m/2 & -c_m s_m & -s_m^2 & -\gamma_m L_0 s_m/2 \\ -c_m \gamma_m L_0/2 & -\gamma_m L_0 s_m/2 & \gamma_m^2 L_0^2/4 & c_m \gamma_m L_0/2 & \gamma_m L_0 s_m/2 & \gamma_m^2 L_0^2/4 \\ -c_m^2 & -c_m s_m & c_m \gamma_m L_0/2 & c_m^2 & c_m s_m & c_m \gamma_m L_0/2 \\ -c_m s_m & -s_m^2 & \gamma_m L_0 s_m/2 & c_m s_m & s_m^2 & \gamma_m L_0 s_m/2 \\ -c_m \gamma_m L_0/2 & -\gamma_m L_0 s_m/2 & \gamma_m^2 L_0^2/4 & c_m \gamma_m L_0/2 & \gamma_m L_0 s_m/2 & \gamma_m^2 L_0^2/4 \end{bmatrix} \\
 \mathbf{K}_M^b &= \frac{E I_0}{L_0} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 1 \end{bmatrix} \\
 \mathbf{K}_M^s &= \frac{G A_0}{L_0} \begin{bmatrix} s_m^2 & -c_m s_m & -a_1 L_0 s_m/2 & -s_m^2 & c_m s_m & -a_1 L_0 s_m/2 \\ -c_m s_m & c_m^2 & c_m a_1 L_0/2 & c_m s_m & -c_m^2 & c_m a_1 L_0/2 \\ -a_1 L_0 s_m/2 & c_m a_1 L_0/2 & a_1^2 L_0^2/4 & a_1 L_0 s_m/2 & -c_m a_1 L_0/2 & a_1^2 L_0^2/4 \\ -s_m^2 & c_m s_m & a_1 L_0 s_m/2 & s_m^2 & -c_m s_m & a_1 L_0 s_m/2 \\ c_m s_m & -c_m^2 & -c_m a_1 L_0/2 & -c_m s_m & c_m^2 & -c_m a_1 L_0/2 \\ -a_1 L_0 s_m/2 & c_m a_1 L_0/2 & a_1^2 L_0^2/4 & a_1 L_0 s_m/2 & -c_m a_1 L_0/2 & a_1^2 L_0^2/4 \end{bmatrix}
 \end{aligned}$$

in which  $a_1 = 1 + e_m$ .

# Material Stiffness "Unlocking"

Nonlinear FEM

Eliminating "shear locking" by MacNeal's  
Residual Bending Flexibility (RBF) device:

$GA_0$  is formally replaced by  $12EI_0/L_0^2$  in shear material stiffness

Bending and shear components merge into a modified bending stiffness:

$$\hat{\mathbf{K}}_M^b = \frac{EI}{L_0^3} \begin{bmatrix} 12s_m^2 & -12c_ms_m & 6a_1L_0s_m & -12s_m^2 & 12c_ms_m & 6a_1L_0s_m \\ -12c_ms_m & 12c_m^2 & -6c_ma_1L_0 & 12c_ms_m & -12c_m^2 & -6c_ma_1L_0 \\ 6a_1L_0s_m & -6c_ma_1L_0 & a_2L_0^2 & -6a_1L_0s_m & 6c_ma_1L_0 & a_3L_0^2 \\ -12s_m^2 & 12c_ms_m & -6a_1L_0s_m & 12s_m^2 & -12c_ms_m & -6a_1L_0s_m \\ 12c_ms_m & -12c_m^2 & 6c_ma_1L_0 & -12c_ms_m & 12c_m^2 & 6c_ma_1L_0 \\ 6a_1L_0s_m & -6c_ma_1L_0 & a_3L_0^2 & -6a_1L_0s_m & 6c_ma_1L_0 & a_2L_0^2 \end{bmatrix}$$

in which  $a_1 = 1 + e_m$ ,  $a_2 = 4 + 6e_m + 3e_m^2$  and  $a_3 = 2 + 6e_m + 3e_m^2$ .

# Geometric Stiffness

The geometric stiffness comes from the variation of strain-displacement matrix **B** while stress resultants (axial forces, shear forces and bending moments) are kept fixed:

$$\mathbf{K}_G = \int_{L_0} \delta \mathbf{B}^T \mathbf{z} d\bar{X}$$

After tons of algebra (see Notes) one arrives at

$$\mathbf{K}_G = \int_{L_0} (\mathbf{W}_N N + \mathbf{W}_V V) d\bar{X} = \mathbf{K}_{GN} + \mathbf{K}_{GV}$$

which evaluated by one point Gauss rule gives the closed form

$$\mathbf{K}_G = \frac{N}{2} \begin{bmatrix} 0 & 0 & s_m & 0 & 0 & s_m \\ 0 & 0 & -c_m & 0 & 0 & -c_m \\ s_m & -c_m & -\frac{1}{2}L_0(1+e_m) & -s_m & c_m & -\frac{1}{2}L_0(1+e_m) \\ 0 & 0 & -s_m & 0 & 0 & -s_m \\ 0 & 0 & c_m & 0 & 0 & c_m \\ s_m & -c_m & -\frac{1}{2}L_0(1+e_m) & -s_m & c_m & -\frac{1}{2}L_0(1+e_m) \end{bmatrix}$$

$$+ \frac{V}{2} \begin{bmatrix} 0 & 0 & c_m & 0 & 0 & c_m \\ 0 & 0 & s_m & 0 & 0 & s_m \\ c_m & s_m & -\frac{1}{2}L_0\gamma_m & -c_m & -s_m & -\frac{1}{2}L_0\gamma_m \\ 0 & 0 & -c_m & 0 & 0 & -c_m \\ 0 & 0 & -s_m & 0 & 0 & -s_m \\ c_m & s_m & -\frac{1}{2}L_0\gamma_m & -c_m & -s_m & -\frac{1}{2}L_0\gamma_m \end{bmatrix}$$