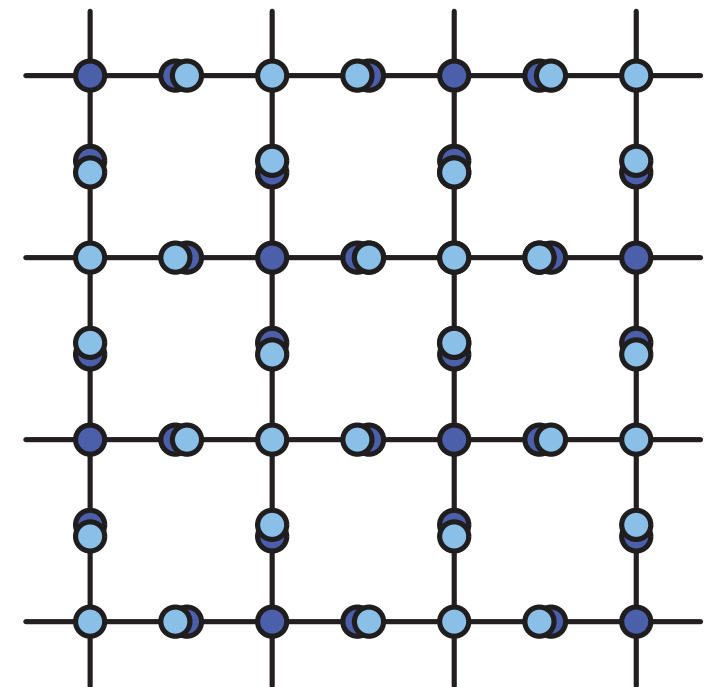


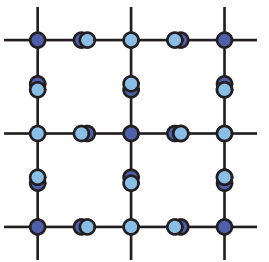
Engineering artificial gauge fields with ultracold atoms - part II

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Ludwig-Maximilians Universität München
Munich Center for
Quantum Science & Technology

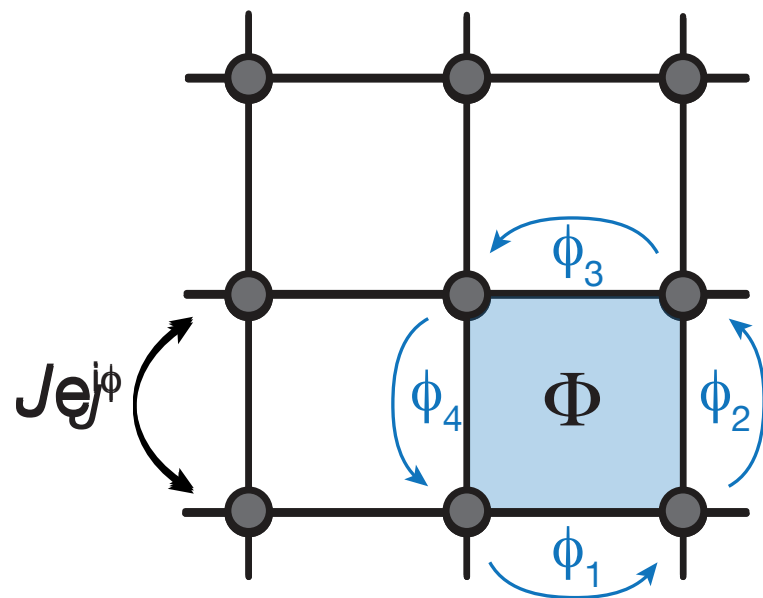


Artificial magnetic fields



Non-interacting
lattice Hamiltonian:

$$\hat{H} = - \sum_{\langle i,j \rangle} J_{ij} \hat{a}_i^\dagger \hat{a}_j + \text{h.c.}$$



Charged particles in magnetic field
→ acquire geometric phase

Peierls substitution: $J_{ij} \rightarrow J_{ij} e^{i\phi_j}$

$$\phi_j = \frac{q}{\hbar} \int_{x_j}^{x_i} \mathbf{A} d\mathbf{l}, \quad \mathbf{B} = \nabla \times \mathbf{A}$$

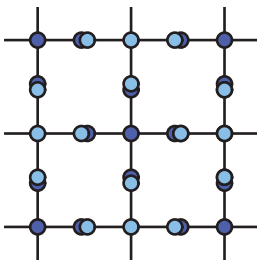
Phase around
closed loop:

$$\Phi = \sum_{\square} \phi_j = 2\pi \frac{\Phi_B}{\Phi_0}$$

Φ_B : magn. flux

$\Phi_0 = h/q$: magn. flux quantum

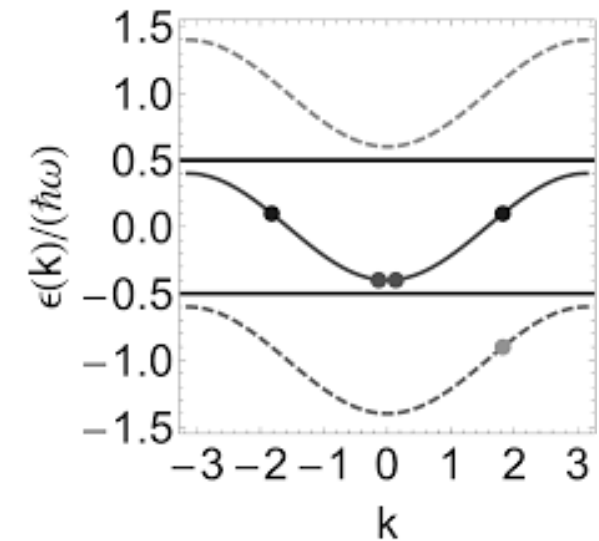
Floquet engineering



Basic idea:

- Time-periodic driven Hamiltonian

$$\hat{H}(t) = \hat{H}(t + T) \quad T: \text{driving cycle}$$

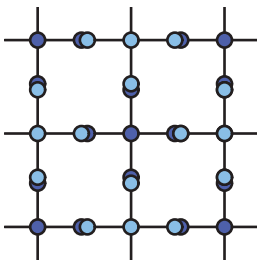


- *Stroboscopic* time evolution governed by *time-independent Floquet Hamiltonian* $\hat{H}^F$

$$\hat{U}(T, 0) = \exp \left(-\frac{i}{\hbar} T \hat{H}^F \right)$$

⇒ Engineer **Floquet Hamiltonian** with desired properties!

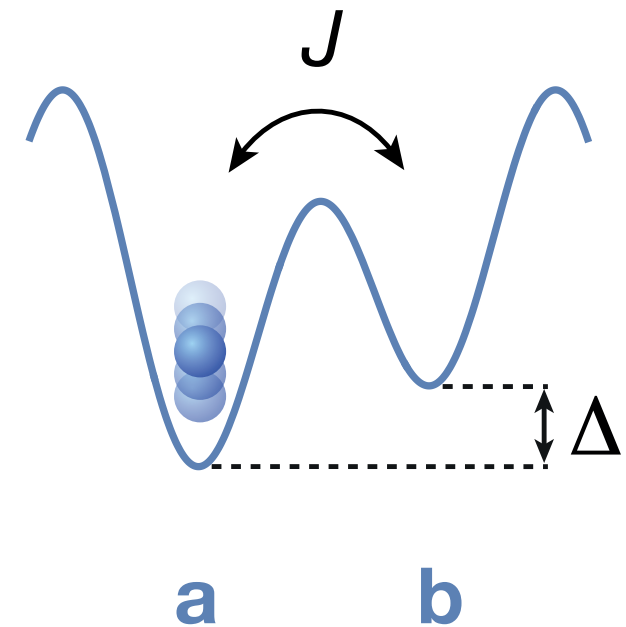
Laser-assisted tunneling



Multi-photon processes:

- Tilted double-well potential

$$\hat{H} = -J \left(\hat{a}^\dagger \hat{b} + \hat{b}^\dagger \hat{a} \right) + \Delta \hat{b}^\dagger \hat{b}$$



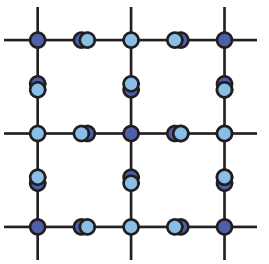
- Resonant* modulation at $\Delta_\nu = \nu \hbar \omega$

$$\hat{V}(t) = V_0 \cos(\omega t + \phi) \hat{a}^\dagger \hat{a}$$

- Restored tunneling (Floquet theory)

$$J_{\text{eff}} = J \mathcal{J}_\nu(\chi) e^{i\nu\phi} \quad \chi = \frac{V_0}{\hbar\omega}$$

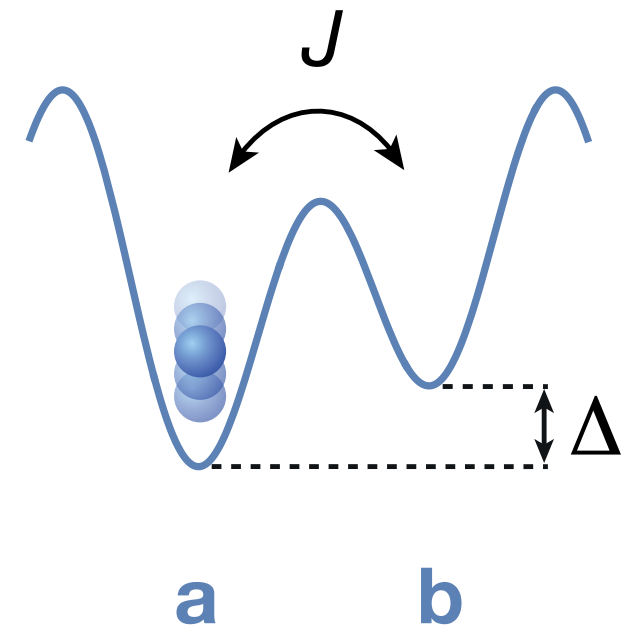
Laser-assisted tunneling



Multi-photon processes:

- Tilted double-well potential

$$\hat{H} = -J \left(\hat{a}^\dagger \hat{b} + \hat{b}^\dagger \hat{a} \right) + \Delta \hat{b}^\dagger \hat{b}$$



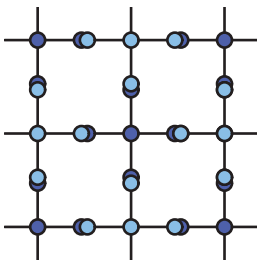
- Resonant* modulation at $\Delta_\nu = \nu \hbar \omega$

$$\hat{V}(t) = V_0 \cos(\omega t + \phi) \hat{a}^\dagger \hat{a}$$

- Reflection properties of Bessel function:

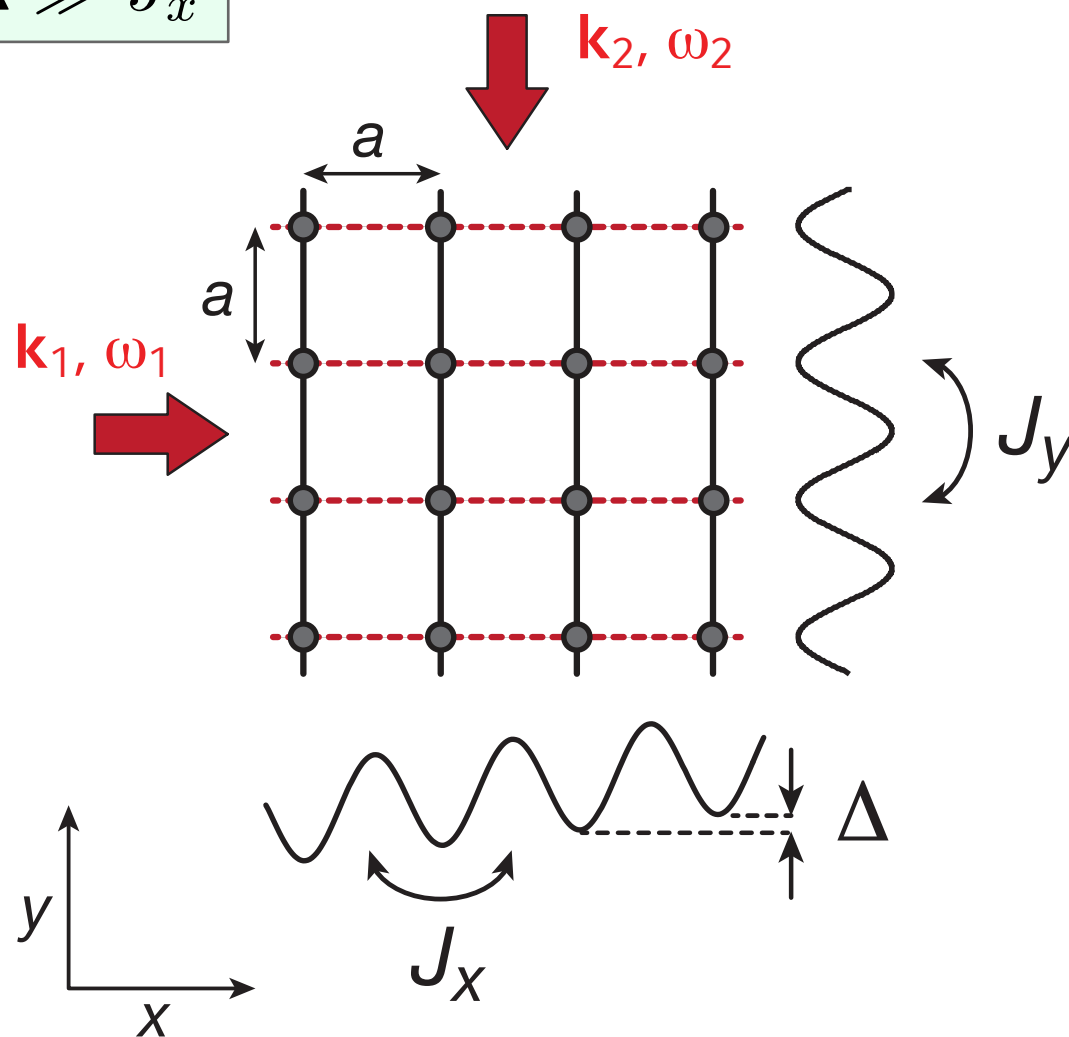
$$\mathcal{J}_{-\nu}(\chi) = (-1)^\nu \mathcal{J}_\nu(\chi)$$

Engineering complex tunneling



Tunneling inhibited along one direction
+ bare tunneling along the orthogonal direction

$$\Delta \gg J_x$$



Apply resonant modulation with
far-detuned running-wave beams:

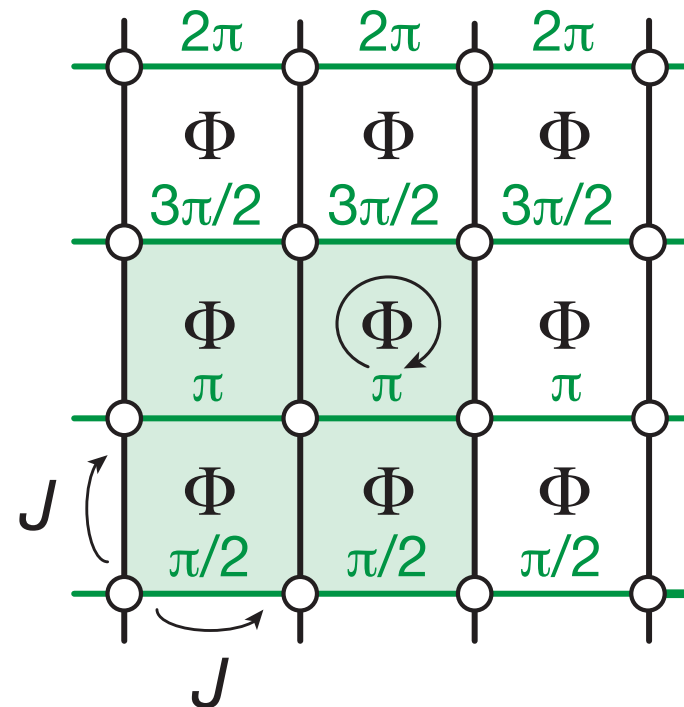
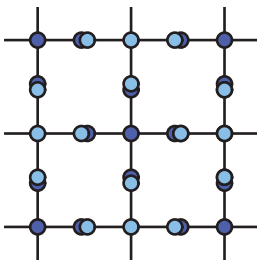
$$\omega_2 - \omega_1 = \Delta/\hbar$$

Local optical potential creates
on-site modulation:

$$V(x, y, t) = V_0 \cos(\omega t + \phi(\mathbf{r}))$$

$$\phi(\mathbf{r}) = \delta \mathbf{k} \cdot \mathbf{r}$$

Hofstadter model

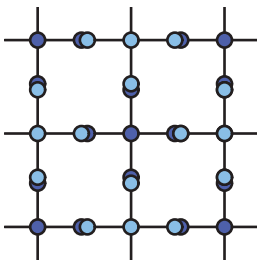


$$\hat{H} = -J \sum_{m,n} \left(e^{in\Phi} \hat{a}_{m+1,n}^\dagger \hat{a}_{m,n} + \hat{a}_{m,n+1}^\dagger \hat{a}_{m,n} + \text{h.c.} \right)$$

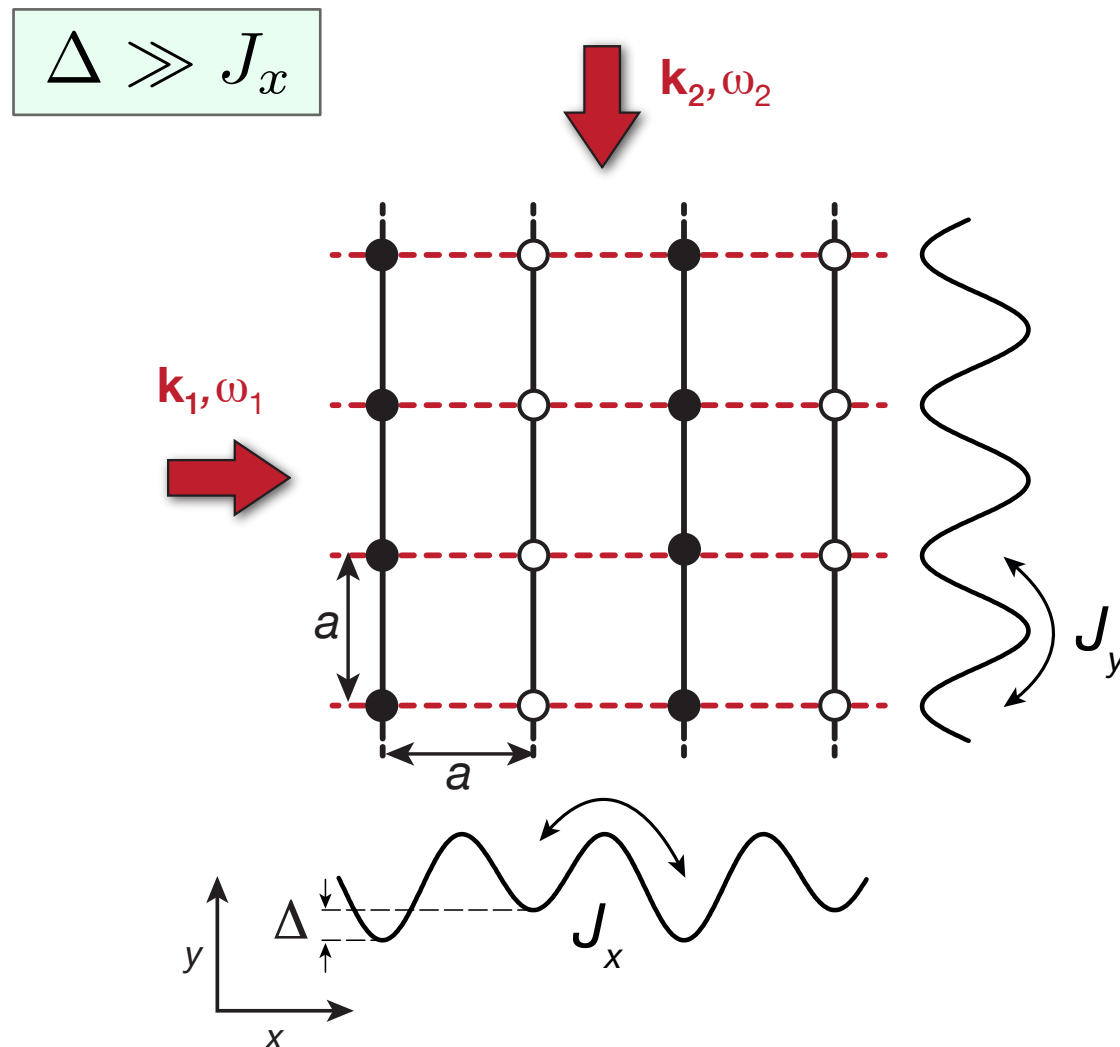
P.G. Harper, Proc. Phys. Soc., Sect.A 68, 874 (1955);

M.Y. Azbel, Zh. Eksp. Teor. Fiz. 46, 929 (1964); D.R. Hofstadter, Phys. Rev. B14, 2239 (1976)

Question



What is the flux configuration for this modulation scheme?



Resonant modulation with *far-detuned* running-wave beams:

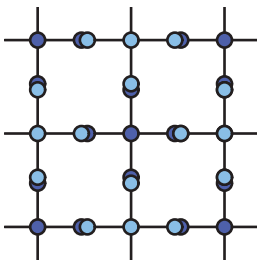
$$\omega_2 - \omega_1 = \Delta/\hbar$$

Local optical potential creates on-site modulation:

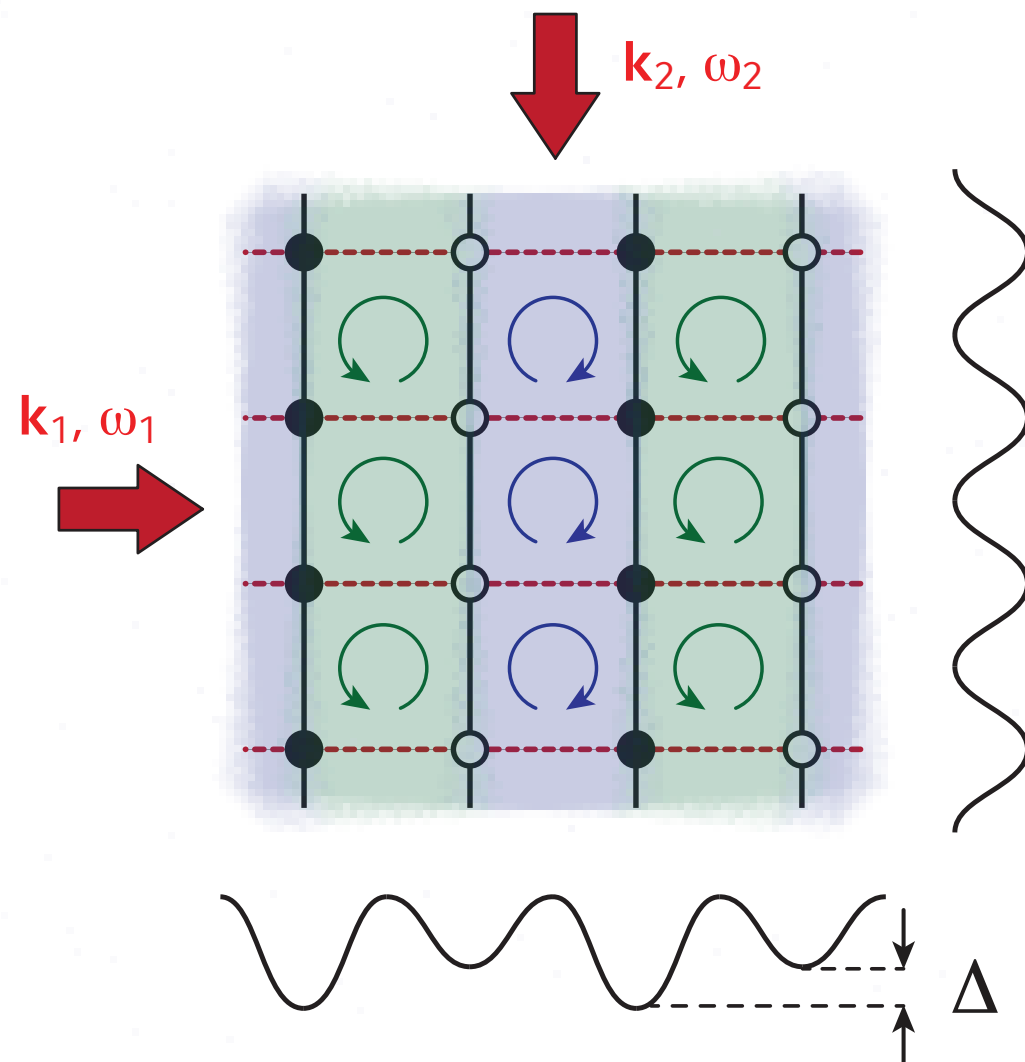
$$V(x, y, t) = V_0 \cos(\omega t + \phi(\mathbf{r}))$$

$$\phi(\mathbf{r}) = \delta \mathbf{k} \cdot \mathbf{r}$$

Answer



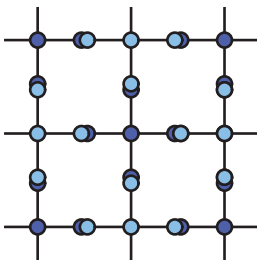
Staggered flux configuration (zero net magnetic field)



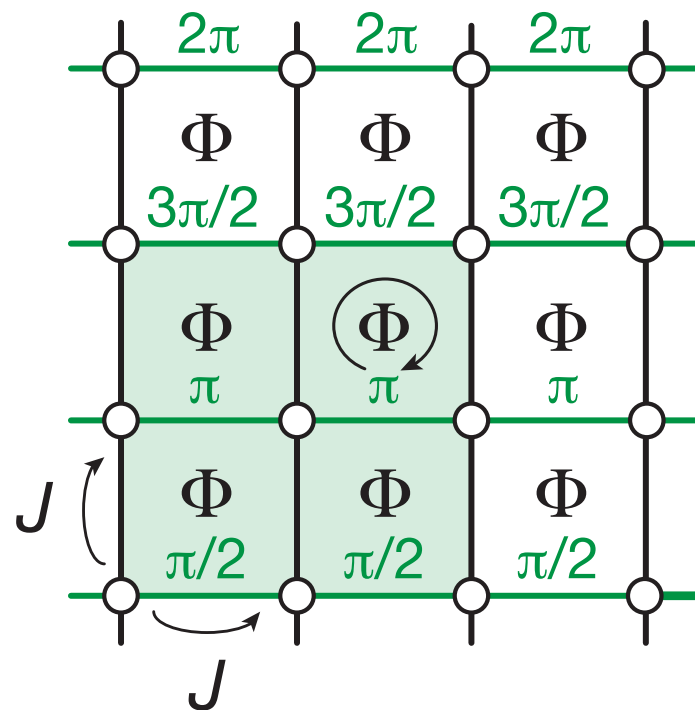
Sign of Peierls phases depends on sign of energy offset:

$$\Delta \Longleftrightarrow -\Delta$$
$$K e^{i\phi_{m,n}} \Longleftrightarrow -K e^{-i\phi_{m,n}}$$

Hofstadter model



$$\hat{H} = -J \sum_{m,n} \left(e^{in\Phi} \hat{a}_{m+1,n}^\dagger \hat{a}_{m,n} + \hat{a}_{m,n+1}^\dagger \hat{a}_{m,n} + \text{h.c.} \right)$$

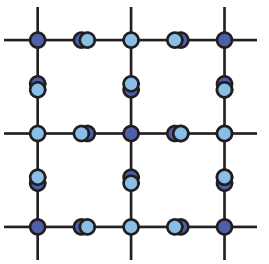


Discrete translational symmetry
no longer determined by lattice vectors

- Magnetic translation operators!
- Increased magnetic unit cell (flux 2π)

How to probe topology?

Probing band topology



Chern number:

$$\nu_\mu = \frac{1}{2\pi} \int_{\text{BZ}} \Omega_\mu d^2q$$

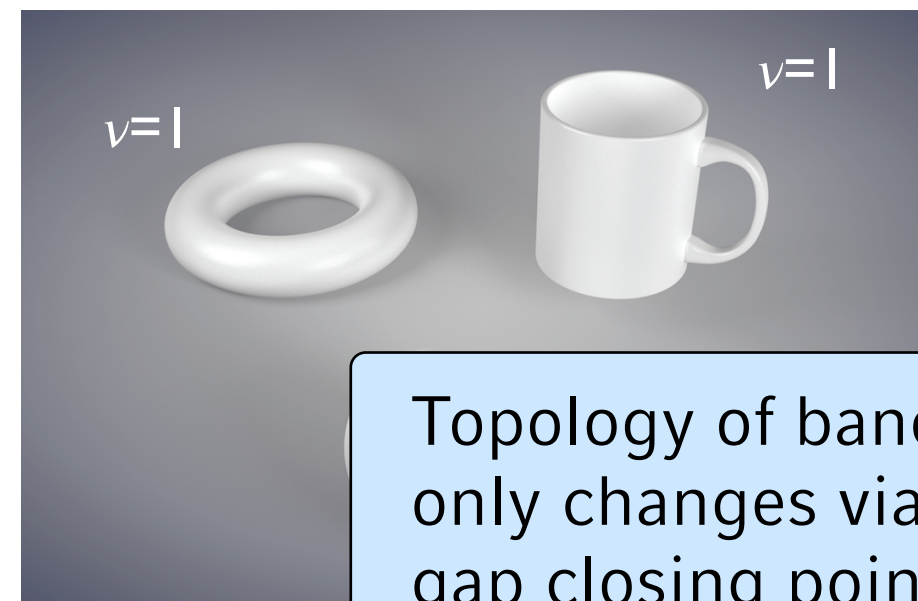
$|u_\mu(\mathbf{q})\rangle$: periodic Bloch function,
 μ : band index

Berry curvature: $\Omega_\mu = i \left(\langle \partial_{q_x} u_\mu | \partial_{q_y} u_\mu \rangle - \langle \partial_{q_y} u_\mu | \partial_{q_x} u_\mu \rangle \right)$

Gauss-Bonnet theorem:

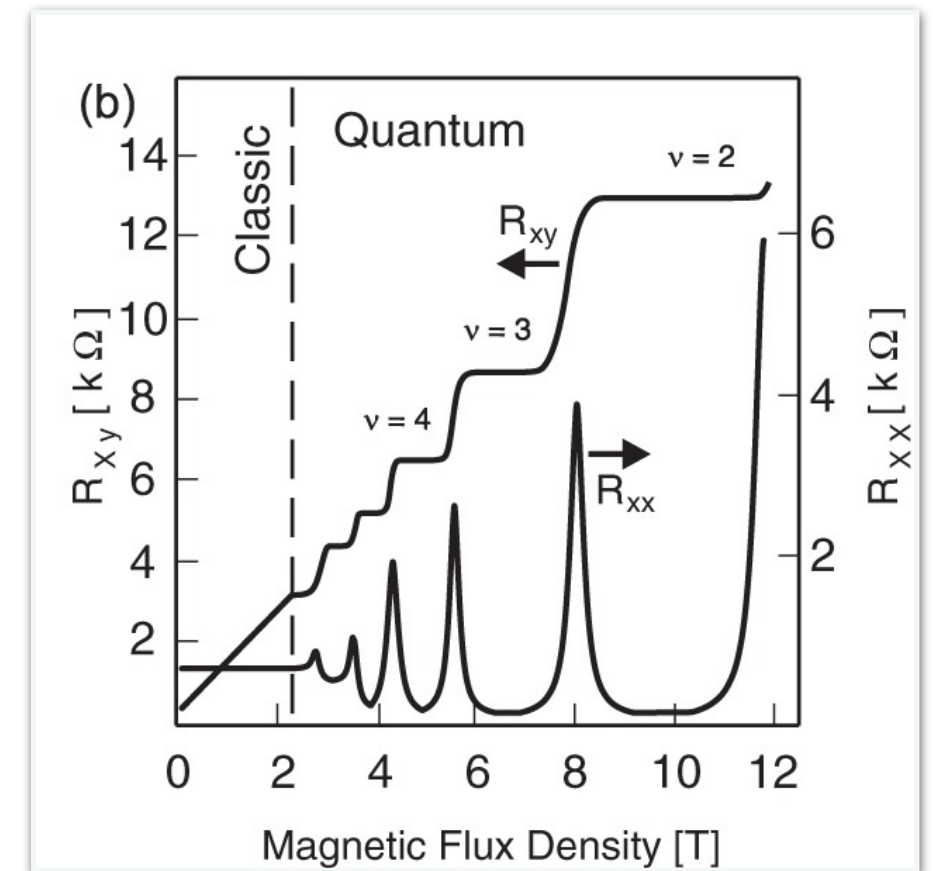
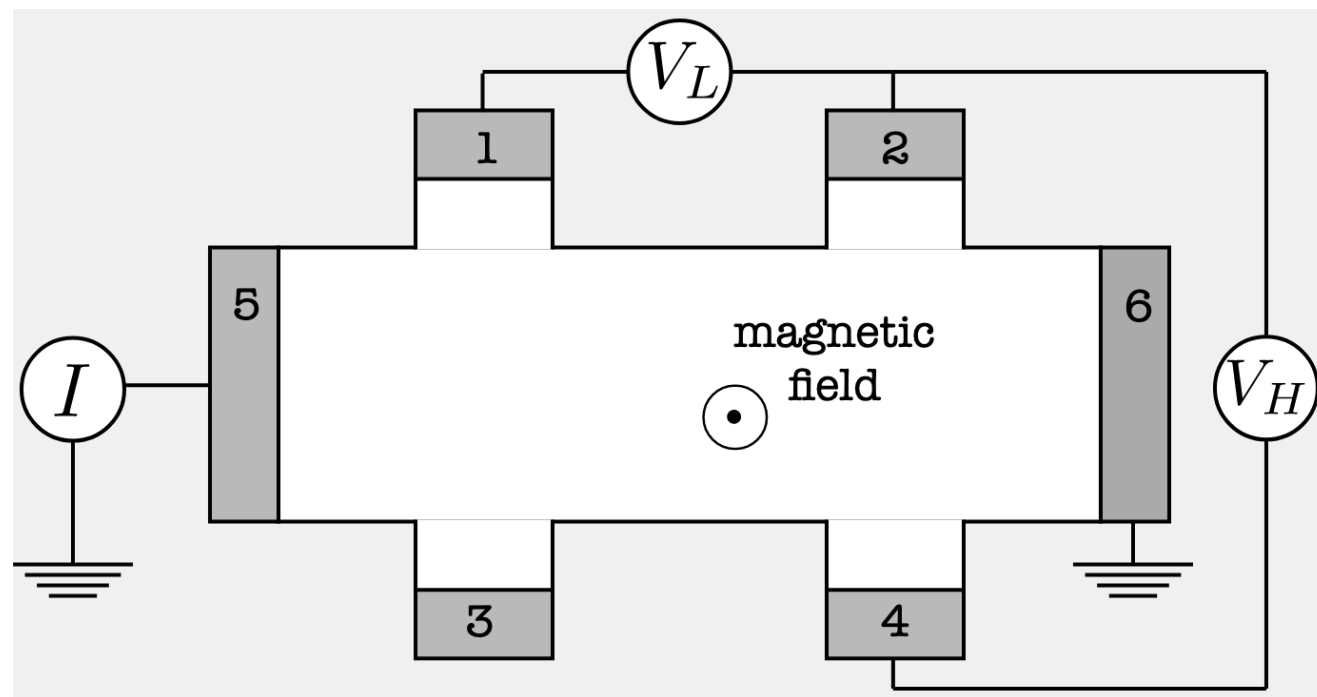
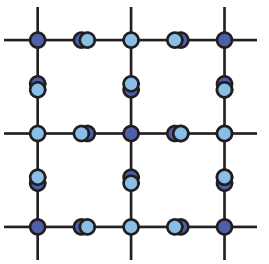
$$\frac{1}{2\pi} \int_{\mathcal{M}} \Omega dS = 2(1 - \nu)$$

local curvature



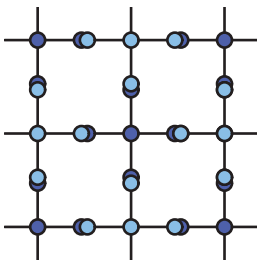
Topology of bands
only changes via
gap closing points!

Transport measurements

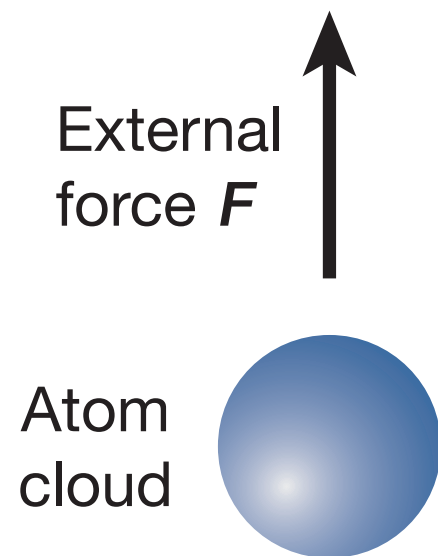


How does this work with ultracold atoms?

Hall deflection with bosonic atoms



Semiclassical dynamics:



Atoms undergo Bloch oscillations captured by **band velocity**

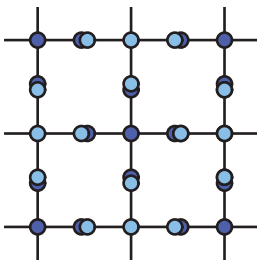
$$\mathbf{v}_{\mu}^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{k}} E_{\mu}$$

μ : band index

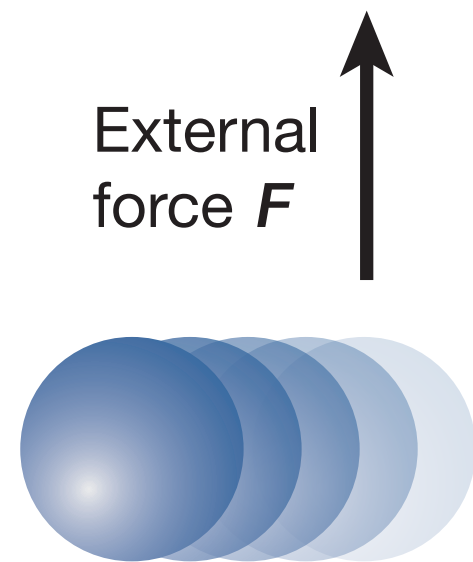
Atoms experience an **anomalous transv.** velocity prop. to Berry curvature Ω_{μ}

$$v_{\mu}^x(\mathbf{k}) = -\frac{F}{\hbar} \Omega_{\mu}(\mathbf{k})$$

Hall deflection with bosonic atoms



Semiclassical dynamics:



Atoms undergo Bloch oscillations captured by **band velocity**

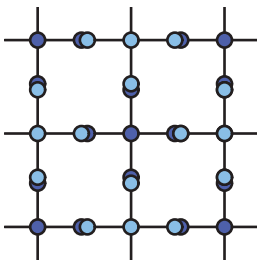
$$\mathbf{v}_{\mu}^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{k}} E_{\mu}$$

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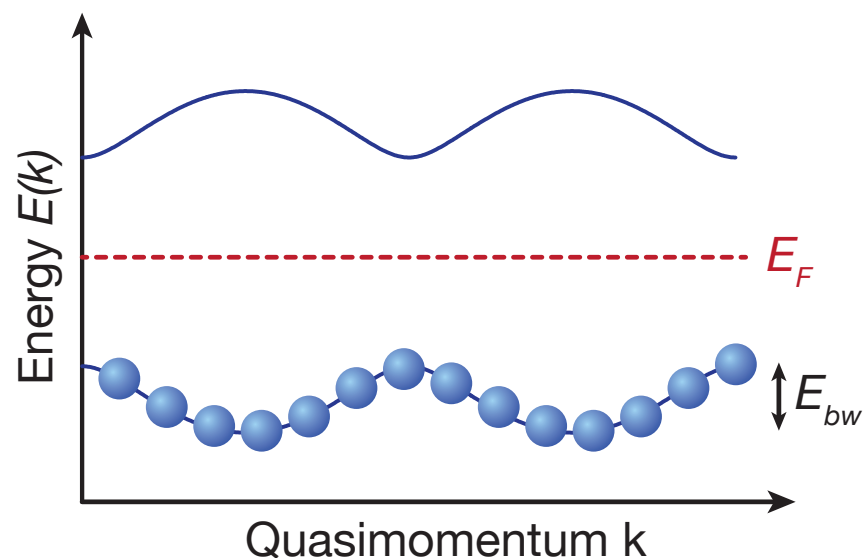
Atoms experience an **anomalous transv.** velocity prop. to Berry curvature Ω_{μ}

$$v_{\mu}^x(\mathbf{k}) = -\frac{F}{\hbar} \Omega_{\mu}(\mathbf{k})$$

Hall deflection with bosonic atoms



Deflection of the cloud for filled band:

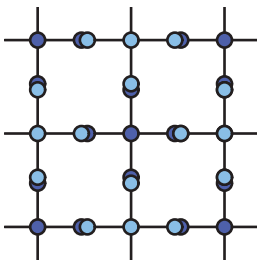


Contribution from band velocity vanishes

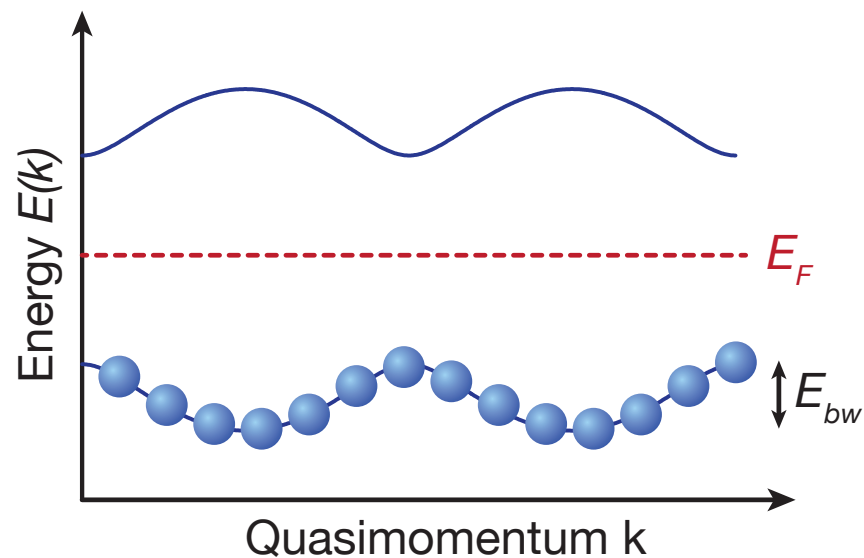
$$\int \partial E_{\mu} / \partial k_{x,y} d^2 k = 0$$

- **Fermions:** Set Fermi energy within a spectral gap
- **Bosons:** Incoherent distribution homogeneously populating the band ($k_B T \gg E_{bw}$)

Hall deflection with bosonic atoms



Deflection of the cloud for filled band:



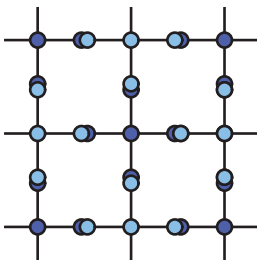
Transverse center-of-mass motion (μ th band):

$$x_{\mu}(t) = -\frac{4a^2 F}{h} \nu_{\mu} t$$

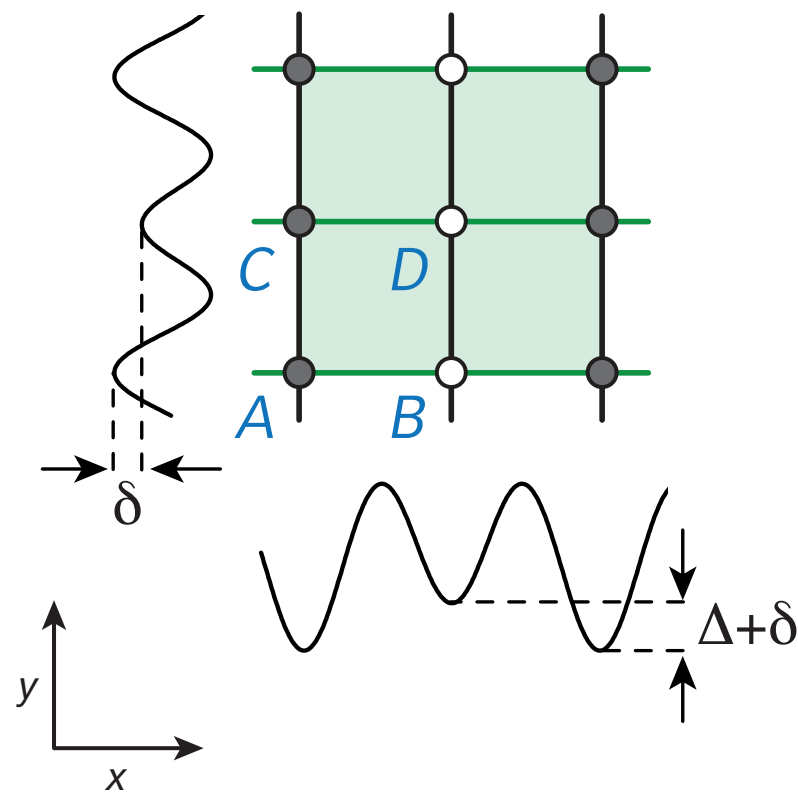
Response is proportional to Chern number of μ th band:

$$\nu_{\mu} = (1/2\pi) \int_{\text{FBZ}} \Omega_{\mu} d^2 k$$

Hall deflection with bosonic atoms



Experimental preparation protocol



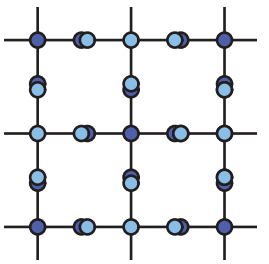
Challenge:

lowest tight-binding band
splits in the **magnetic subbands**

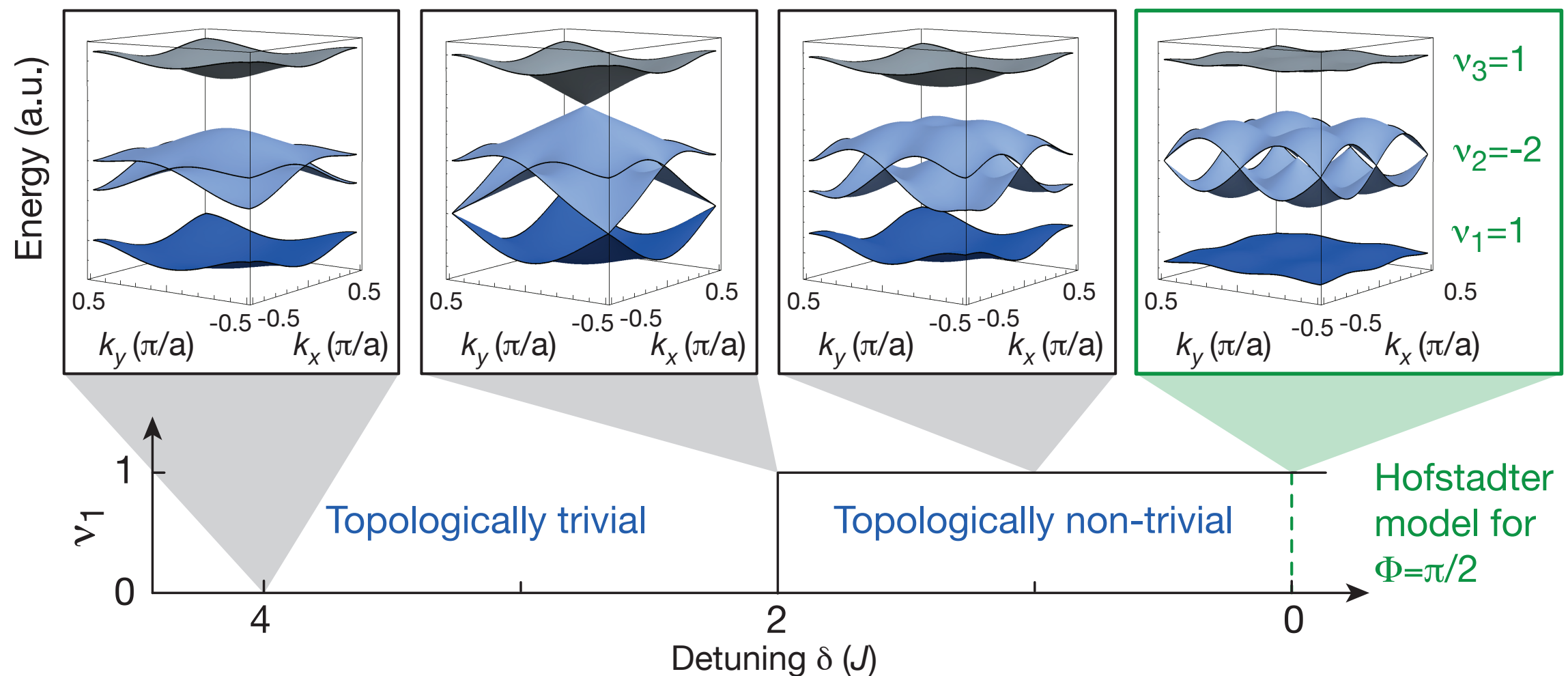
Solution:

preserve number of bands during
loading sequence

Hall deflection with bosonic atoms

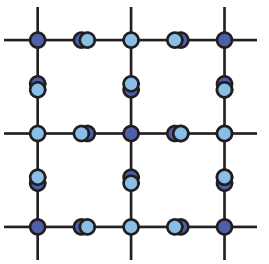


Experimental preparation protocol

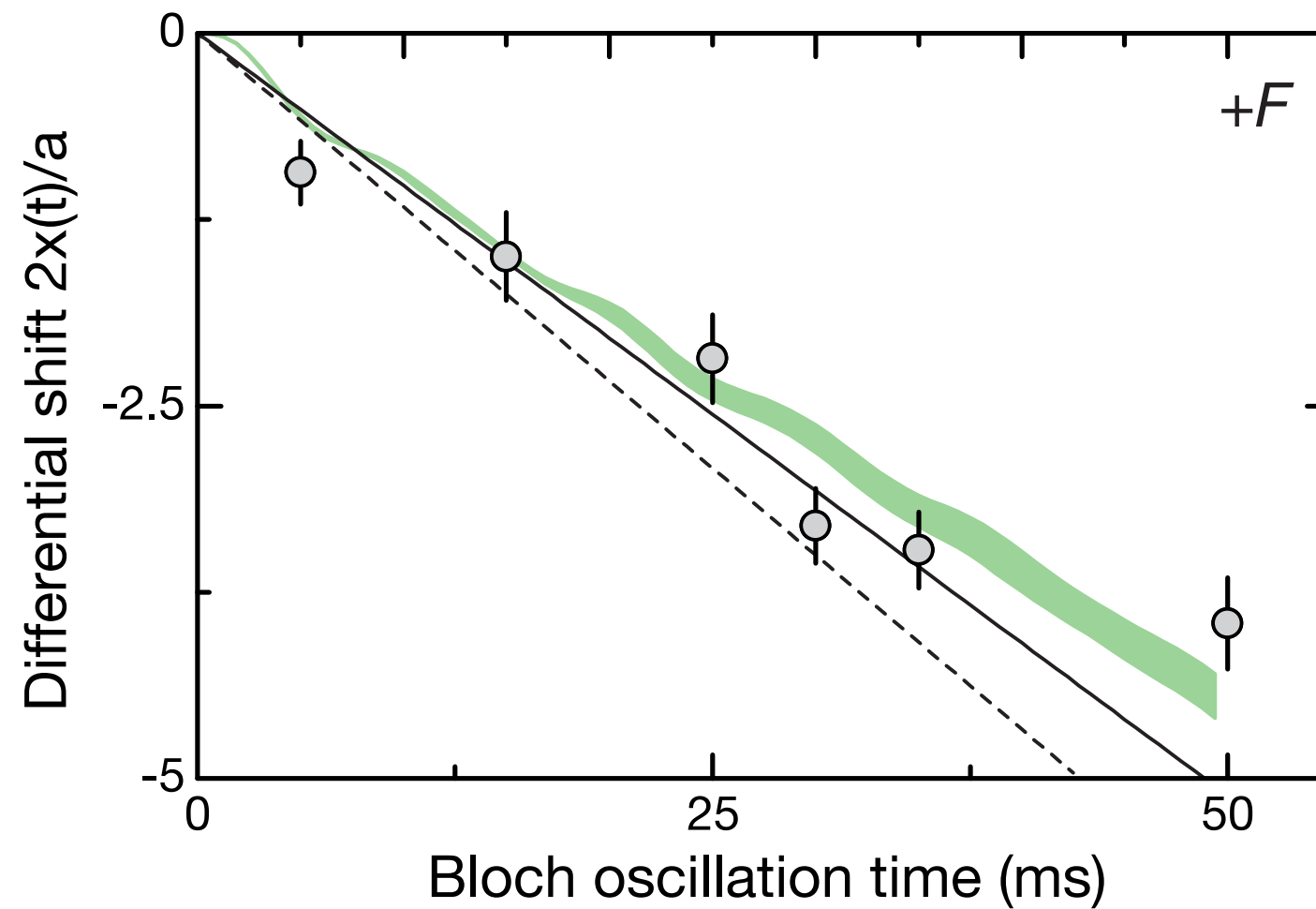


Topological phase transition (gap-closing point) at $\delta = 2J$

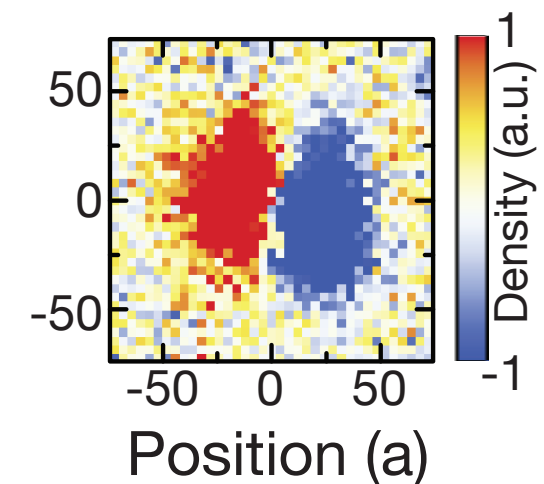
Hall deflection with bosonic atoms



Differential transv. displacement: $x(+\Phi, t) - x(-\Phi, t) = 2x(t)$

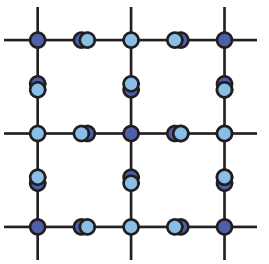


Difference image for $\pm\Phi$

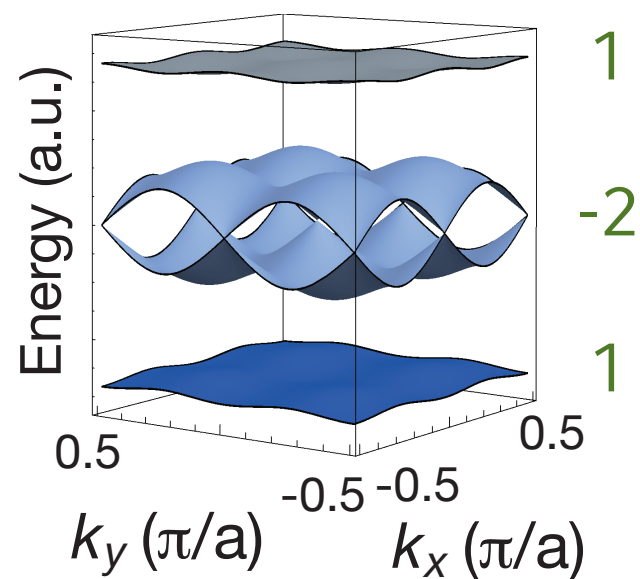


Observed deflection reduced by factor $\gamma \simeq 0.37$

Hall deflection with bosonic atoms



Higher-band populations:



- If only the lowest band is filled:

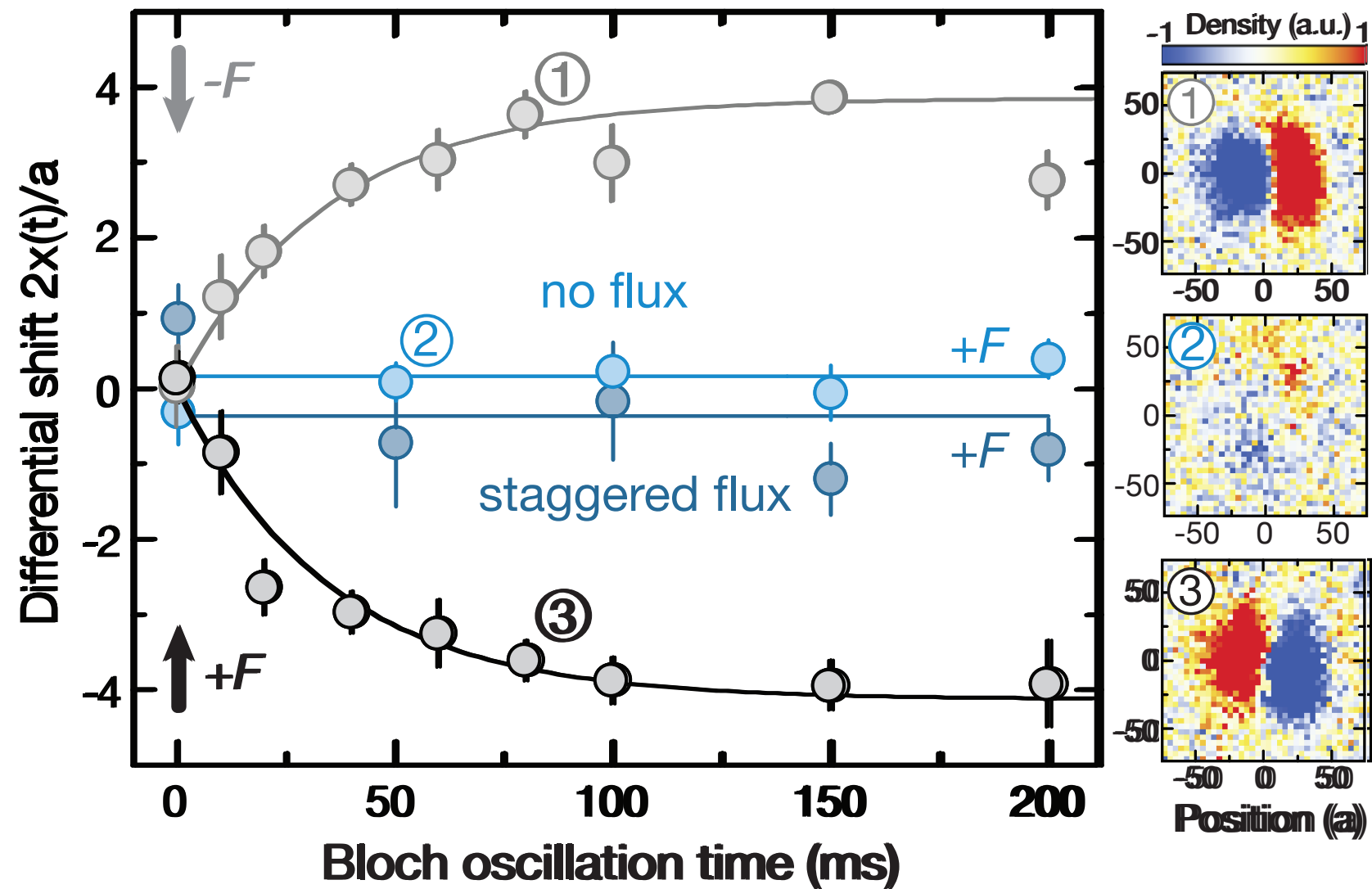
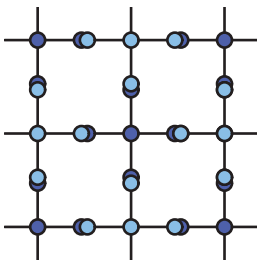
$$x(t) = -\frac{4a^2 F}{h} \nu_1 t$$

linear evolution prop. to Chern number

- Higher-bands: $x(t) \rightarrow \gamma x(t)$ with $\gamma = \eta_1 - \eta_2 + \eta_3$,
 η_i : rel. band populations

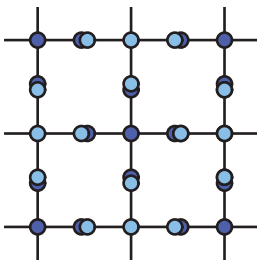
Derivation relies on particle-hole symmetry & $\sum_{\mu} \nu_{\mu} = 0$

Hall deflection with bosonic atoms

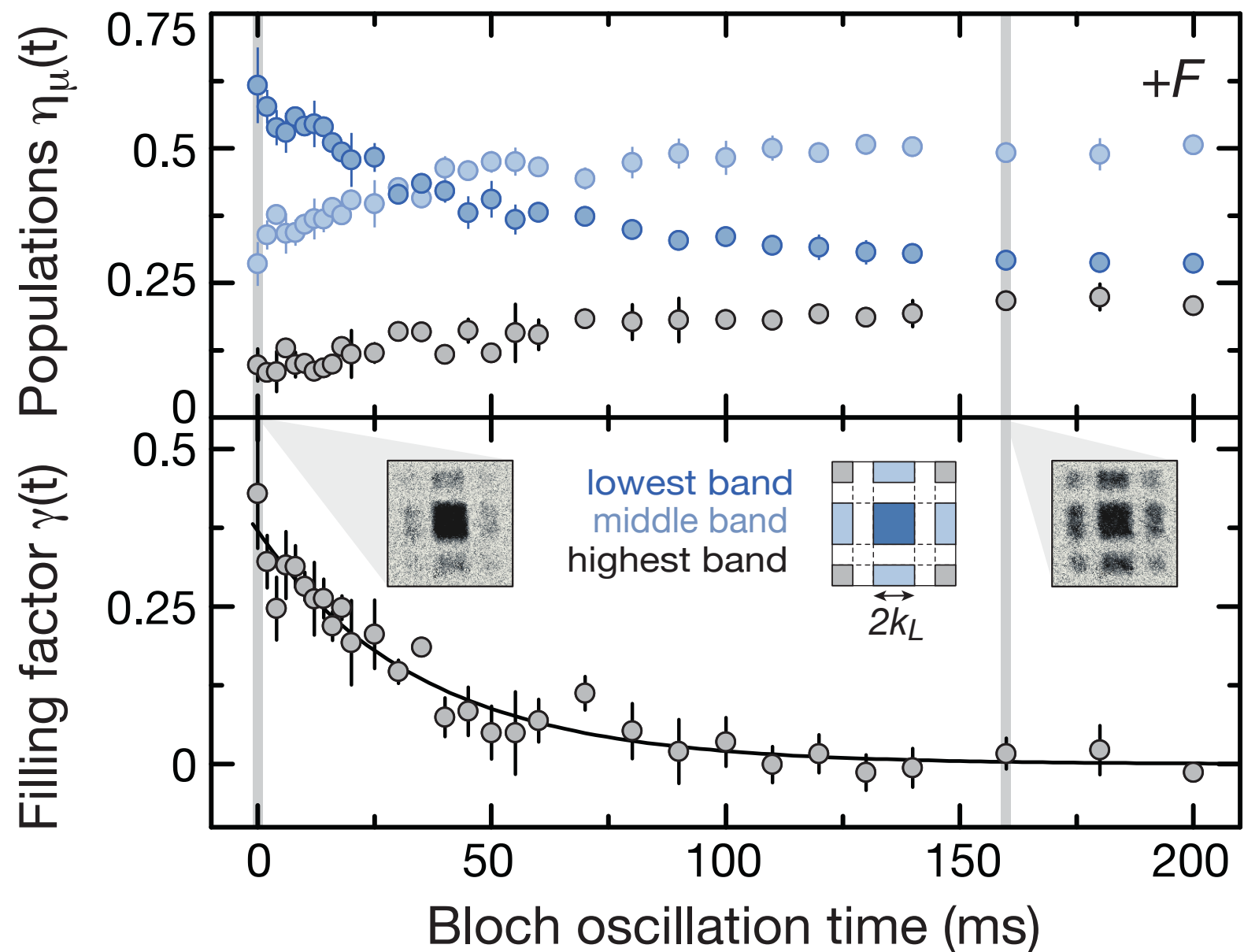


Saturation for longer evolution times

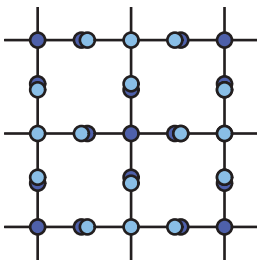
Hall deflection with bosonic atoms



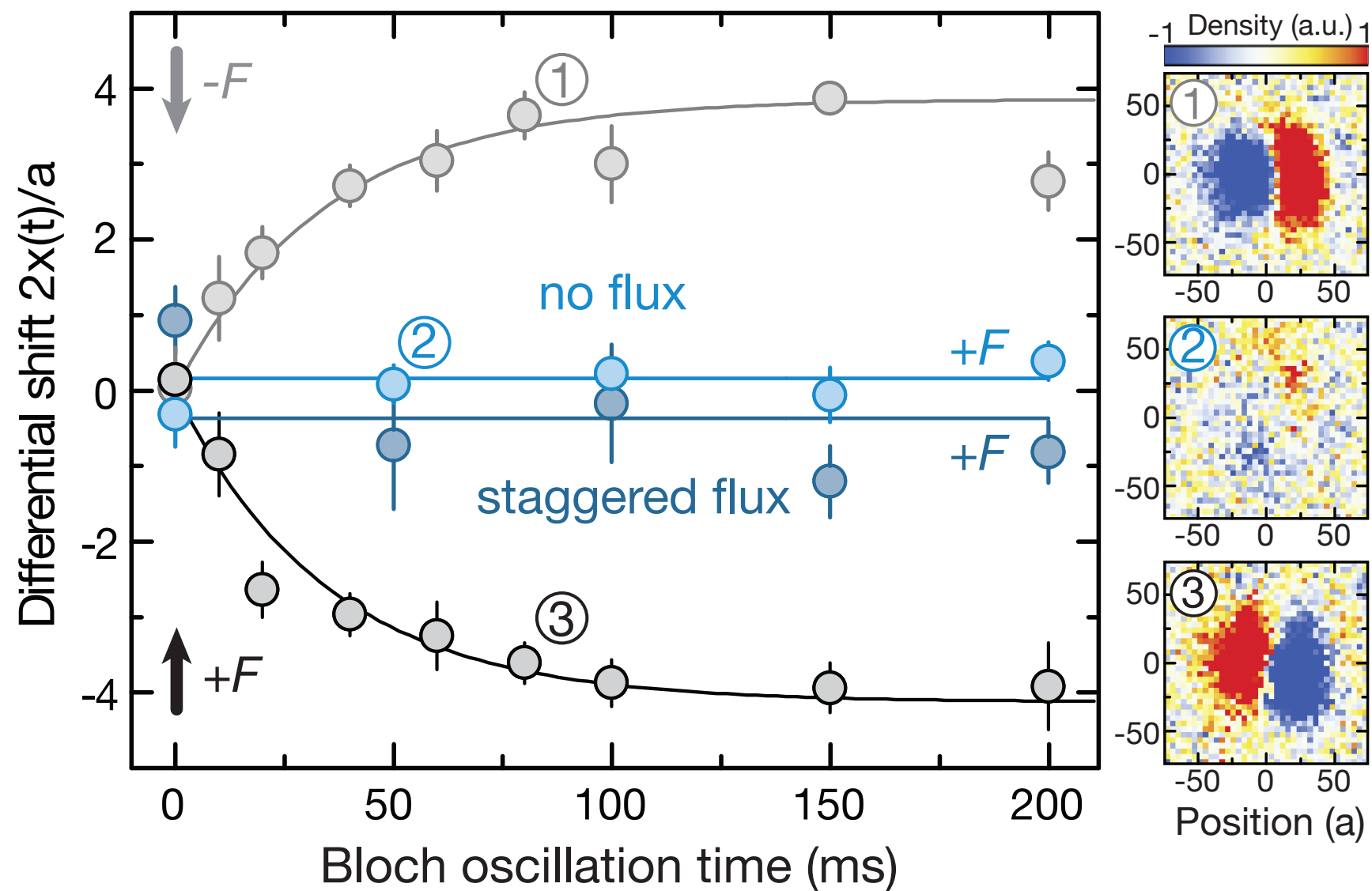
Measured band populations:



Hall deflection with bosonic atoms

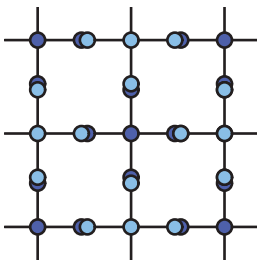


Combination of both:
$$x(t) = -\frac{4Fa^2}{h} \nu_1 \int_0^t \gamma(t') dt'$$



$\nu_{\text{exp}} = 0.99(5)$

Probing band topology

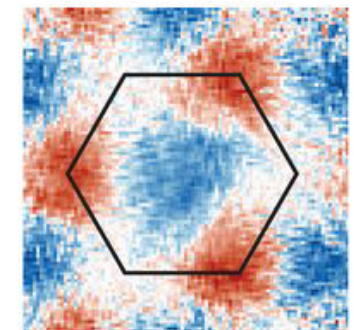
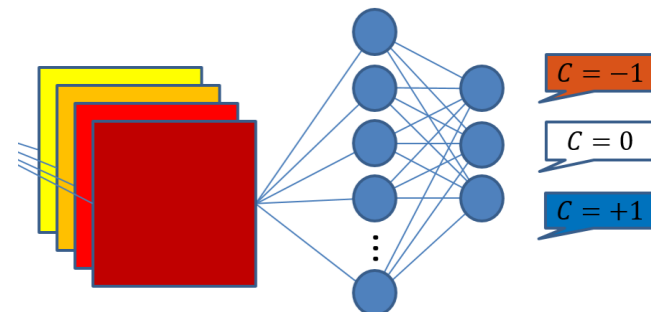
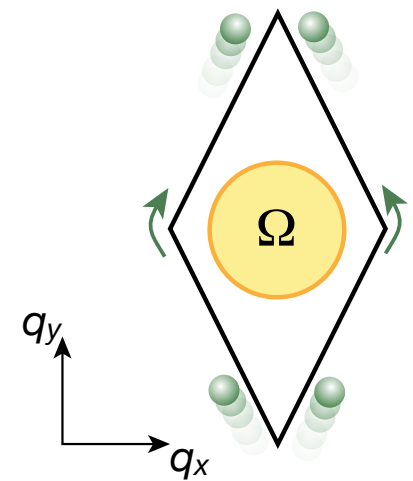


Chern number:

$$\mathcal{C}^\mu = \frac{1}{2\pi} \int_{\text{BZ}} \Omega^\mu d^2q$$

$|u_\mu(\mathbf{q})\rangle$: periodic Bloch function,
 μ : band index

- Interferometric measurements / transport:
M. Atala *et al.*, Nat. Phys. (2013); L. Duca *et al.*, Science (2015)
G. Jotzu *et al.*, Nature (2014); M. A. *et al.*, Nature Phys. (2015)
- State tomography:
N. Fläschner *et al.*, Science (2016); T. Li *et al.*, Science (2016)
- Linking number / circular dichroism / machine learning:
Tarnowski *et al.*, Nat. Comm. (2019);
L. Asteria *et al.*, Nat. Phys. (2019);
B. Rem *et al.*, Nat. Phys. (2019)



see also: experiments in synth. dimensions