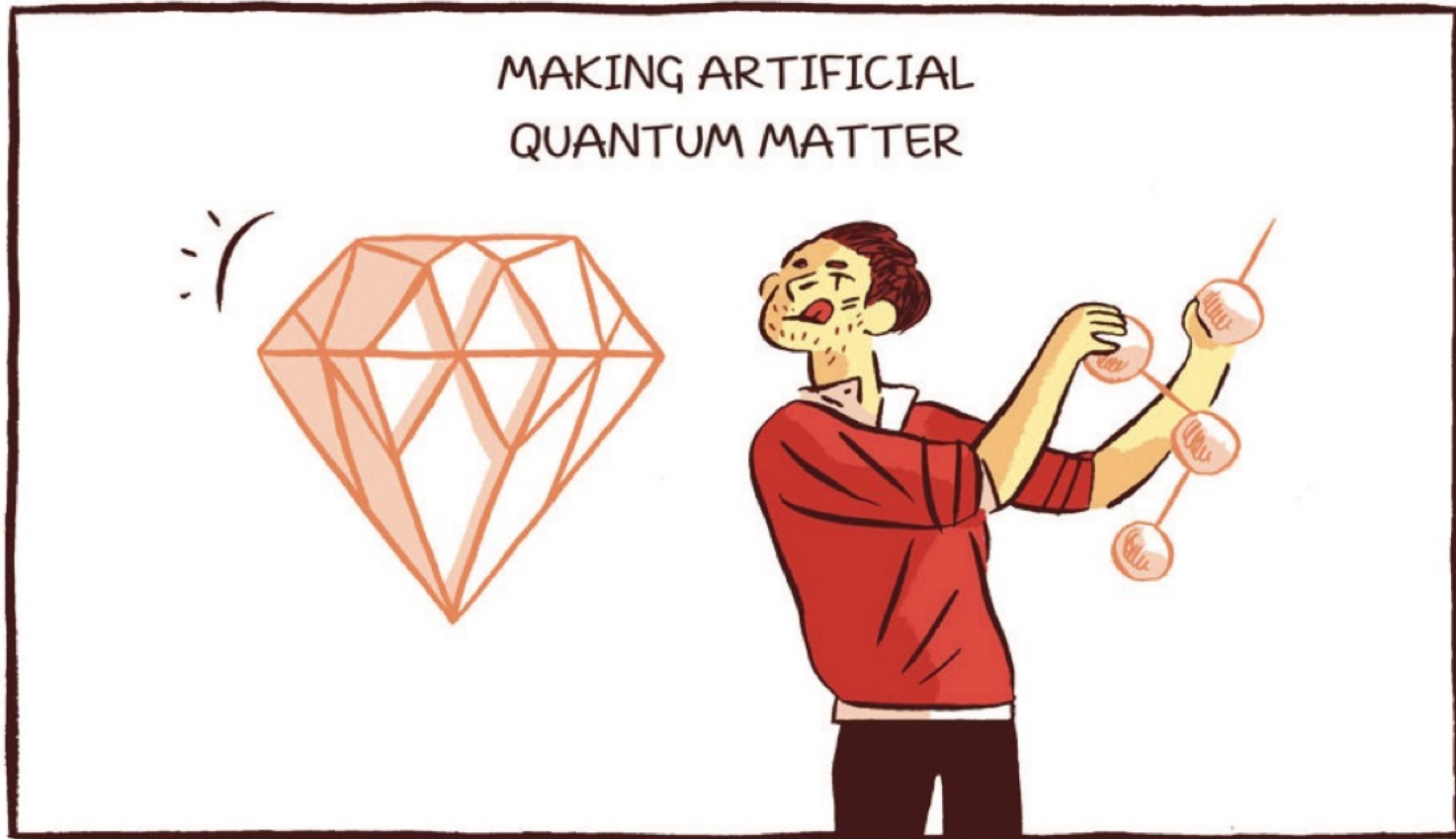


Rydberg atoms in optical tweezers: many-body physics and QIP

Antoine Browaeys

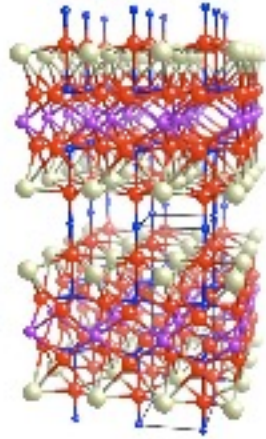
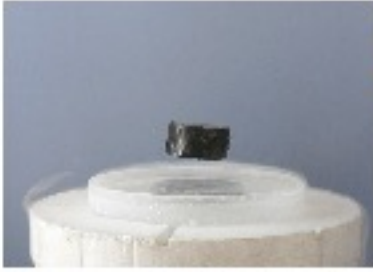
*Laboratoire Charles Fabry,
Institut d'Optique, CNRS, FRANCE*

Quantum state engineering: create controlled quantum states



Quantum state engineering: create controlled quantum states

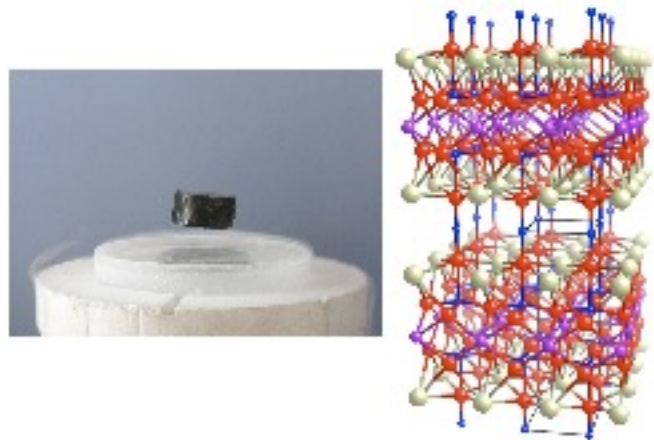
Many-body physics



Quantum simulation.

Quantum state engineering: create controlled quantum states

Many-body physics



Quantum metrology

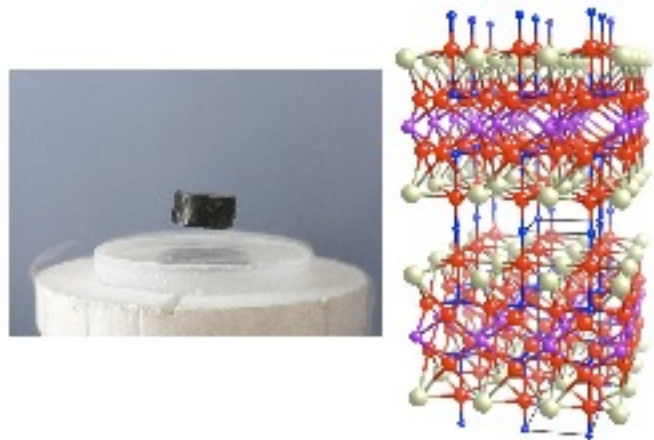


$\frac{\Delta\nu}{\nu_0} \propto \frac{1}{\sqrt{N}}$ **qu. states** $\frac{\Delta\nu}{\nu_0} \propto \frac{1}{N}$

$\omega_a \uparrow \begin{matrix} -e \\ -g \end{matrix}$ $\begin{matrix} v & \vdots & v \\ v & \vdots & v \end{matrix}$

Quantum state engineering: create controlled quantum states

Many-body physics

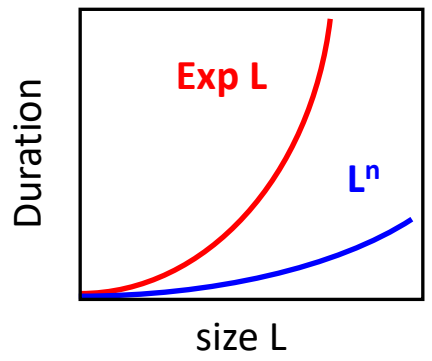


Quantum metrology

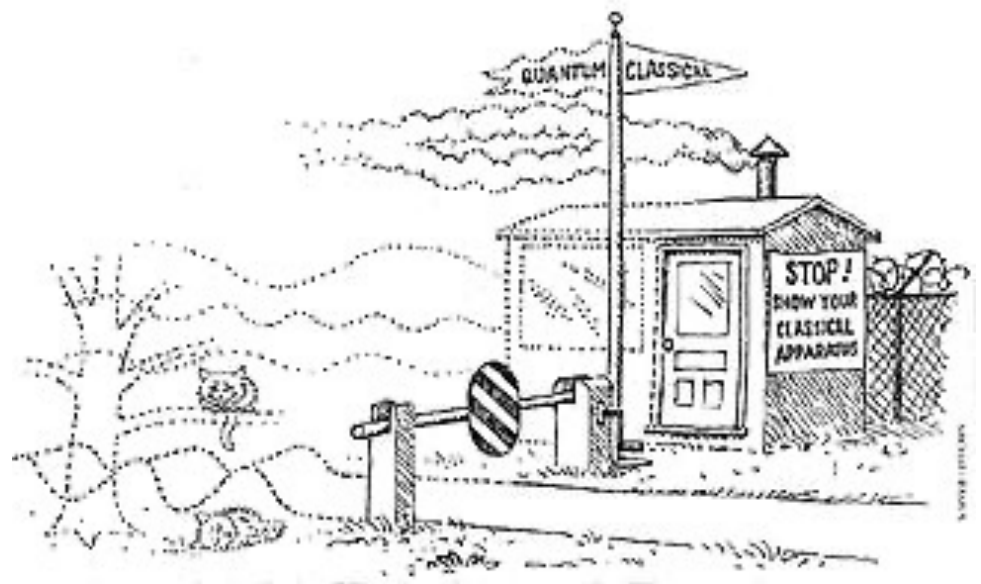


$\frac{\Delta\nu}{\nu_0} \propto \frac{1}{\sqrt{N}}$ **qu. states** $\rightarrow$ $\frac{\Delta\nu}{\nu_0} \propto \frac{1}{N}$

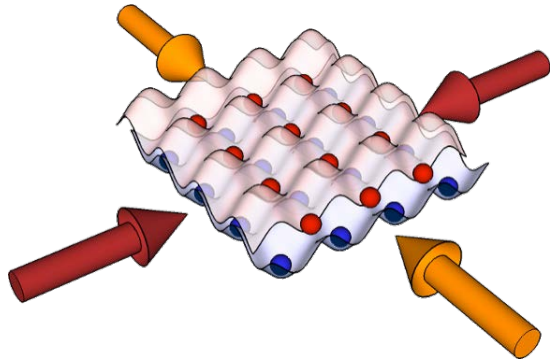
Quantum information



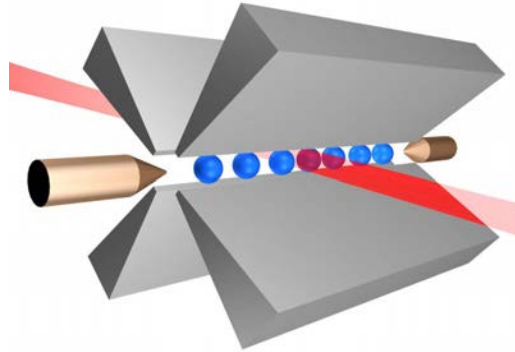
Transition quantum / classical



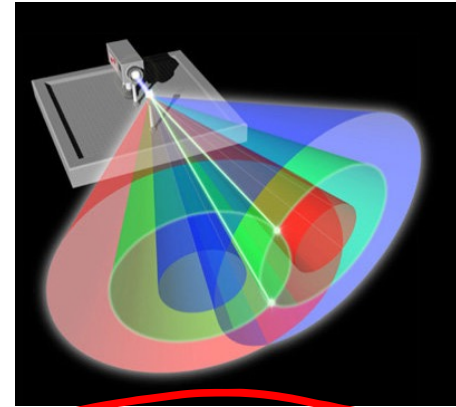
Quantum state engineering with individual systems



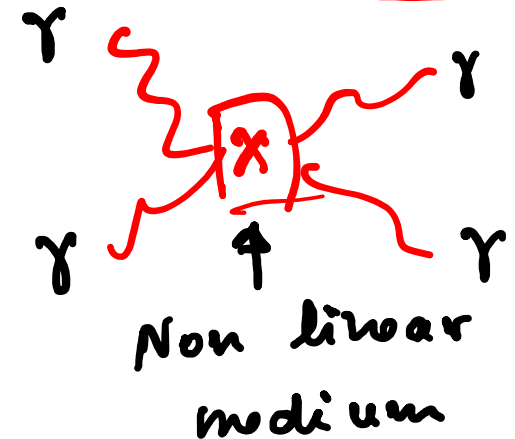
Cold atoms and molecules



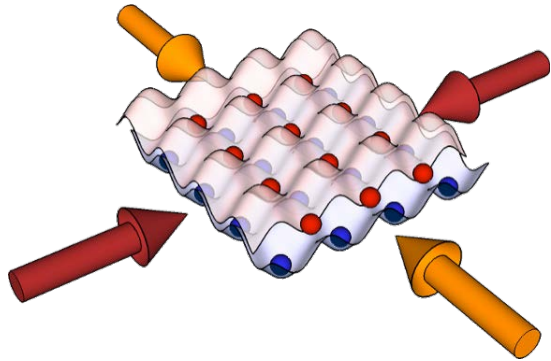
Trapped ions



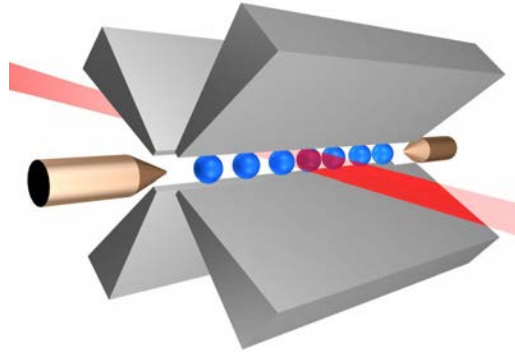
Photons in NL media



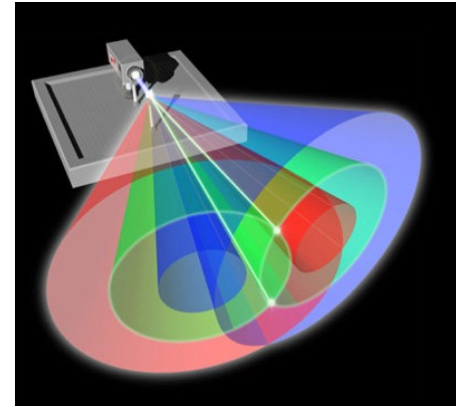
Quantum state engineering with individual systems



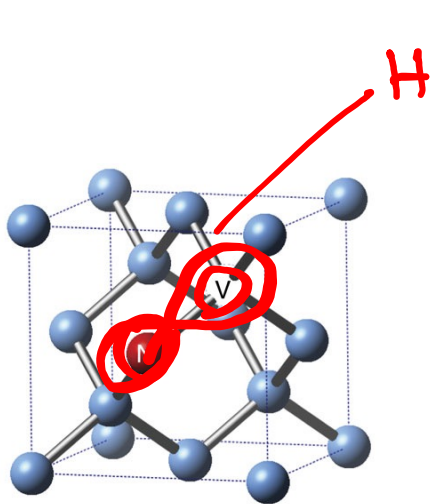
Cold atoms and molecules



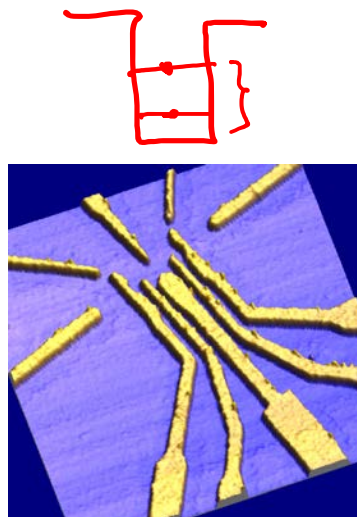
Trapped ions



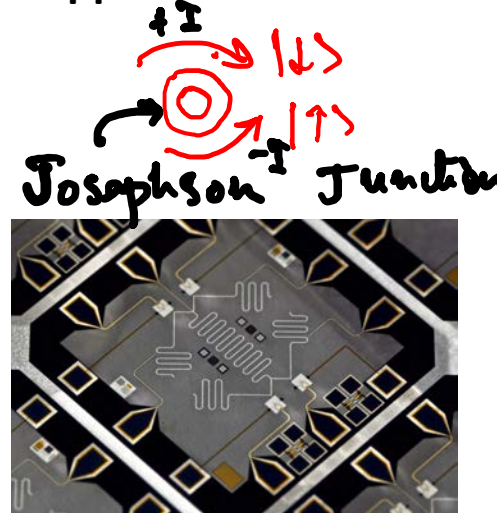
Photons in NL media



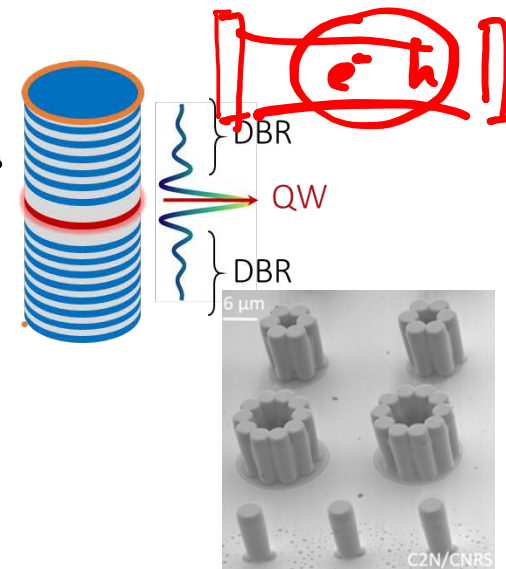
NV centers



Quantum dots

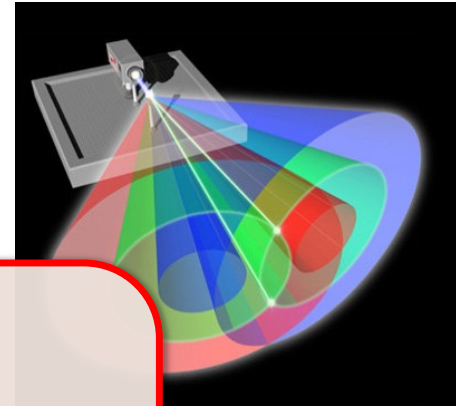
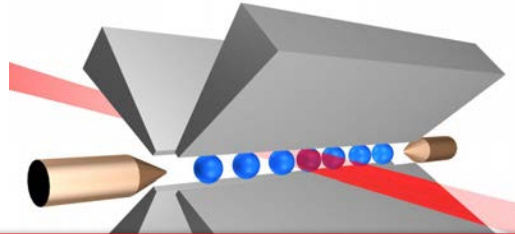
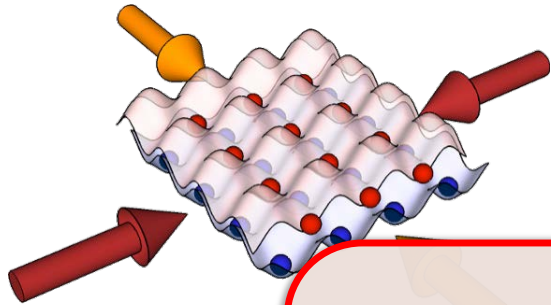


Superconducting qubits



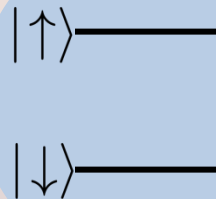
Polaritons in 1D cond.

Quantum state engineering with individual systems

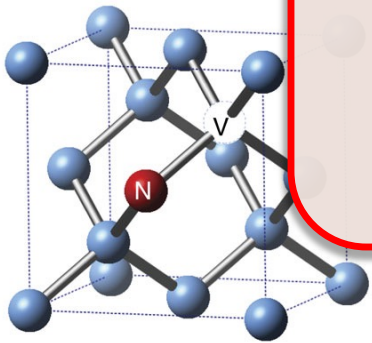


Cold atoms

Two-level systems to encode a spin:



Addressable + controlled interactions



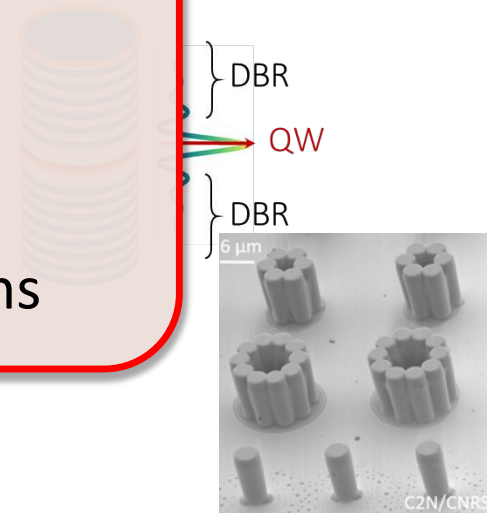
NV centers



Quantum dots

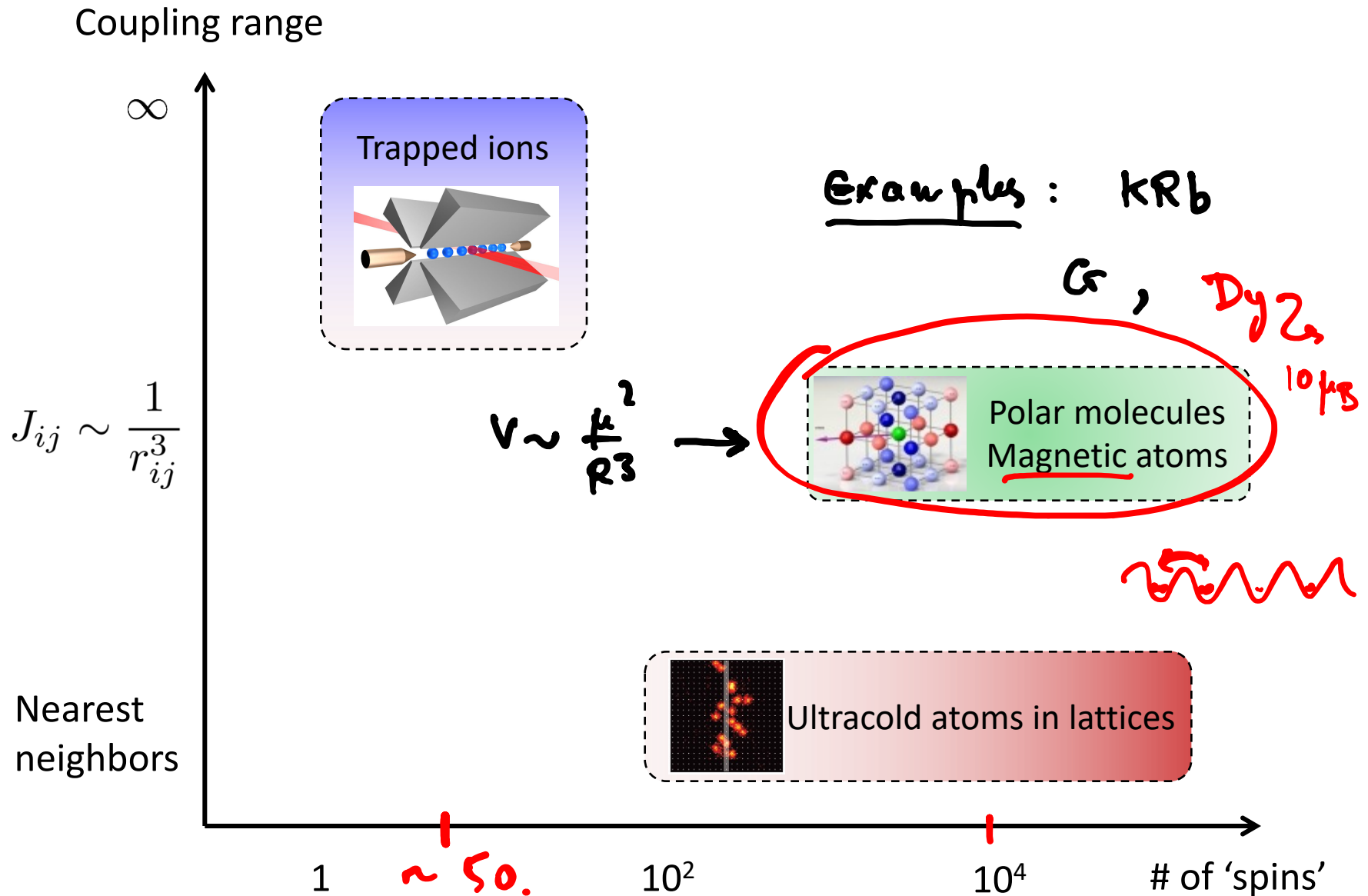


Superconducting qubits



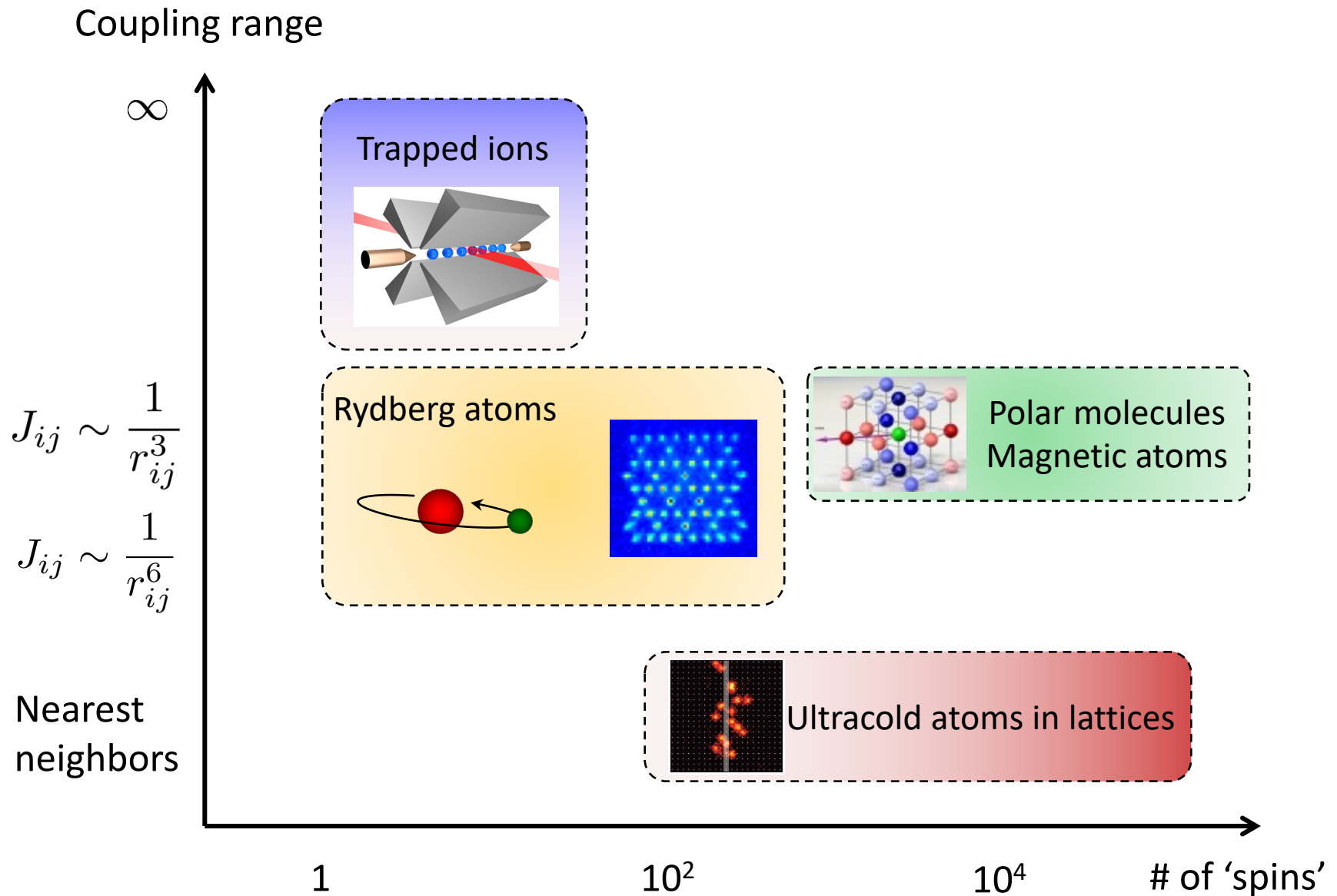
Polaritons in $\frac{1}{2}$ cond.

"Atomic" arrays (ignoring bulk gases...)



See e.g. Hazzard *et al.*, PRA **90**, 063622 (2014)

“Atomic” arrays (ignoring bulk gases...)



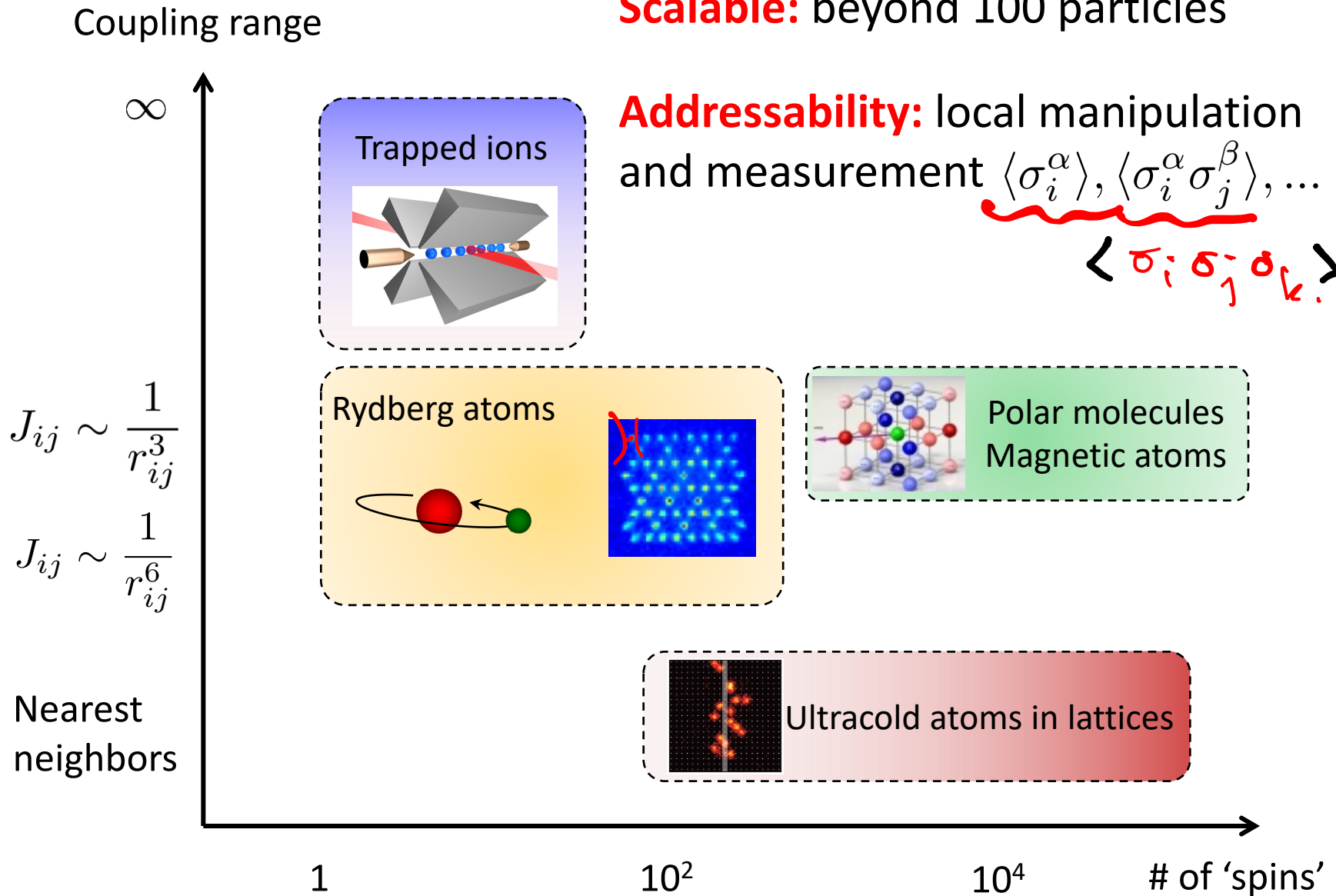
See e.g. Hazzard *et al.*, PRA **90**, 063622 (2014)

“Atomic” arrays (ignoring bulk gases...)

Scalable: beyond 100 particles

Addressability: local manipulation and measurement $\langle \sigma_i^\alpha \rangle, \langle \sigma_i^\alpha \sigma_j^\beta \rangle, \dots$

$$\langle \sigma_i \sigma_j \sigma_k \rangle$$

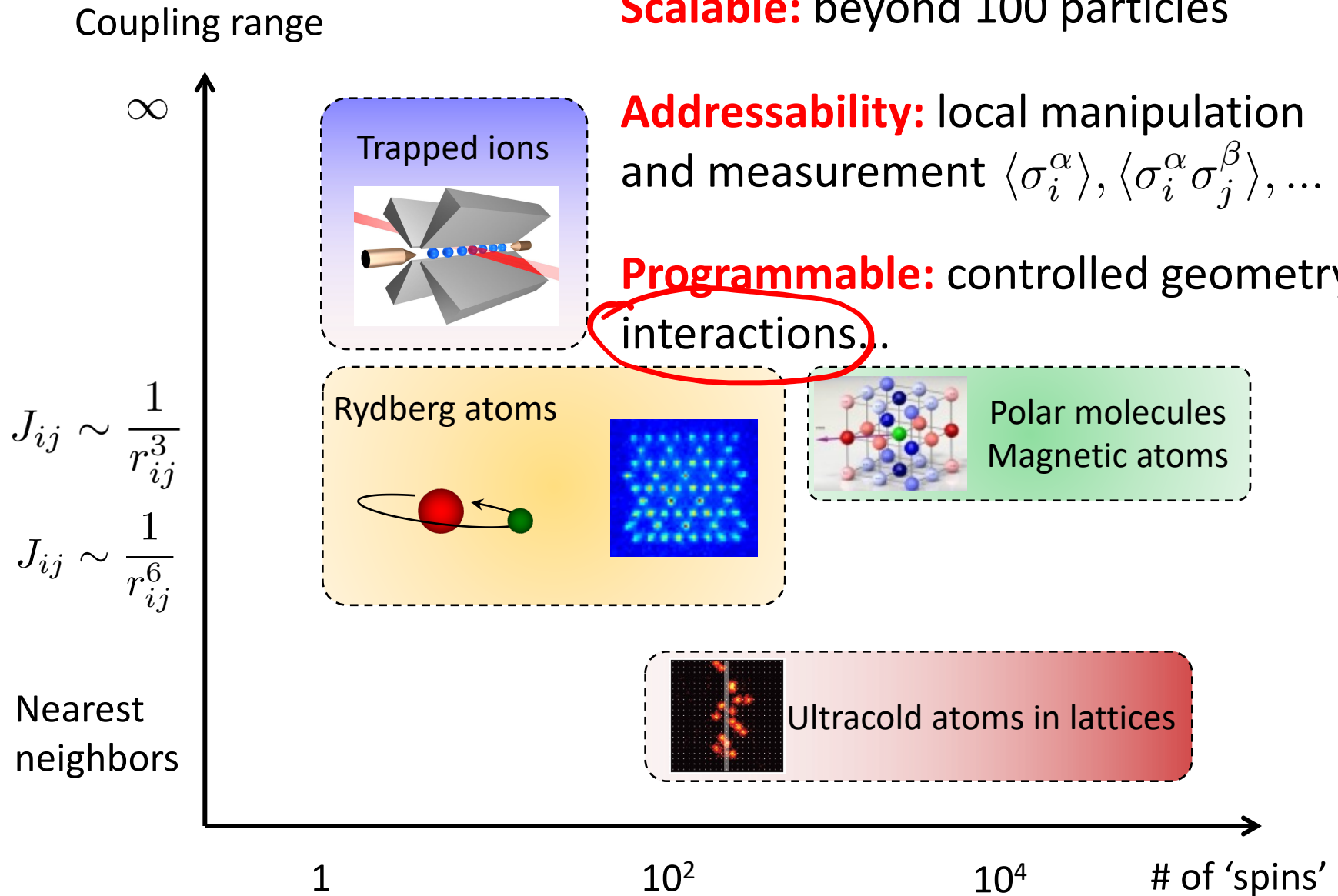


“Atomic” arrays (ignoring bulk gases...)

Scalable: beyond 100 particles

Addressability: local manipulation and measurement $\langle \sigma_i^\alpha \rangle, \langle \sigma_i^\alpha \sigma_j^\beta \rangle, \dots$

Programmable: controlled geometry, interactions..

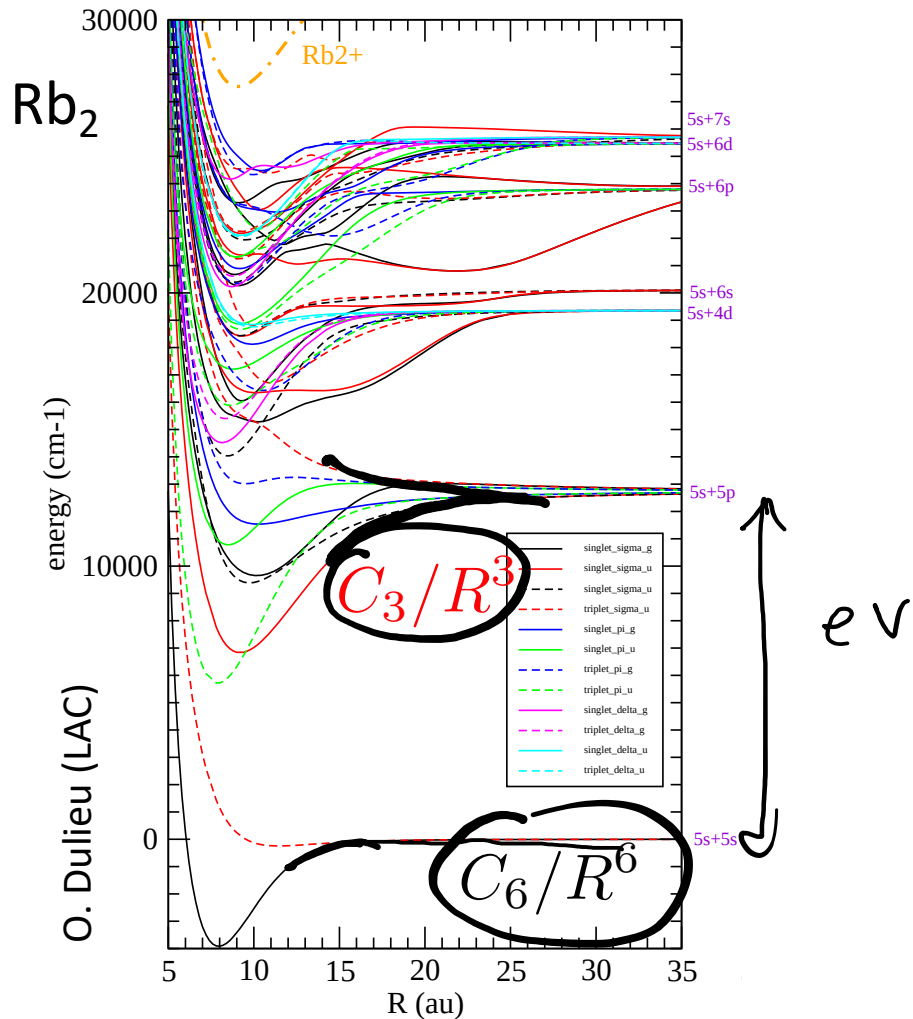


Outline

- Lecture 1: Dipolar interactions between atoms
- Lecture 2: Basics of Rydberg physics. Arrays of atoms. Rydberg blockade (2-body physics) & gates
- Lecture 3: Many-body physics with Rydberg atoms: spin models and transport

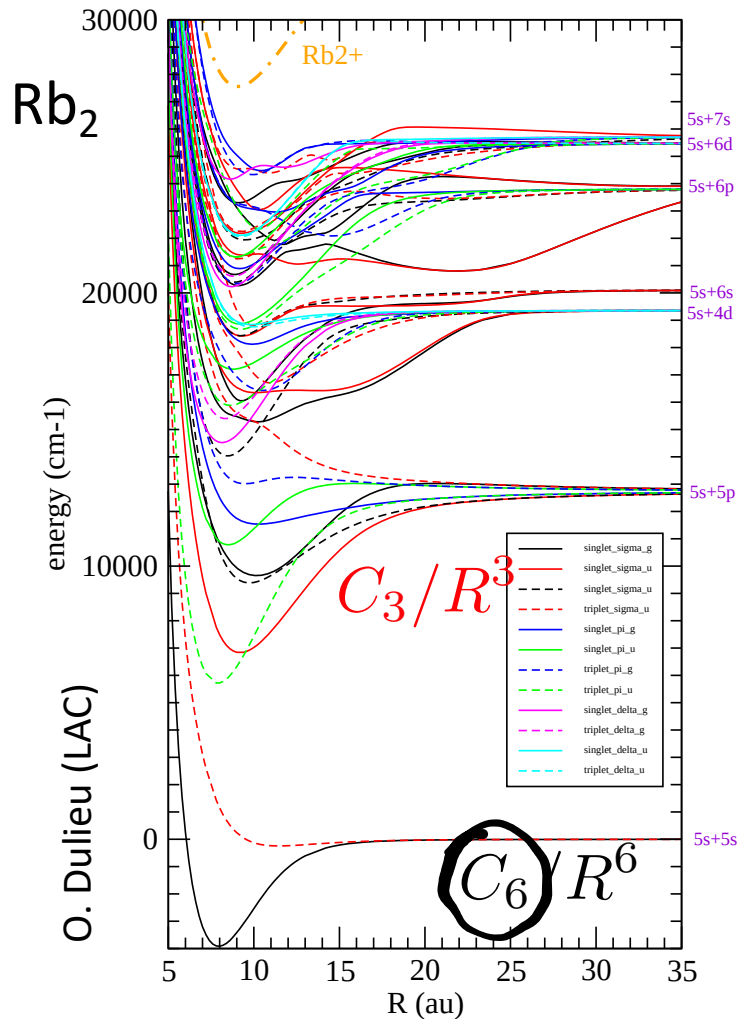
Interaction between atoms

Goal: understand *long-range* interactions between atoms



Interaction between atoms

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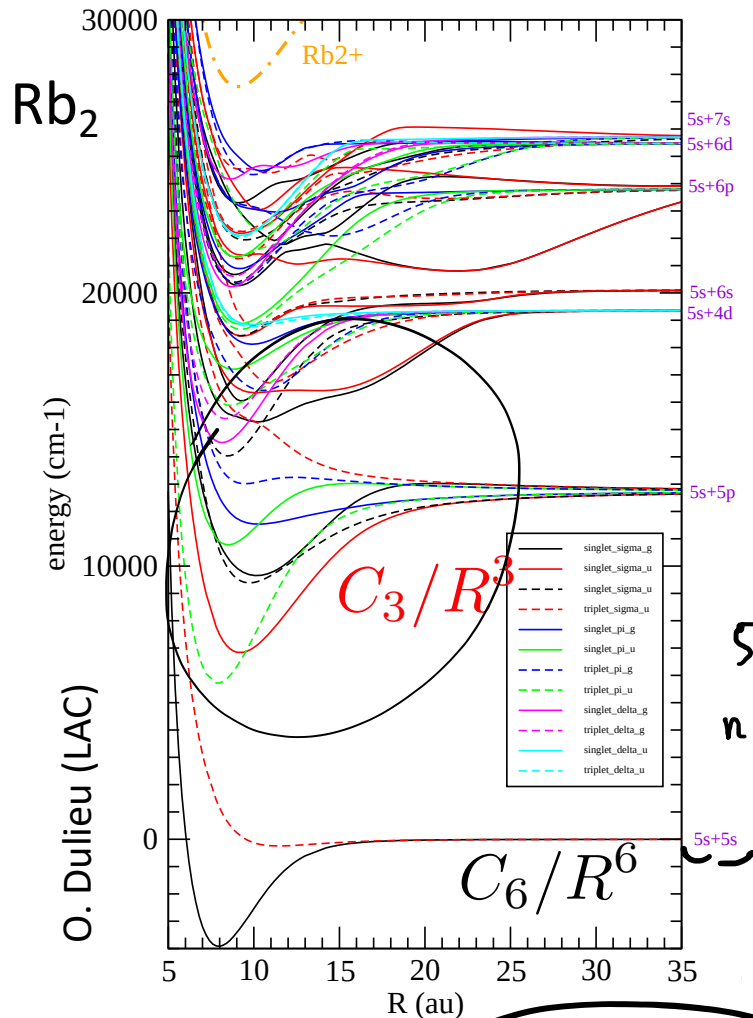
Van der Waals: C_6/R^6

système	C_6 [u.a.]	$\Gamma/2\pi$ [MHz]	λ [nm]
Li-Li	1389	5.87	671
Na-Na	1556	9.80	589
K-K	3897	6.04	767
Rb-Rb	4691	6.07	780
Cs-Cs	6870	5.22	852
Mg-Mg	627	80.9	235
Ca-Ca	2121	34.6	423
Sr-Sr	3103	32.0	461
Er-Er	1760	29.7	401
Dy-Dy	2275	32.2	421
Yb-Yb	1929	29	399

J. Dalibard, Collège de France, 2021

Interaction between atoms

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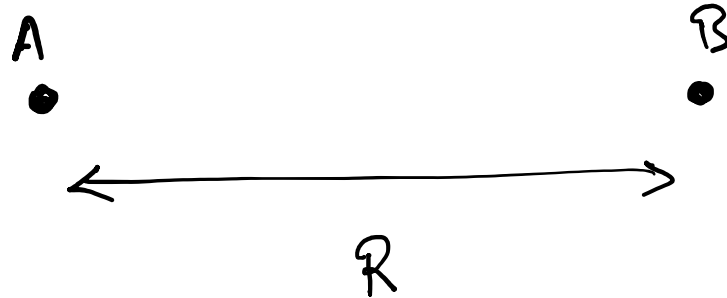
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J. Dalibard, Collège de France, 2021

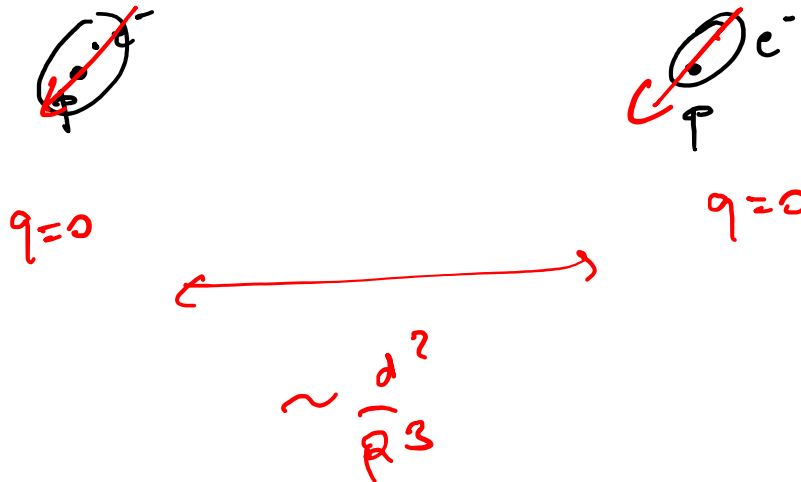
Useful for: scattering length (quantum gases), Rydberg physics...

Interaction between atoms



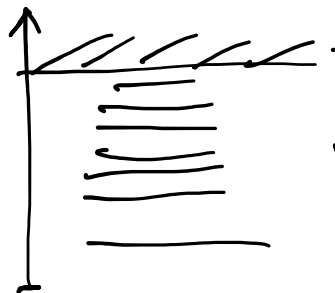
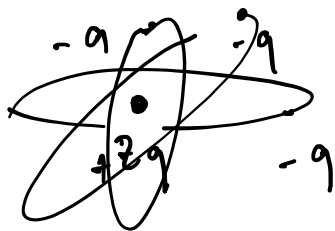
Point-like
neutral particle
 $\Rightarrow$ no interaction!

But:
atoms
are not
point-like!



$\Rightarrow$ dipole

Atomic dipoles (1)



$\psi_{nlm}(\vec{r})$

$$\vec{D} = \sum_{i=1}^Z q (\vec{r}_i - \vec{r})$$

(1e⁻ $\vec{d} = q \vec{r}$)

Quantum :

$\hat{D}$

Matrix element of $\hat{D}$:

$\hat{D} \rightarrow -\hat{D}$

$|nlm\rangle$

$\langle nlm | \hat{D} | nlm \rangle = 0$

$\int \psi_{nlm}^* q \vec{r} \psi_{nlm}(\vec{r}) d^3r$

$l=0$
 $l=1$

$\int |\psi_{nlm}|^2 q \vec{r} = 0$

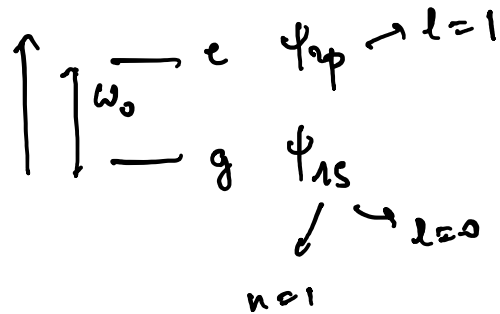


orthifical
atom $[H, L^2] \neq 0$

$\langle nlm | \hat{D} | nl \pm 1 m \rangle \neq 0$

Atomic dipoles (2)

Average dipole:



$$|\psi(t)\rangle$$

$$\langle \psi(t) | \vec{D} | \psi(t) \rangle = \langle \vec{D} \rangle$$

$$\langle e | D_z | g \rangle = d_{eg} \quad (\triangle \quad \langle e | D | e \rangle = \langle g | D | g \rangle = 0)$$

$$|\psi(t)\rangle = \alpha |g\rangle + \beta e^{-i\omega_0 t} |e\rangle$$

$$\alpha^2 + \beta^2 = 1$$

$$\Rightarrow \langle \vec{D} \rangle = 2\alpha\beta d_{eg} \cos\omega_0 t$$



Do not

mix

①

$$\langle n l m | D | n' l' m' \rangle$$

(Matrix element)

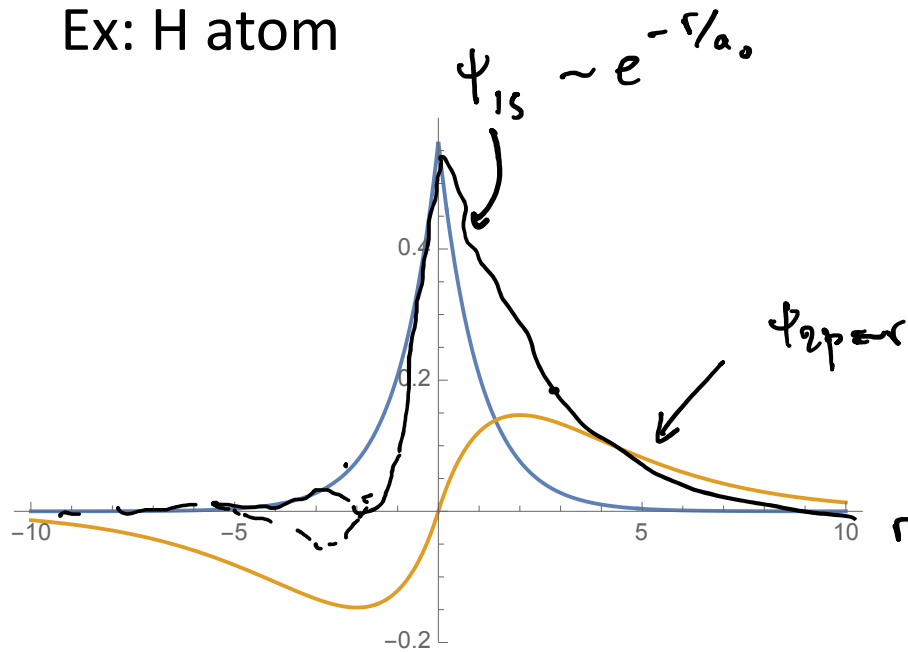
②

Average dipole

$$\langle \psi(t) | D | \psi(t) \rangle$$

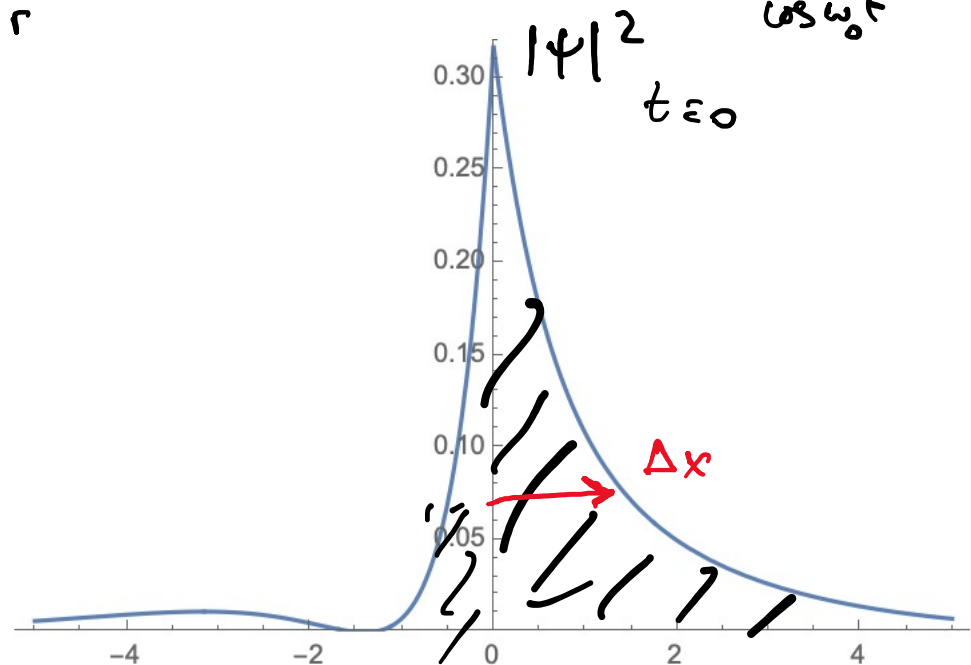
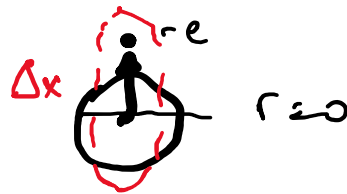
Atomic dipoles (3)

Ex: H atom



$$\psi(\vec{r}, t) = \psi_{1s} + \psi_{2p} e^{-i\omega_0 t}$$

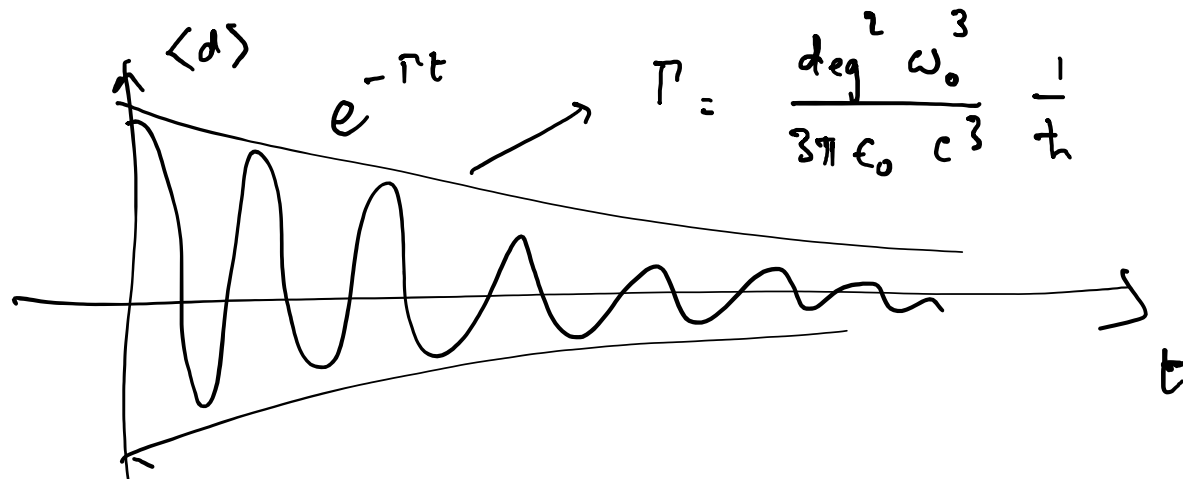
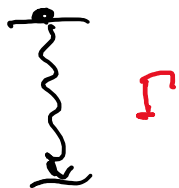
$$|\psi(\vec{r}, t)|^2 = |\psi_{1s}|^2 + |\psi_{2p}|^2 + 2\psi_{1s}\psi_{2p} \cos \omega_0 t$$



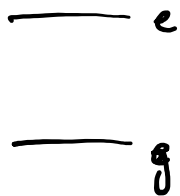
Atomic dipoles (4)

Dipole radiates a field!

$\uparrow$ $d(t), \omega_0$



①



μW $\omega_0 \sim GHz$
 Γ

" $\Gamma \rightarrow 0$ " $\omega_0 \downarrow$

Many oscillations

②

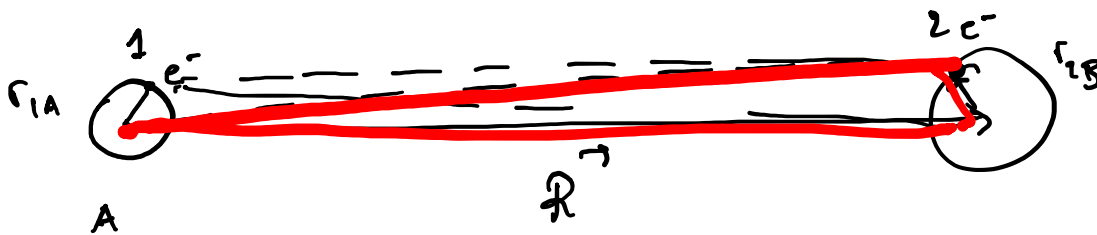


$\lambda \sim \mu m.$ $\omega_0 \sim 400 THz$



$$Q = \frac{\omega_0}{\Gamma} \rightarrow \frac{4 \times 10^{12}}{10 \text{ THz}} \sim 10^6.$$

Dipolar interaction between 2 atoms



$$e^2 = \frac{q^2}{4\pi\epsilon_0}$$

$$H = \frac{p_A^2}{2m} - \frac{e^2}{r_{1A}} + \frac{p_B^2}{2m} - \frac{e^2}{r_{2B}} + \underbrace{\frac{e^2}{R} + \frac{e^2}{r_{12}} - \frac{e^2}{r_{1B}} - \frac{e^2}{r_{2A}}}_{\text{dipole-dipole interaction.}}$$

$$r_{12} = |\vec{R} - \vec{r}_1 + \vec{r}_2|$$

$$\frac{1}{r_{12}} \sim \frac{1}{R} \left(1 - \frac{\vec{r}_1 \cdot \vec{R}}{R^2} - \frac{r_2^2}{R^2} + \frac{3}{2} \left(\frac{\vec{r}_2 \cdot \vec{R}}{R^2} \right)^2 + \dots \right)$$

dipole-dipole interaction.

$$r_{1B} = |\vec{R} + \vec{r}_1|$$

H_I

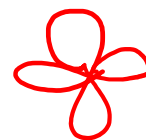
Size of atom $\sim a_0 \ll R$

$$\Rightarrow \dots \Rightarrow H_I = \frac{q^2}{4\pi\epsilon_0} \frac{\vec{r}_1 \cdot \vec{r}_2 - 3(\vec{r}_1 \cdot \vec{u})(\vec{r}_2 \cdot \vec{u})}{R^3}$$

$$\vec{u} = \frac{\vec{R}}{R}$$

$$\vec{D}_i = q \vec{r}_i$$

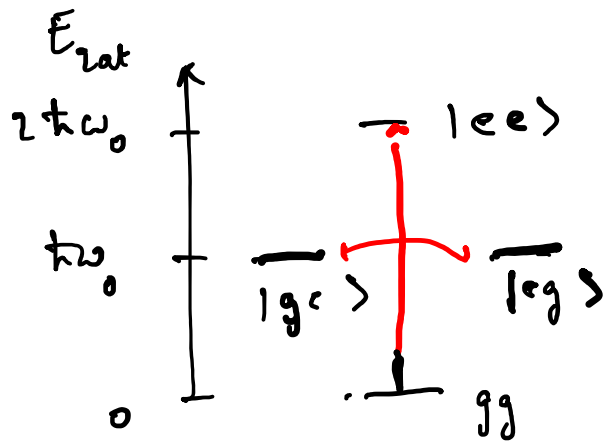
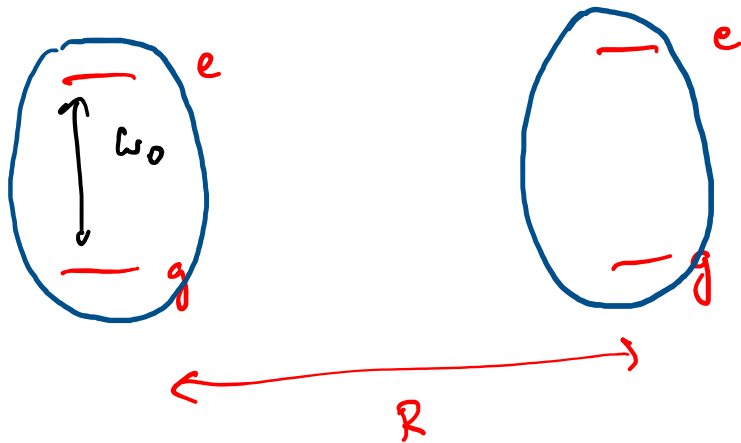
$$H_I = \frac{\vec{D}_1 \cdot \vec{D}_2 - 3(\vec{D}_1 \cdot \vec{u})(\vec{D}_2 \cdot \vec{u})}{4\pi\epsilon_0 R^3}$$



$$\Sigma = \Sigma_1 \otimes \Sigma_2 \rightarrow$$

$$|n_1, l_1, m_1\rangle \otimes |n_2, l_2, m_2\rangle$$

A toy model for two interacting atoms



$$\langle gg | D_1 D_2 | ee \rangle = \frac{d_{eg}^2}{R^3}$$

$$\langle eg | D_1 D_2 | eg \rangle = \frac{d_{eg}^2}{R^3}$$

$\vec{D}_1 \cdot \vec{u} = 0$

$d_{eg} \langle e | D_z | g \rangle$

$\vec{D}_1 \quad \vec{R} \quad \vec{D}_2$

$H = \frac{\vec{D}_1 \cdot \vec{D}_2}{R^3}$

$H =$

$\begin{pmatrix} 0 & \frac{d_{eg}^2}{R^3} \\ \frac{d_{eg}^2}{R^3} & 2\hbar\omega_0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$
$\begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$	$\begin{pmatrix} \hbar\omega_0 & \frac{d_{eg}^2}{R^3} \\ \frac{d_{eg}^2}{R^3} & \hbar\omega_0 \end{pmatrix}$

$\begin{matrix} gg \\ ee \\ eg \\ ge \end{matrix}$

2 Blocks !

The off-resonant block: van der Waals

$$H^{(1)} = \begin{pmatrix} 0 & d_{eg}^2/r^3 \\ \frac{d_{eg}^2}{r^3} & 2\hbar\omega_0 \end{pmatrix}$$

Recall:

$$H = \begin{pmatrix} 0 & V \\ V & \Delta \end{pmatrix} \text{ when } \boxed{\Delta \gg V}$$

$$\Delta - \frac{V^2}{\Delta} \rightarrow |\tilde{e}\rangle = |e\rangle - \frac{V}{\Delta} |g\rangle$$

$$0 - \frac{V^2}{\Delta} \rightarrow |\tilde{g}\rangle = |g\rangle + \frac{V}{\Delta} |e\rangle$$

$$2\hbar\omega_0 + \left(\frac{d_{eg}^2}{r^3}\right)^2 \frac{1}{2\hbar\omega_0}$$

$$0 - \left(\frac{d_{eg}^2}{r^3}\right)^2 \frac{1}{2\hbar\omega_0} = \Delta E^{(1)}$$



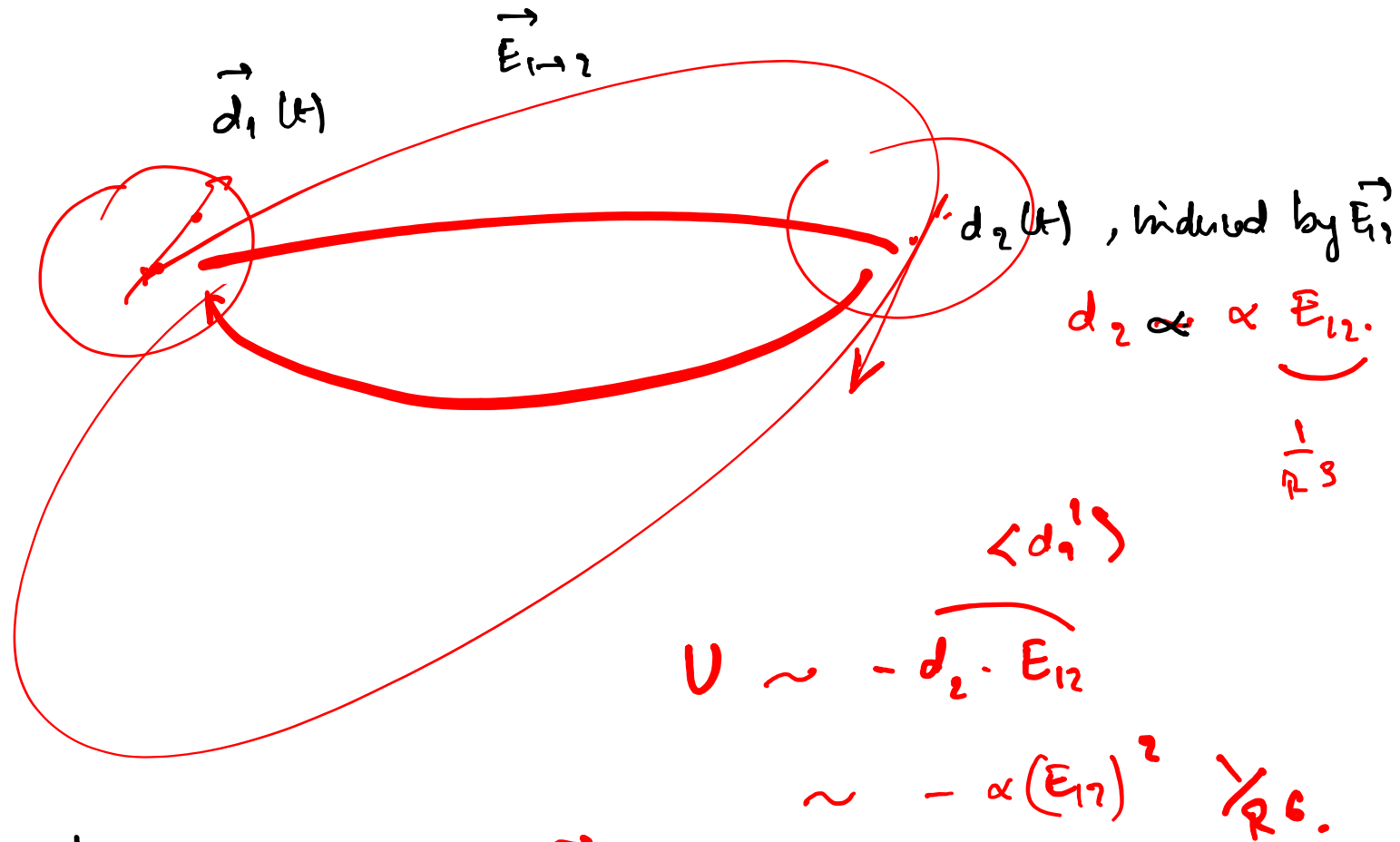
$$\langle gg | \frac{d_1 d_2}{r^3} | gg \rangle$$

$$= \Delta E^{(1)} = 0$$

$$\hookrightarrow C_6/r^6.$$

Van der Waals

van der Waals from a "classical" viewpoint



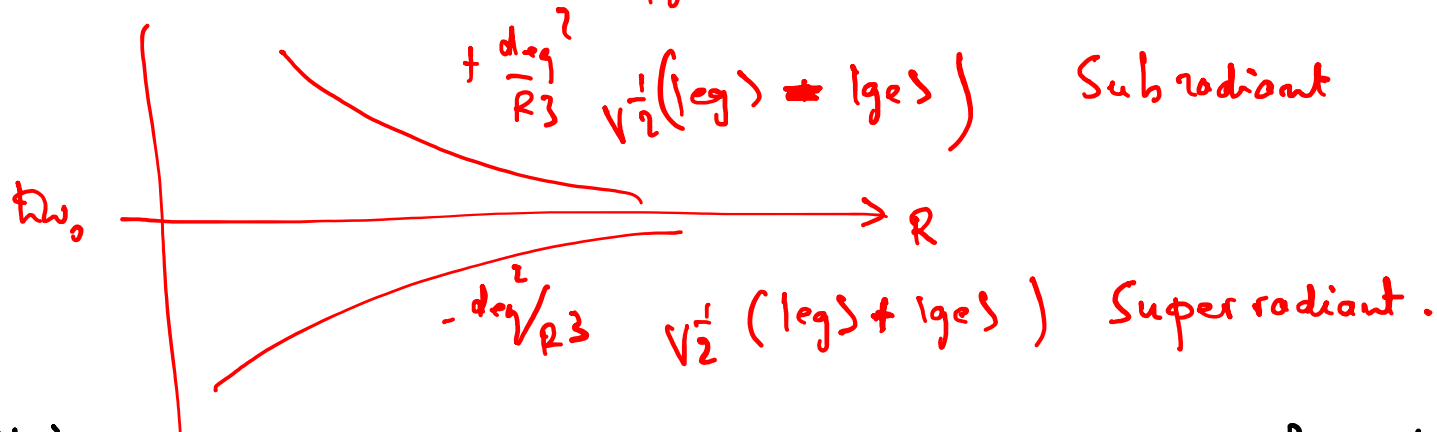
Van der Waals;
correlations between
dipoles!

$$\begin{aligned}\langle \vec{d}_1 \rangle &= 0 \\ \langle d_1^2 \rangle &= \overline{d_1^2}\end{aligned}$$

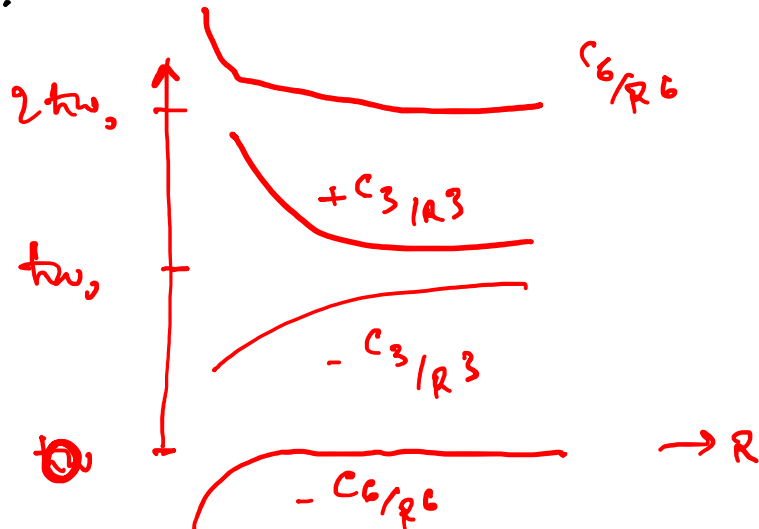
$$\langle \vec{d}_1 \vec{d}_2 \rangle$$

The resonant block: resonant dipole interaction

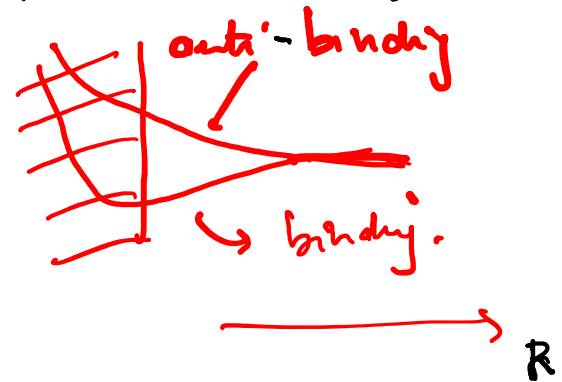
$$H^{(2)} = \begin{pmatrix} \hbar\omega_0 & \frac{d_{eg}^2}{R^3} \\ \frac{d_{eg}^2}{R^3} & -\hbar\omega_0 \end{pmatrix} \begin{matrix} |eg\rangle \\ |ge\rangle \end{matrix} \rightarrow \hbar\omega_0 \pm \frac{d_{eg}^2}{R^3}$$



Recap:



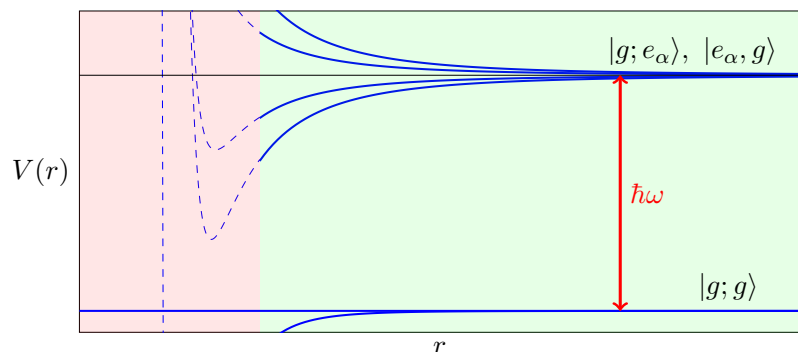
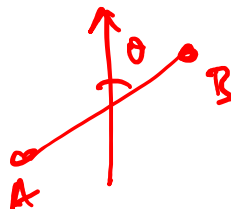
Note: in chemistry:



Molecular potentials in real life: example of Rb_2

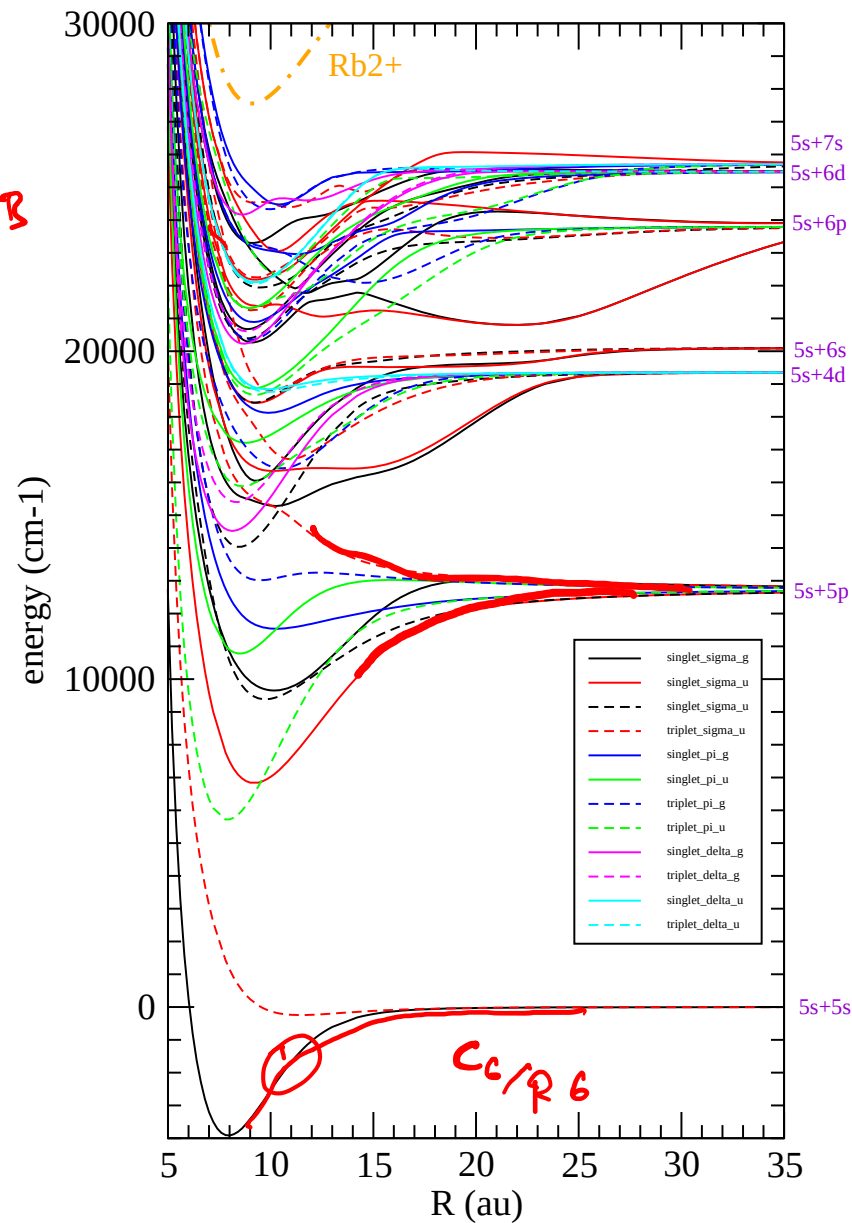
keeping $d_1, d_2 \dots$

$$C_3 \propto 1 - 3 \cos^2 \theta$$

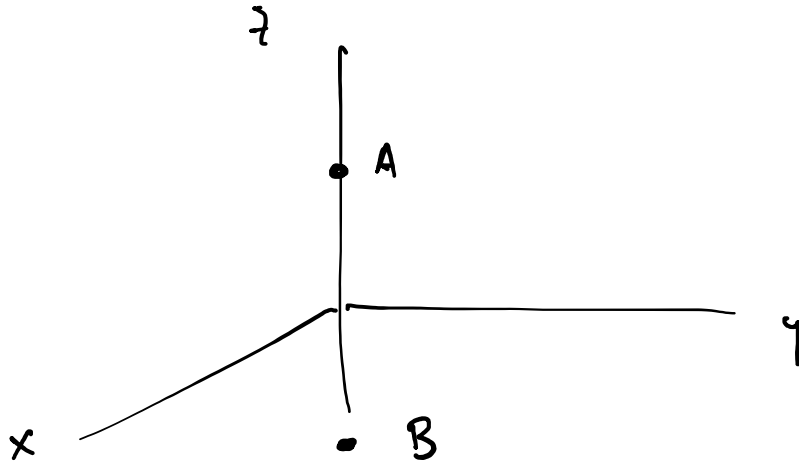


e
 $5p$ —
 $5s$ —
 g

e
 $5p$
 $5s$
 g



Ground state of "real" atoms

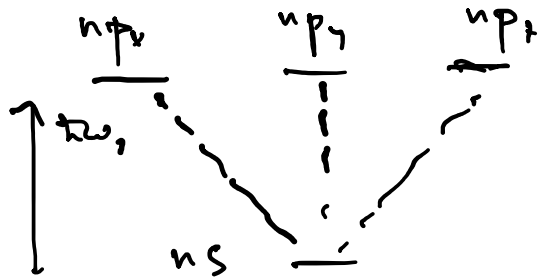


$$H = \frac{1}{4\pi\epsilon_0 r^3} \left[\hat{d}_{Ax} \hat{d}_{Bx} + \hat{d}_{Ay} \hat{d}_{By} - 2 \hat{d}_{Az} \hat{d}_{Bz} \right]$$

$$|ns, ns\rangle$$

$$np_x \quad np_y \quad np_z$$

$$\begin{aligned} p &\rightarrow l=0 \\ x &\rightarrow m=+1 \\ &0 \\ &-1 \end{aligned}$$



$$\Delta E_{ns, ns}^{(1)} = - \left(\frac{1}{4\pi\epsilon_0 r^3} \right)^2 \frac{1}{2\hbar\omega_0}$$

$$AS: \begin{array}{c} np \\ \downarrow \\ ns \end{array} \quad \hbar\Gamma = \frac{|d_{ps}|^2 \hbar^3}{3\pi\epsilon_0}, \quad k = \frac{\omega_0}{c}$$

$$\left[|d_{pxs}|^4 + |d_{py s}|^4 + 4 |d_{pz s}|^4 \right]$$

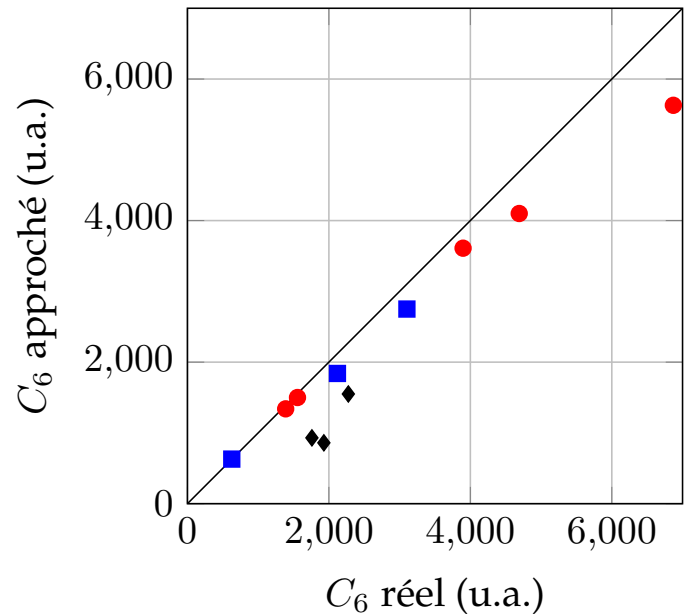


$$\Delta E_{ns, ns} = -\frac{27}{16} \frac{\hbar^3}{(\hbar r)^6} \frac{1}{\omega_0}$$

Ground state of “real” atoms

$$C_6 = -\frac{27}{16} \frac{\hbar \Gamma}{k^6} \frac{\Gamma}{\omega_0}$$

système	C_6 [u.a.]	$\Gamma/2\pi$ [MHz]	λ [nm]	C_6 approché [u.a.]
Li-Li	1389	5.87	671	1340
Na-Na	1556	9.80	589	1500
K-K	3897	6.04	767	3610
Rb-Rb	4691	6.07	780	4100
Cs-Cs	6870	5.22	852	5629
Mg-Mg	627	80.9	235	630
Ca-Ca	2121	34.6	423	1840
Sr-Sr	3103	32.0	461	2750
Er-Er	1760	29.7	401	930
Dy-Dy	2275	32.2	421	1550
Yb-Yb	1929	29	399	860



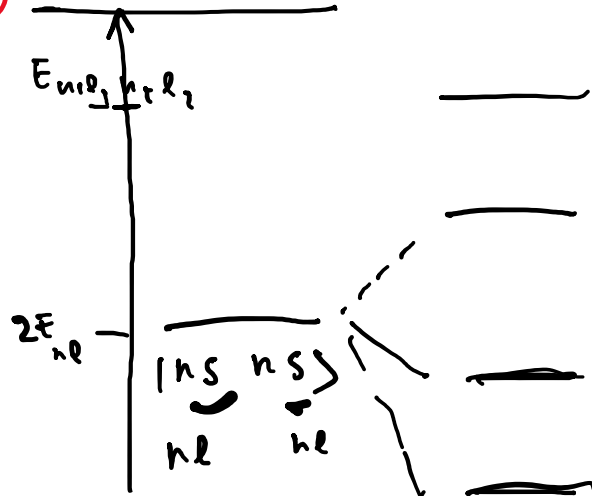
ee →

gg ⊗ $\left(\frac{d^2}{R^3}\right)^2 \frac{1}{2\hbar\omega_0}$

Excited states of "real" atoms: the Rydberg case

①

Van der Waals



$n_1 l_1, n_2 l_2$ $l_1 = \pm 1$

$$\langle n_1 l_1, n_1 l_1 | H_{dd} | n_l n_l \rangle$$

$C_6 > 0$

OR

$C_6 < 0$

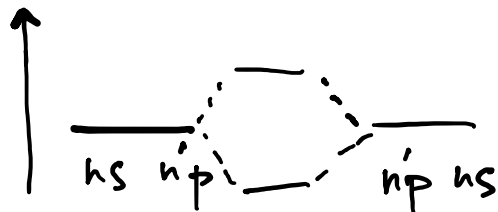
$\frac{1}{R^3}$

$$\Rightarrow \Delta E_{n_s n_s}^{(2)} = \sum_{\substack{n_1 l_1 \\ n_2 l_2}}$$

$$\frac{|\langle n_1 l_1, n_2 l_2 | \widehat{H_{dd}} | n_l n_l \rangle|^2}{2 E_{n_l} - E_{n_1 l_1} - E_{n_2 l_2}} \propto \frac{C_6}{R^6}$$

② Resonant

E_{res}



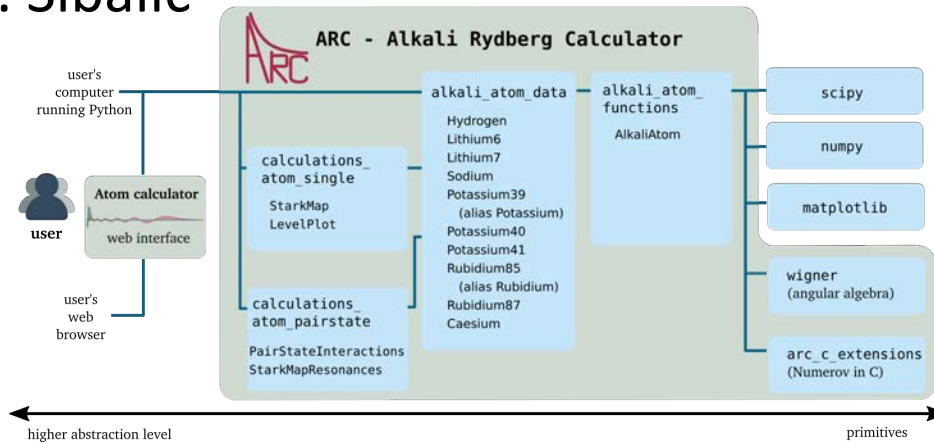
$\frac{C_3}{R^3}$

2

Resonant dip-dip interaction.

On-line interaction calculator for Rydberg atoms

N. Sibalic



<https://arc-alkali-rydberg-calculator.readthedocs.io/en/latest/>

S. Weber

[Docs](#) » Pairinteraction - A Rydberg Interaction Calculator

Pairinteraction - A Rydberg Interaction Calculator

build passing build passing codecov 67% pypi v0.9.5a0 arXiv 1612.08053
License GPLv3

The *pairinteraction* software calculates properties of Rydberg systems. The software consists of a C++/Python library and a graphical user interface for pair potential calculations. For usage examples visit the [tutorials](#) section of the documentation. Stay tuned by [signing up](#) for the newsletter so whenever there are updates to the software or new publications about pairinteraction we can contact you. If you have a question that is related to problems, bugs, or suggests an improvement, consider raising an [issue](#) on [GitHub](#).

<https://pairinteraction.github.io/pairinteraction/sphinx/html/index.html>

On-line interaction calculator for Rydberg atoms

<https://pairinteraction.github.io/pairinteraction/sphinx/html/index.html>

$61D_{3/2}$

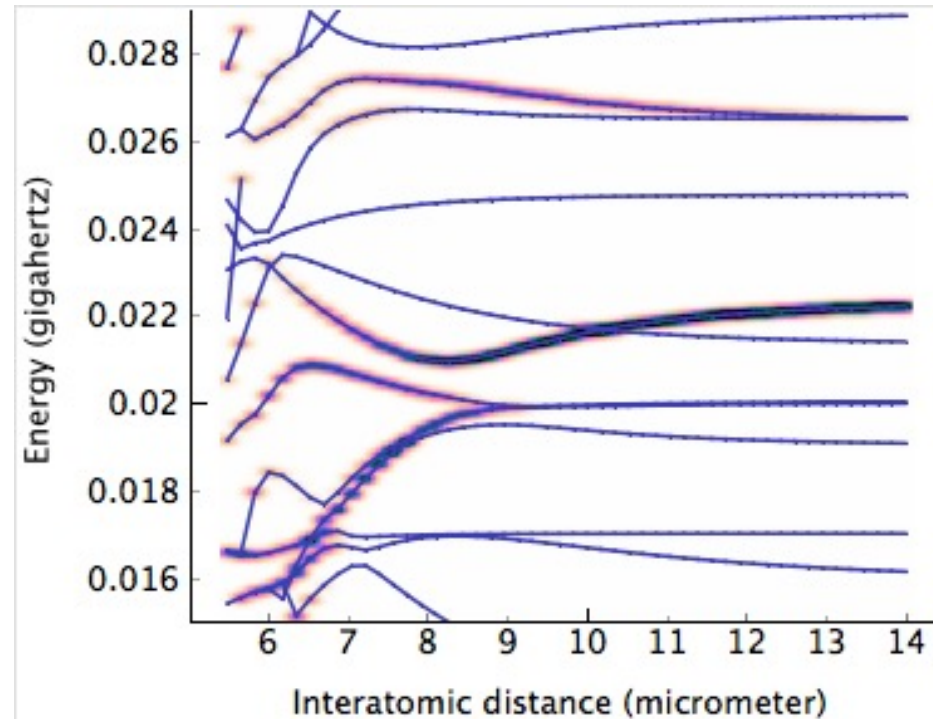
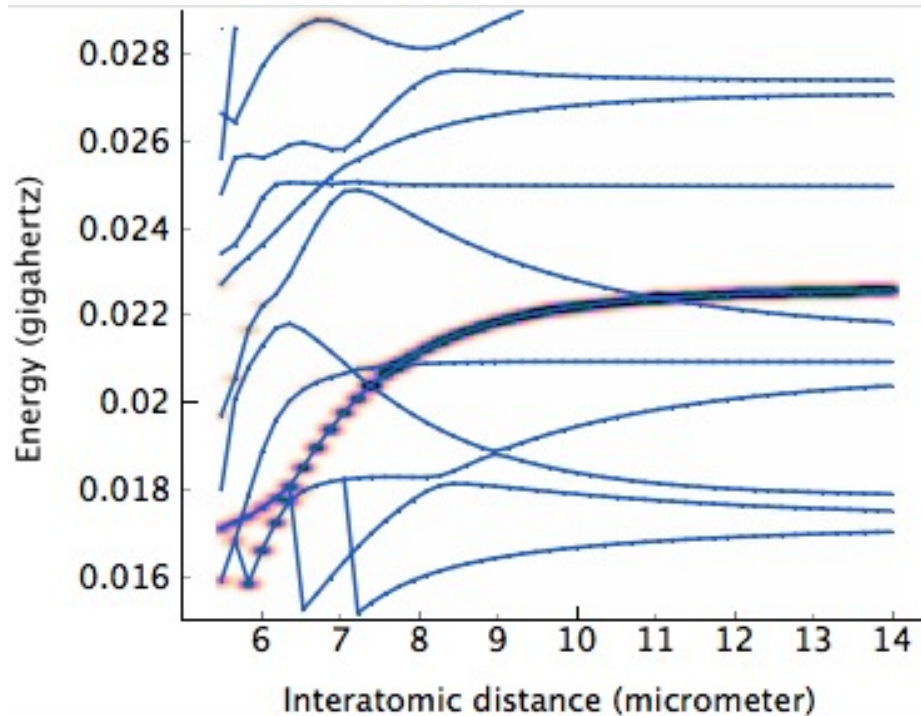
$B = 6.9 \text{ G}$
 $E = 0 \text{ mV/cm}$

Quantization
axis (B)

^{87}Rb

77°

$B = 6.9 \text{ G}$
 $E = 20 \text{ mV/cm}$



Outline

- Lecture 1: Dipole interactions between atoms
- Lecture 2: Arrays of atoms. Basics of Rydberg physics. Rydberg blockade (2-body physics) & gates
- Lecture 3: Many-body physics with Rydberg atoms: spin models and transport