Resonant atomic gases

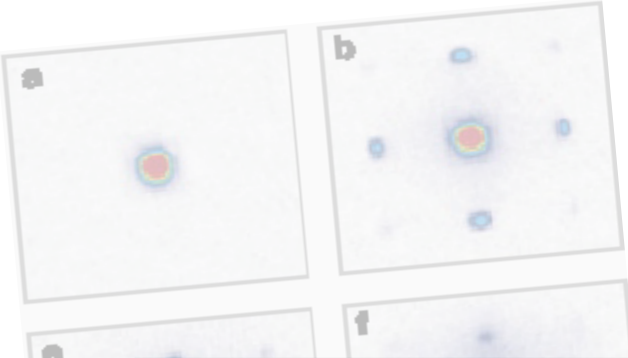
Leo Radzihovsky

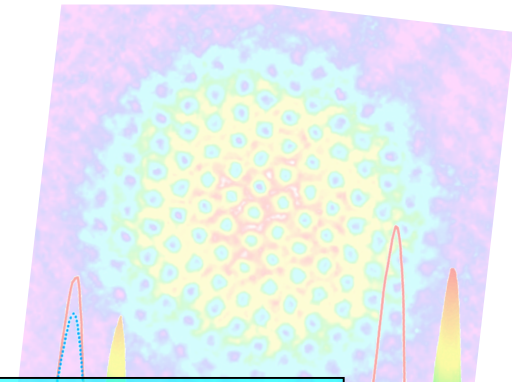
for details see: *Gurarie, L.R., Annals of Physics '07*
Sheehy, L.R., Annals of Physics '07
L.R., Weichman, Park, Annals of Physics '08
Veillette, Sheehy, L.R., PRA, '07
Nikolic, Sachdev, PRA, '07

$: NSF, Packard

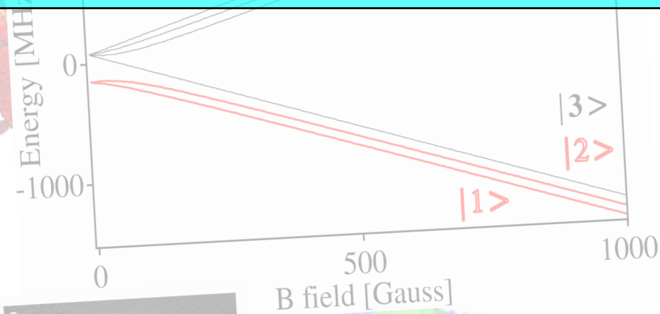
Boulder-2008

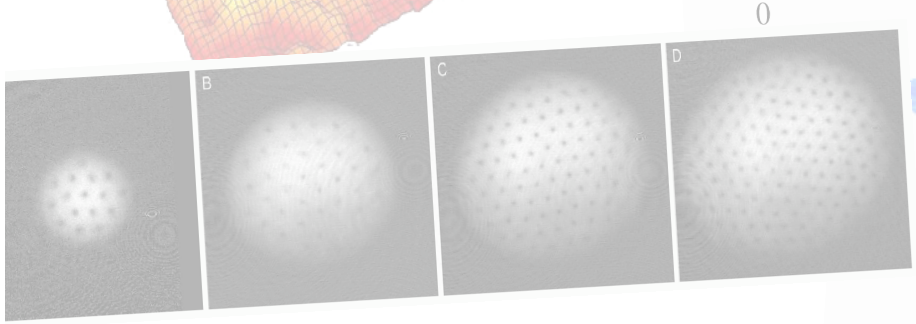
Outline





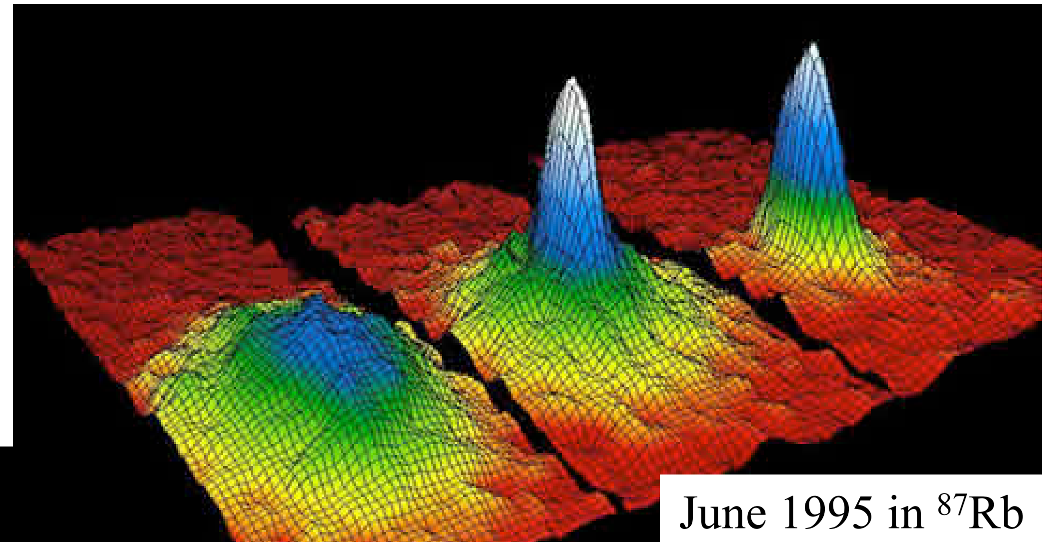
- AMO renaissance overview
- S-wave Feshbach resonant superfluidity
- P-wave Feshbach resonant superfluidity
- Polarized Fermi gases
- Universality of a unitary Fermi gas



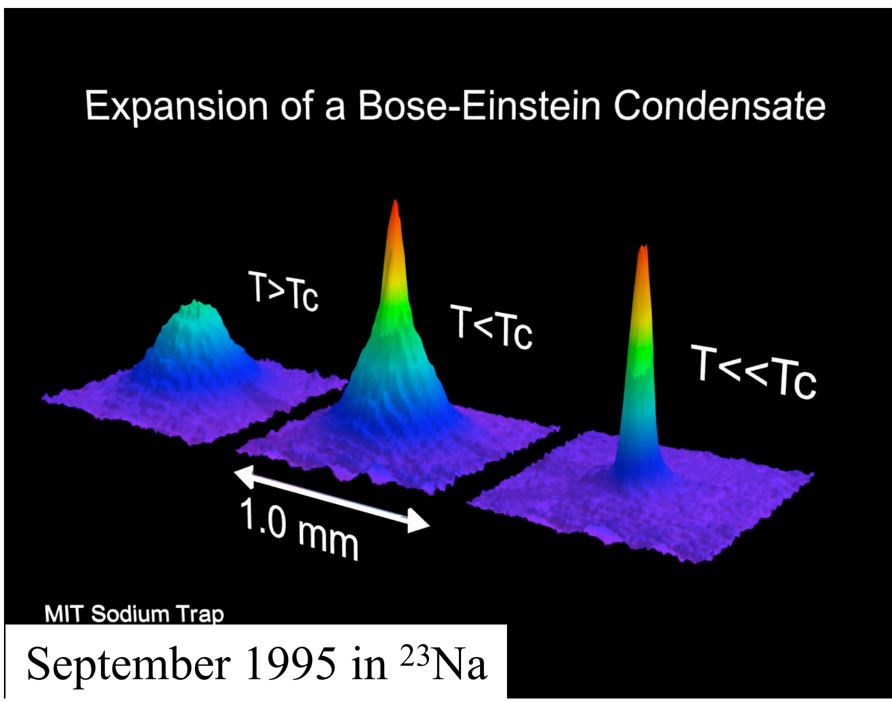


Revolution in AMO physics

- degenerate Bose and Fermi atomic gases

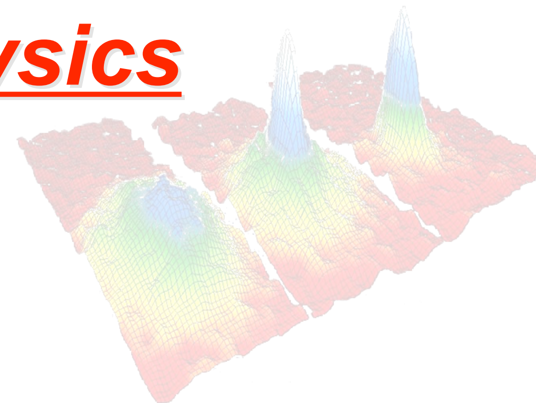


June 1995 in ^{87}Rb

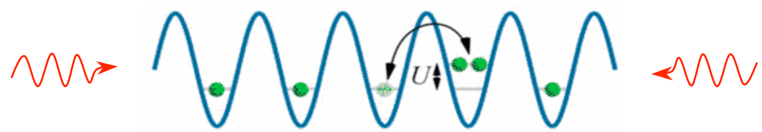


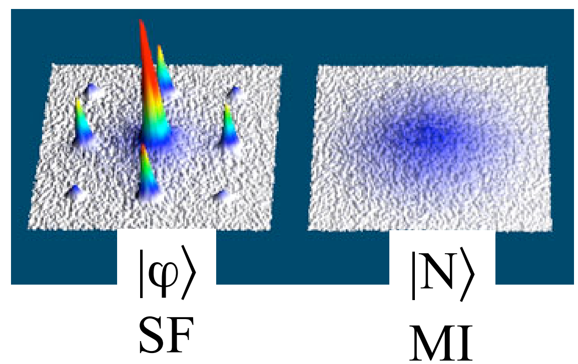
September 1995 in ^{23}Na

Revolution in AMO physics

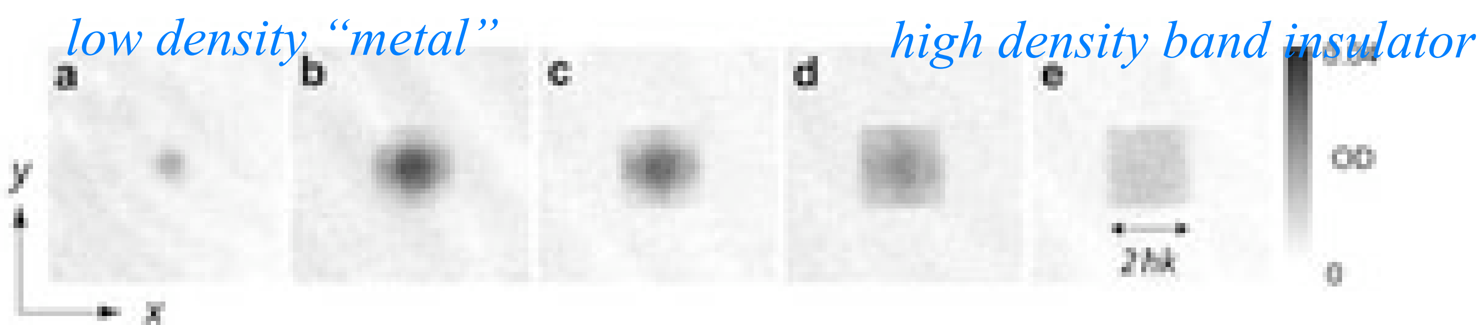


- degenerate Bose and Fermi atomic gases
- optical lattices



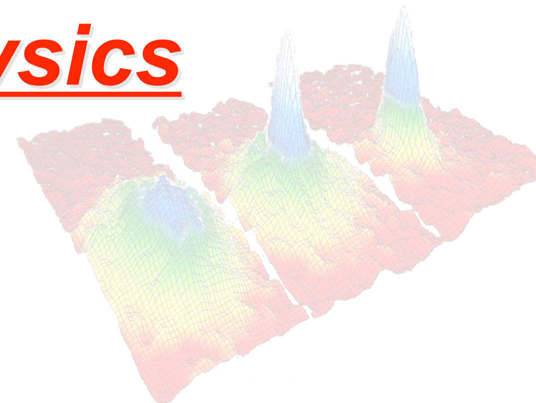


Greiner, et al.



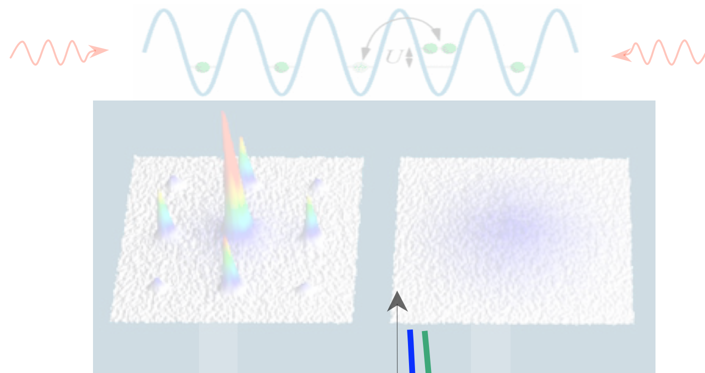
Kohl, Esslinger, et al. '05

Revolution in AMO physics

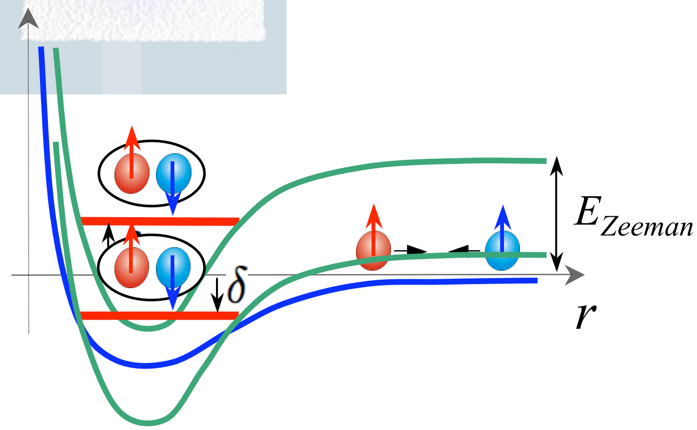


- degenerate Bose and Fermi atomic gases

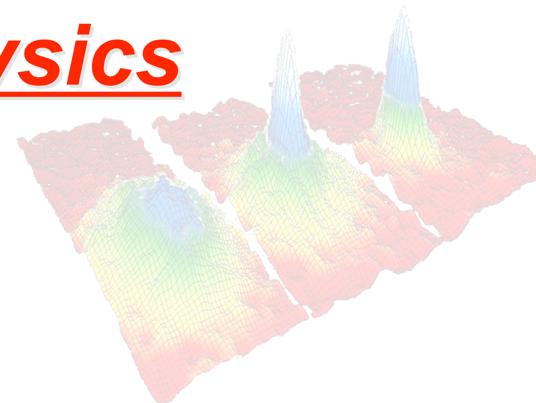
- optical lattices



- Feshbach resonance



Revolution in AMO physics



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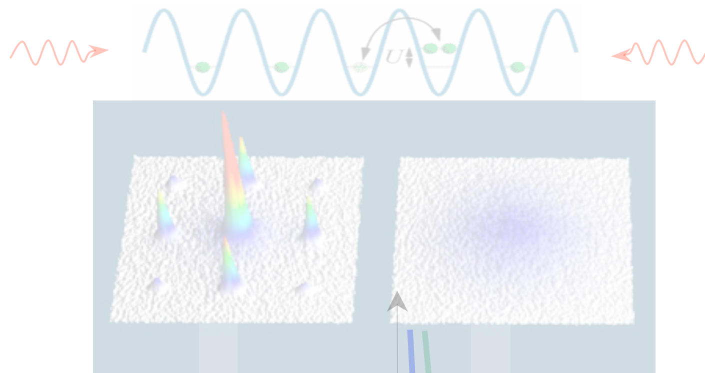
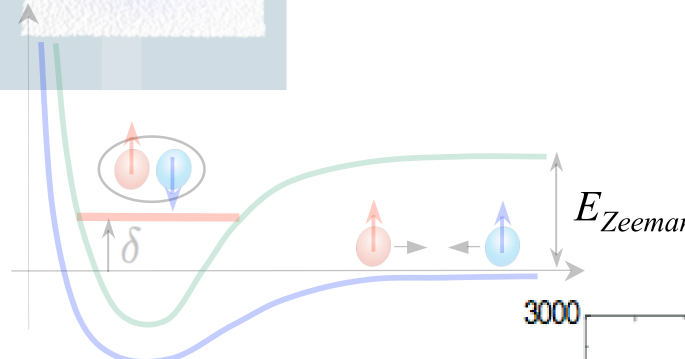
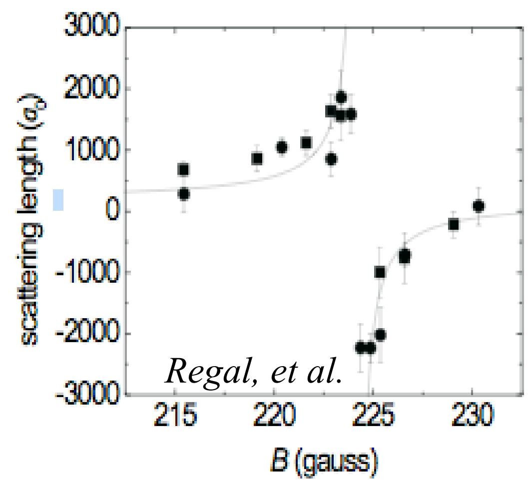
- optical lattices

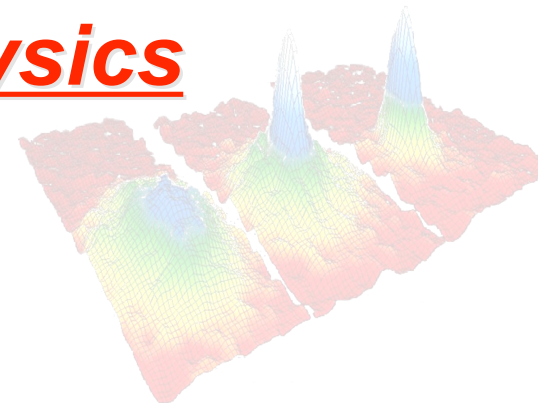


- Feshbach resonance
  - *weak to strong interactions*



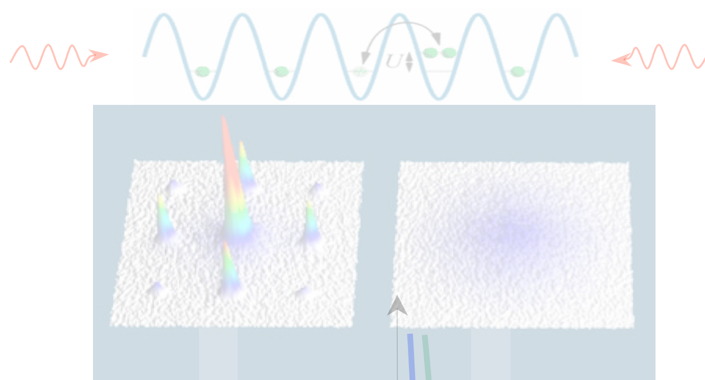


Revolution in AMO physics

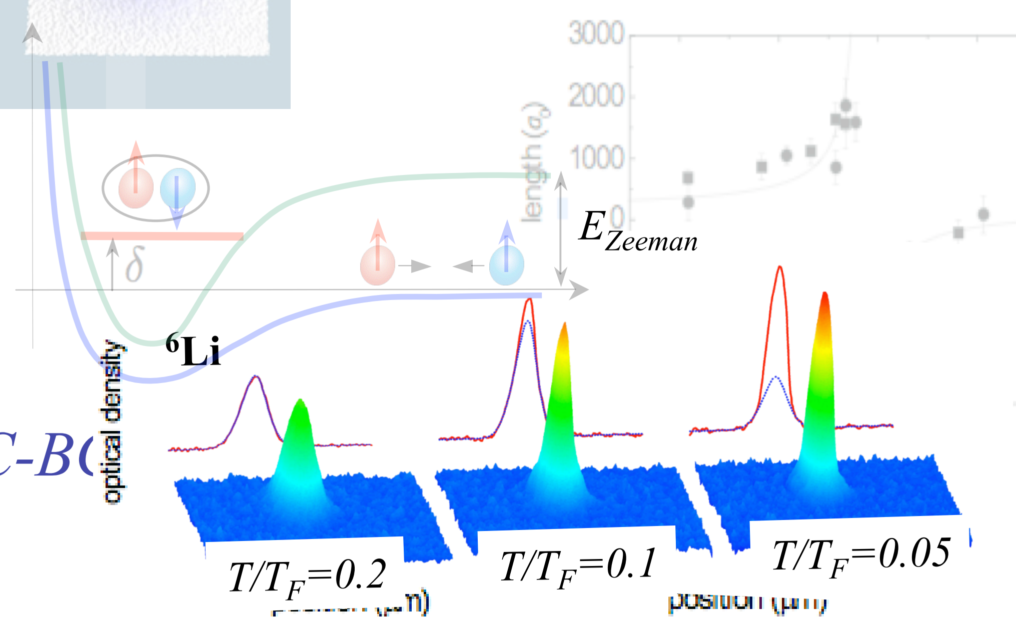


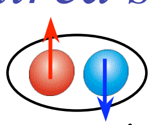
- degenerate Bose and Fermi atomic gases

- optical lattices



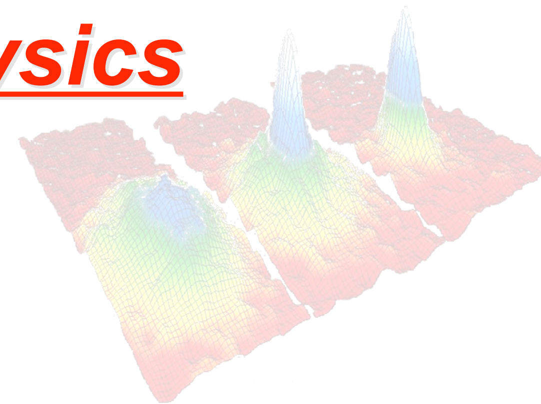
- Feshbach resonance
  - *weak to strong interactions*
  - *paired superfluidity and BEC-B(*





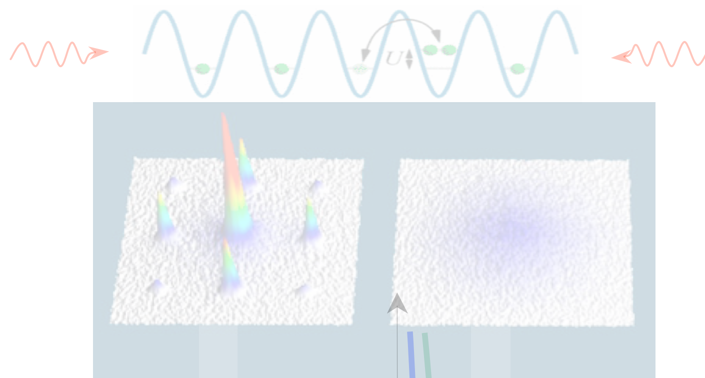
K_2 , Li_2

Revolution in AMO physics

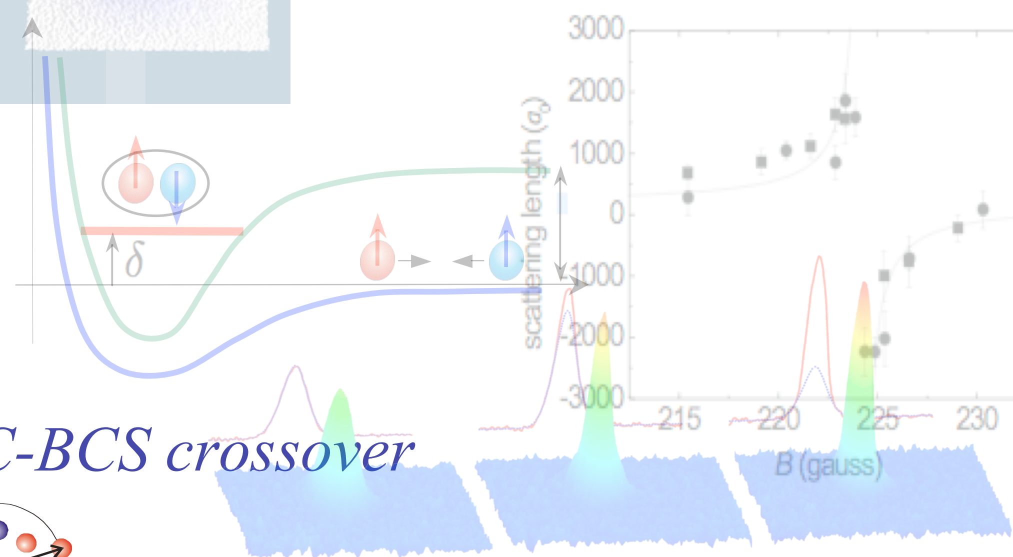


- degenerate Bose and Fermi atomic gases

- optical lattices

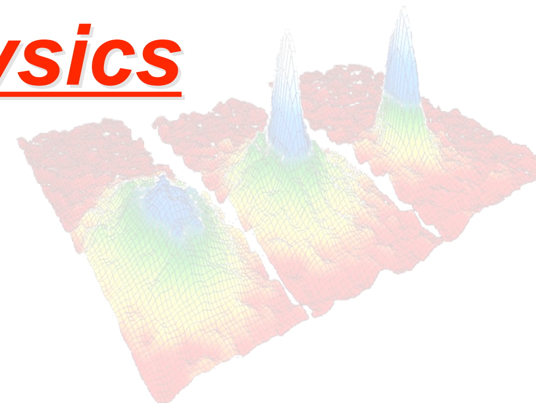


- Feshbach resonance
  - *weak to strong interactions*



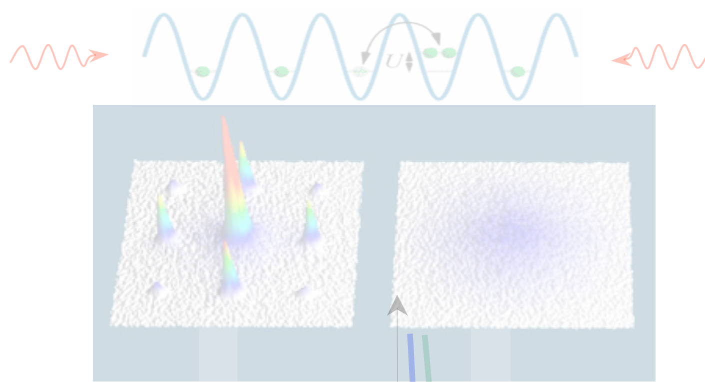
  - *paired superfluidity and BEC-BCS crossover*

Revolution in AMO physics



• degenerate Bose and Fermi atomic gases

• optical lattices



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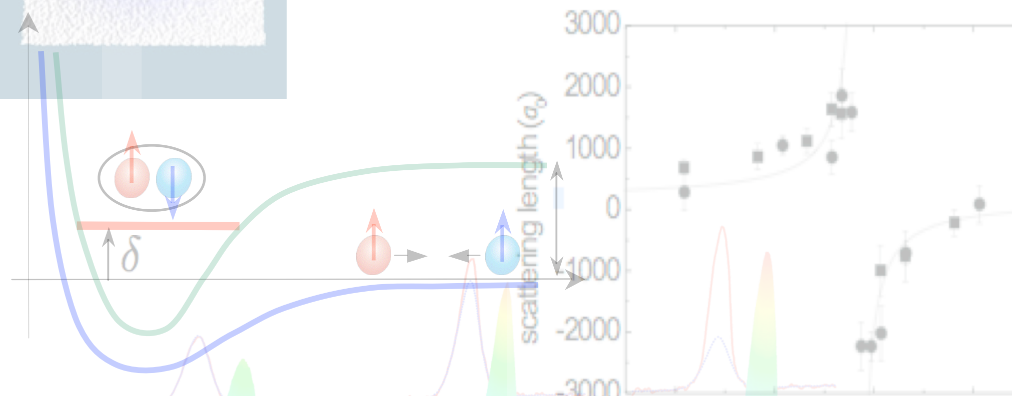
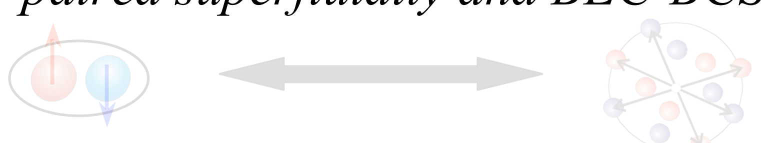
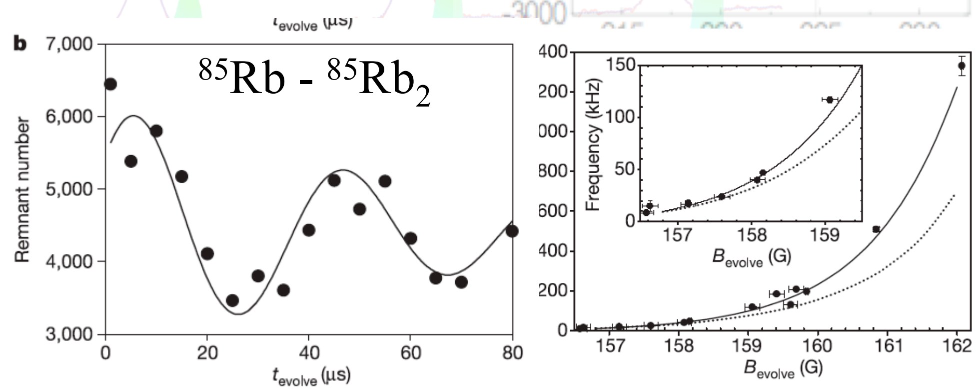


- *weak to strong interactions*
- *paired superfluidity and BEC-BCS cr*

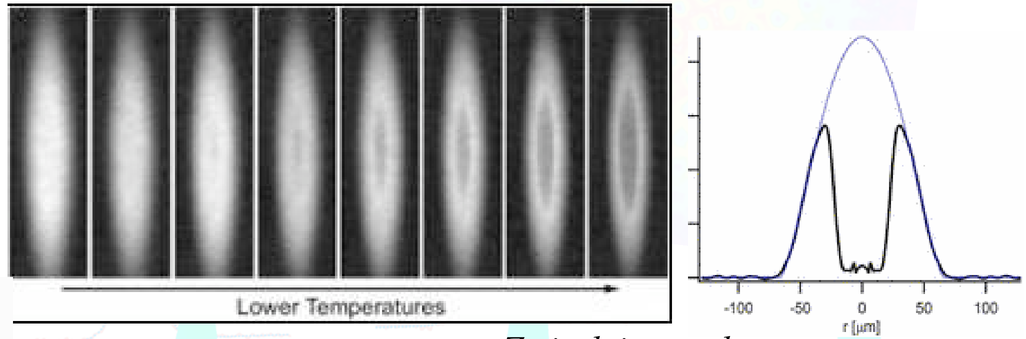


- *quantum nonequilibrium CMP*

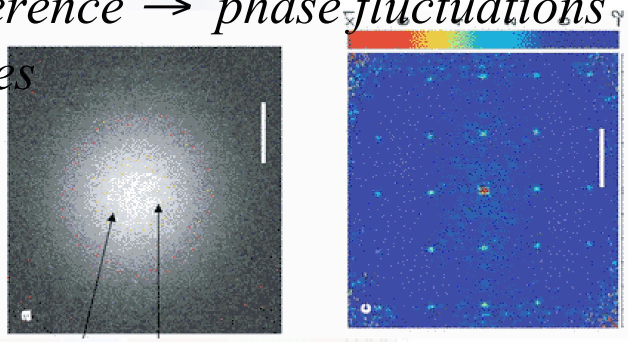


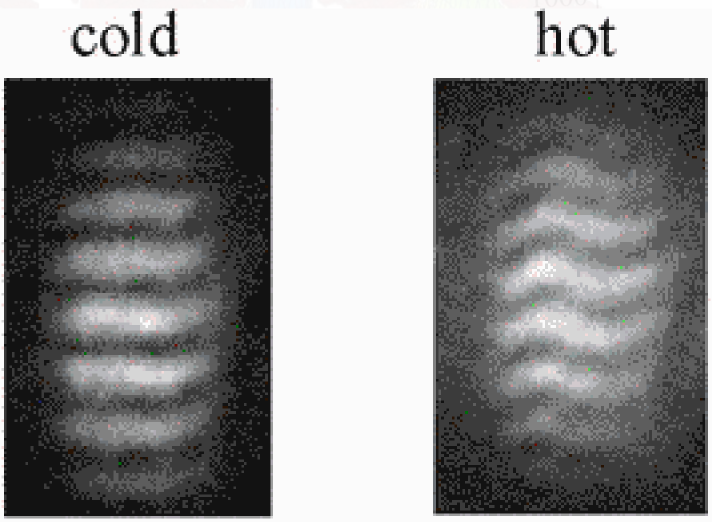
Variety of experimental probes

- Time-of-flight density imaging
  - *momentum distribution function*
  - *scattering length*
  - *temperature*
  - *noise → pairing correlations*
  - *interference → phase fluctuations*
  - *vortices*



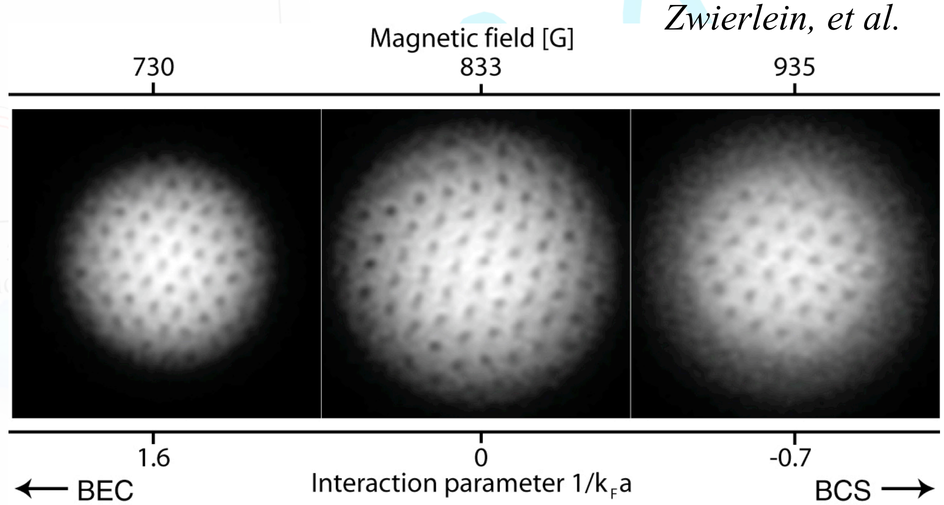
Zwierlein, et al.





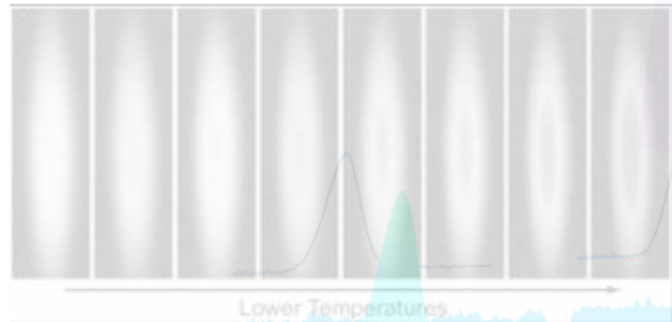
Hadzibabic, et al.

Zwierlein, et al.



Variety of experimental probes

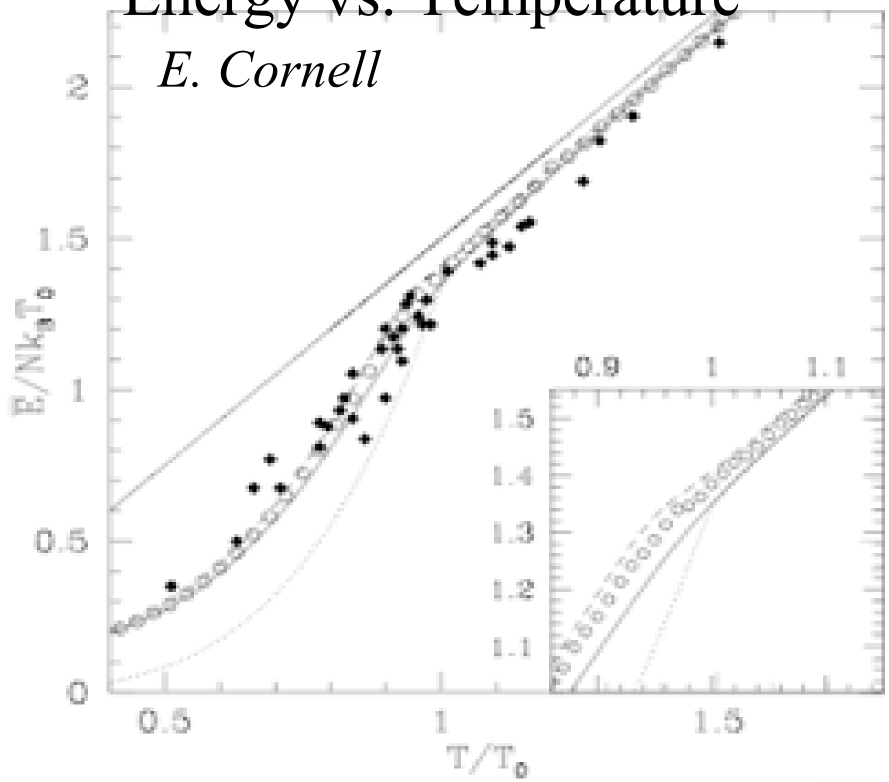
- Time-of-flight density imaging
  - *scattering length*
  - *temperature*
  - *noise → pairing correlations*
  - *quantum phase fluctuations*



- Thermodynamics

Energy vs. Temperature

E. Cornell



Variety of experimental probes

- Time-of-flight density imaging
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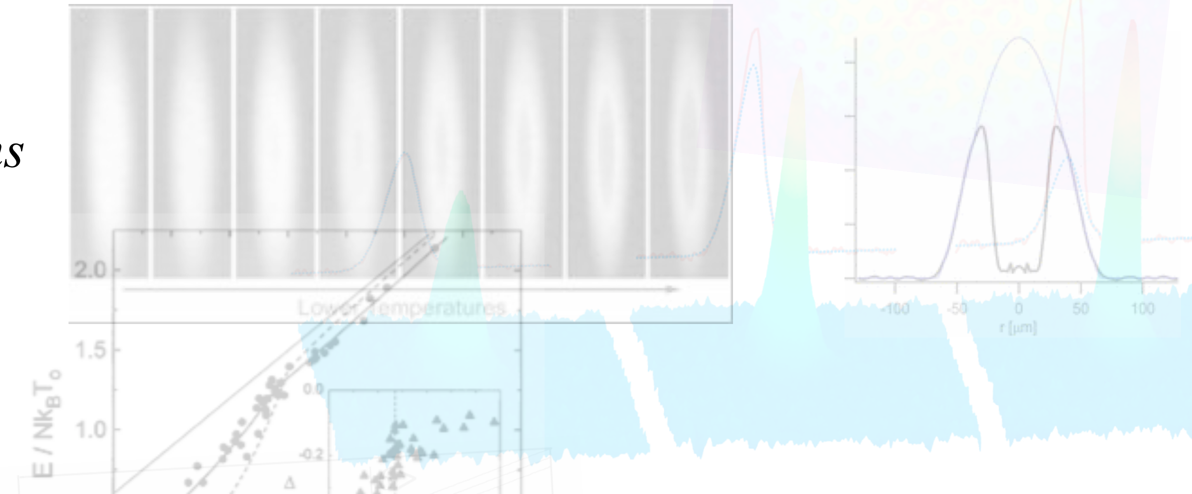
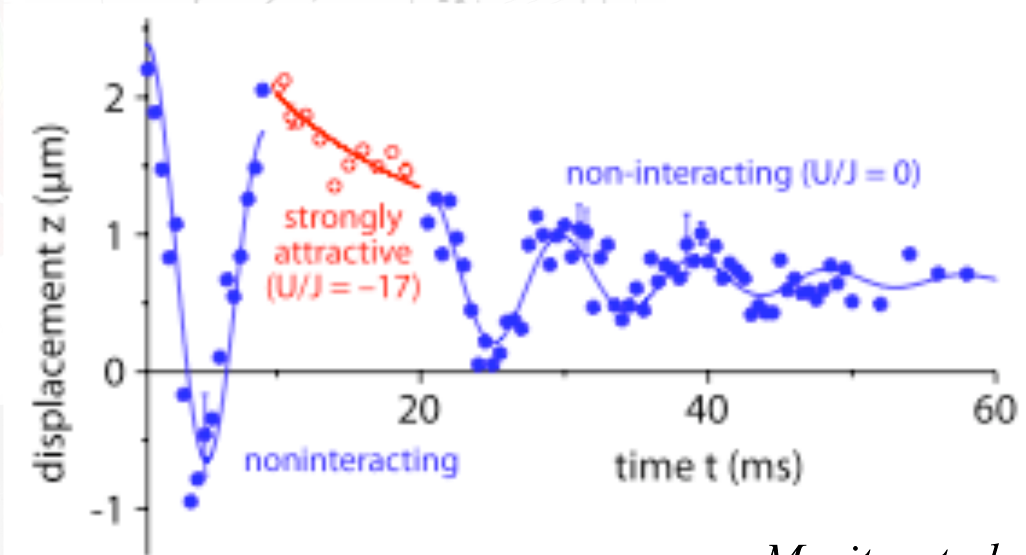


- Thermodynamics

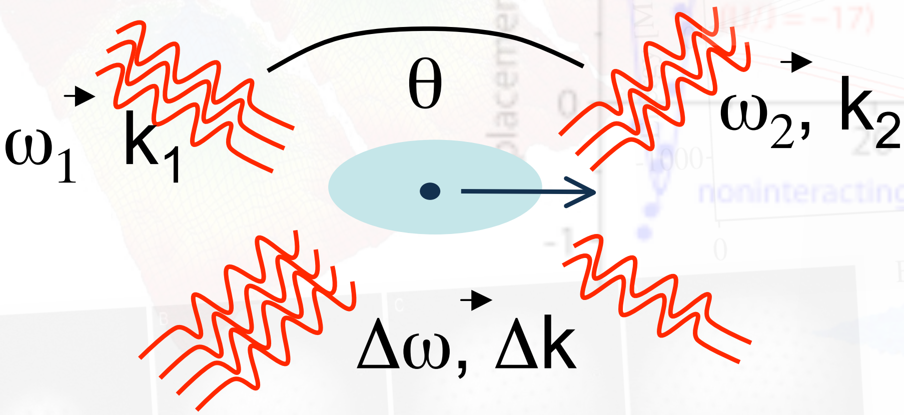
- Transport



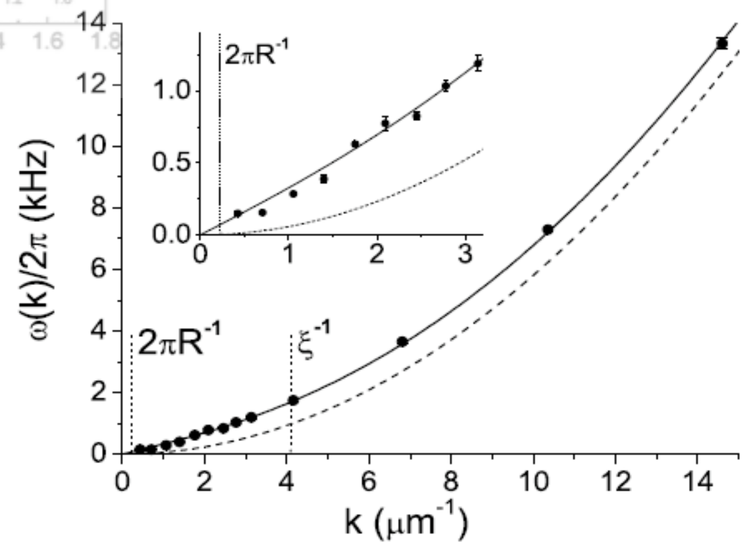
Moritz, et al.

Variety of experimental probes

- Time-of-flight density imaging
  - *scattering length*
  - *temperature*
  - *noise → pairing correlations*
  - *quantum phase fluctuations*
- Thermodynamics
- Transport
- Bragg spectroscopy: *collective excitations*



$$\Delta k = 2\,k \sin \theta/2$$



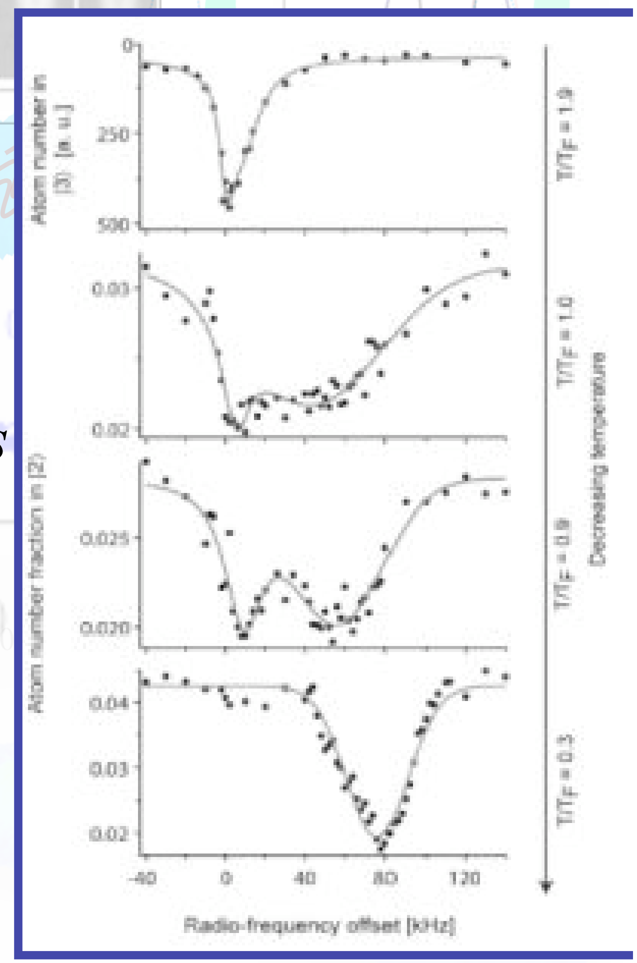
Steinhauer et al., *PRL* **88**, 2002

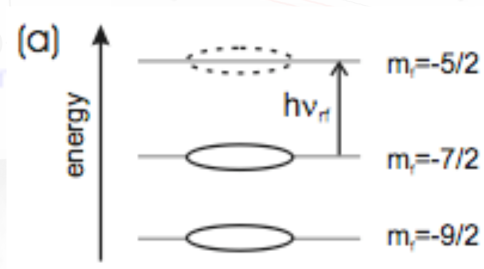
Variety of experimental probes

- Time-of-flight density imaging
  - *scattering length*
  - *temperature*
  - *noise → pairing correlations*
  - *quantum phase fluctuations*
- Thermodynamics
- Transport
- Bragg spectroscopy: *collective excitations*

- RF spectroscopy: *single atom excitations*

Schunck, *et al.*





Regal, Jin'03

<u>*Variety of experimental probes*</u>

- Time-of-flight density imaging
  - *scattering length*
  - *temperature*
  - *noise → pairing correlations*
  - *quantum phase fluctuations*

- Thermodynamics

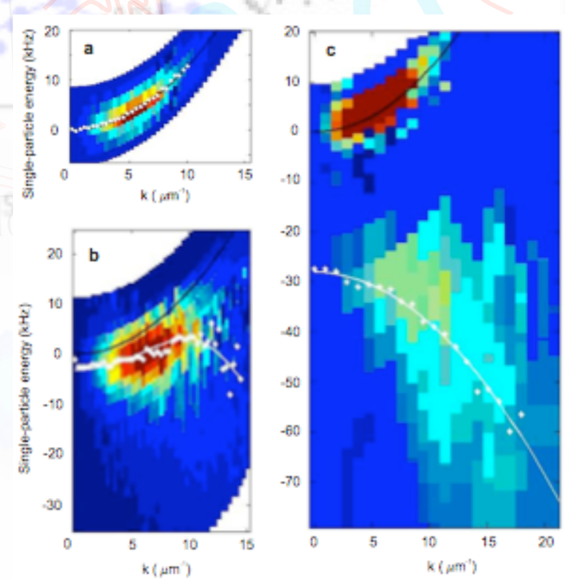
- Transport

- Bragg spectroscopy: *collective excitations*

- RF spectroscopy: *single atom excitations*

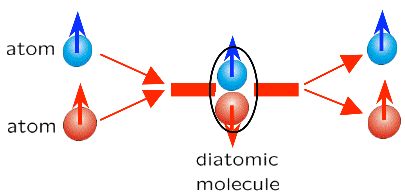
- k-resolved photoemission

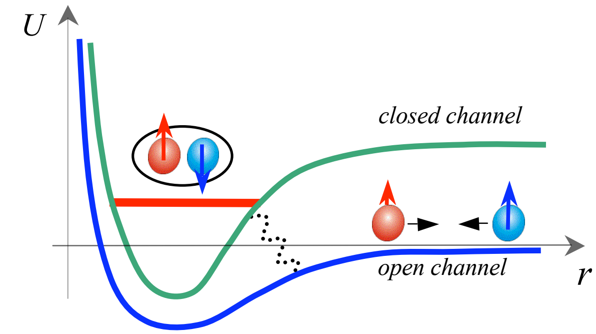
Stewart, et al



S-wave Feshbach resonant scattering

- tunability (strength and sign) of interactions (sudden and adiabatic)

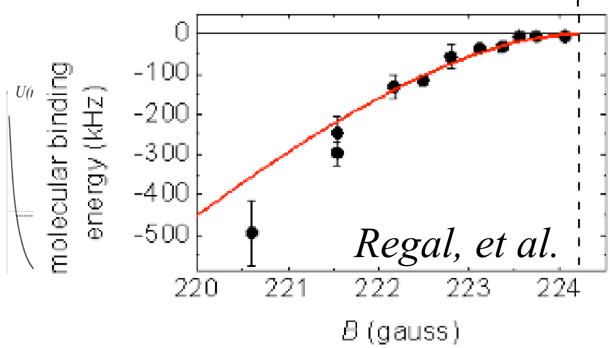


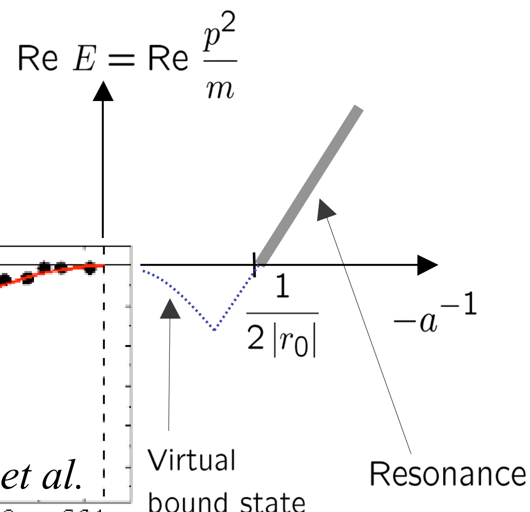


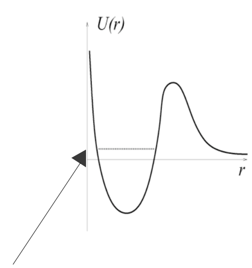
$$\mathcal{H}_{2ch} = \psi^\dagger_\sigma \frac{\hat{p}^2}{2m}\psi_\sigma + \phi^\dagger\left(\frac{\hat{p}^2}{4m} + \epsilon_0\right)\phi - g\phi\psi^\dagger_\uparrow\psi^\dagger_\downarrow$$

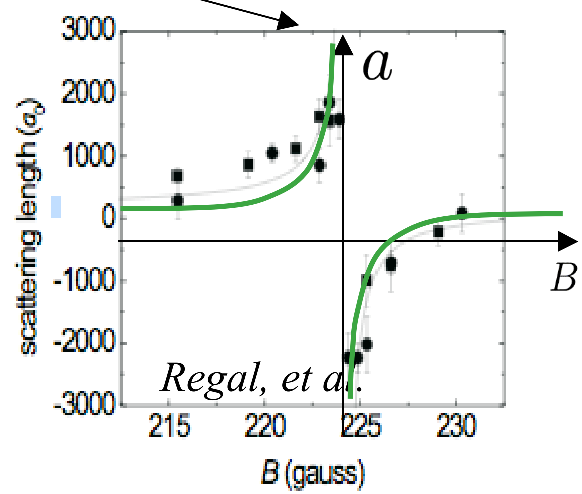
$$\longrightarrow \quad f_s(p) = \frac{1}{-a^{-1} + \frac{r_0}{2}p^2 - ip}, \quad \text{with} \;\; a \sim -\frac{g^2}{\omega_0} \sim a_{bg}\frac{\Delta B}{B_0 - B}, \quad r_0 \sim -\frac{1}{g^2}$$

$$f_s(p) = \frac{1}{F(p^2) - ip}$$



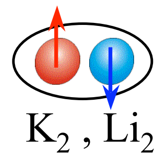




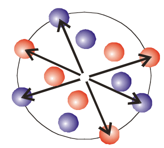


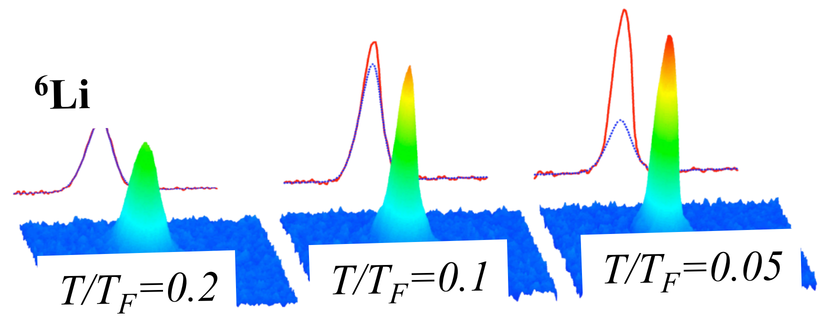
S-wave resonant fermionic superfluidity

- molecular BEC *(Regal, Jin '03)*



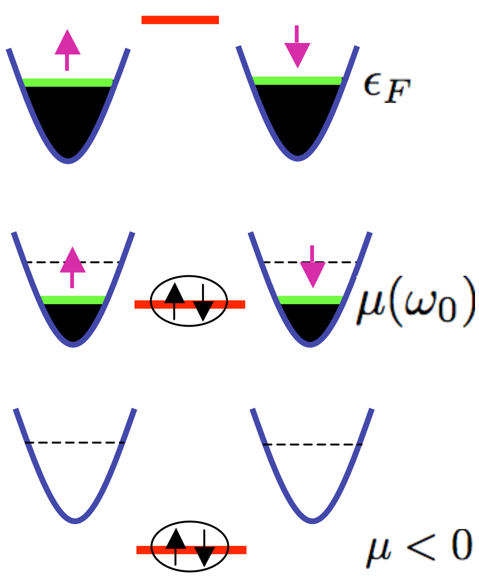
- BCS superfluid *(Regal, Jin 04*
  Zwierlein, Ketterle '04)

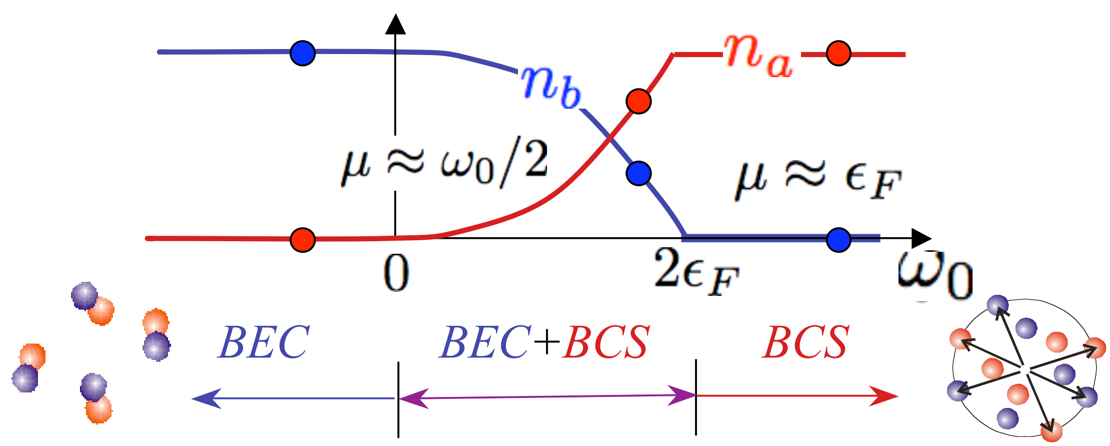




- BCS-BEC crossover:

$$\mathcal{H}_{2ch} = \psi_\sigma^\dagger\left(\frac{\hat{p}^2}{2m} - \mu\right)\psi_\sigma + \phi^\dagger\left(\frac{\hat{p}^2}{4m} - 2\mu + \epsilon_0\right)\phi - g\phi\psi_\uparrow^\dagger\psi_\downarrow^\dagger$$





S-wave resonant fermionic superfluidity

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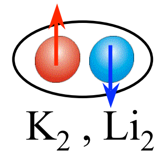

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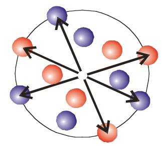

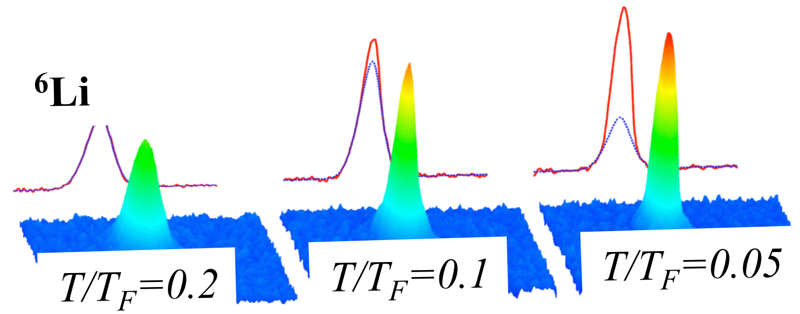

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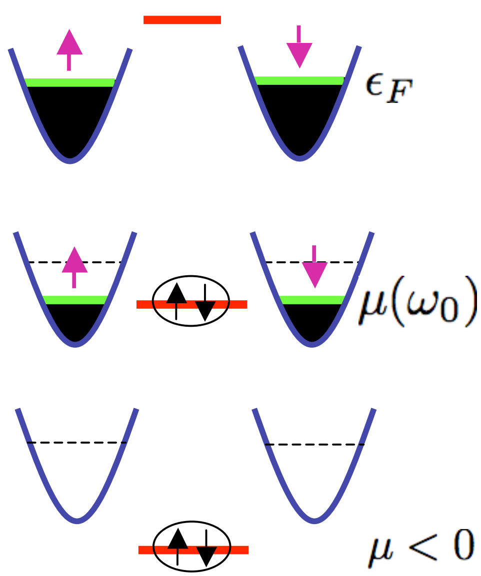

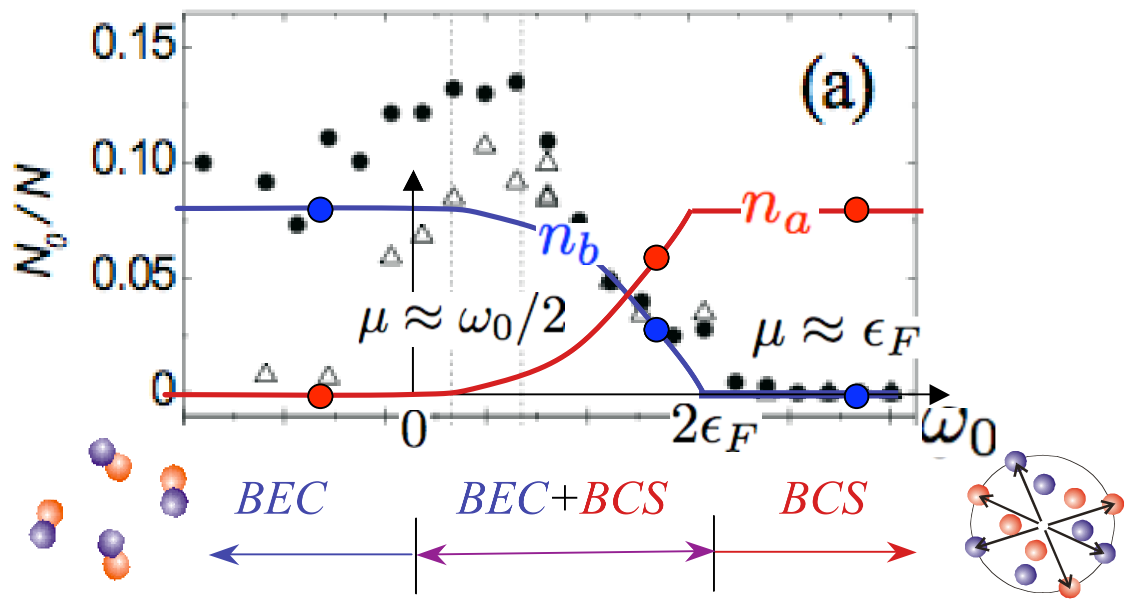

S-wave resonant fermionic superfluidity

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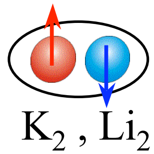

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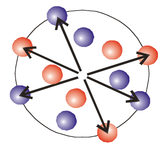

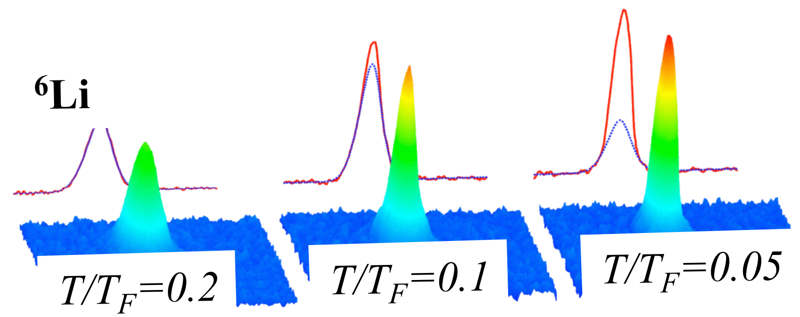

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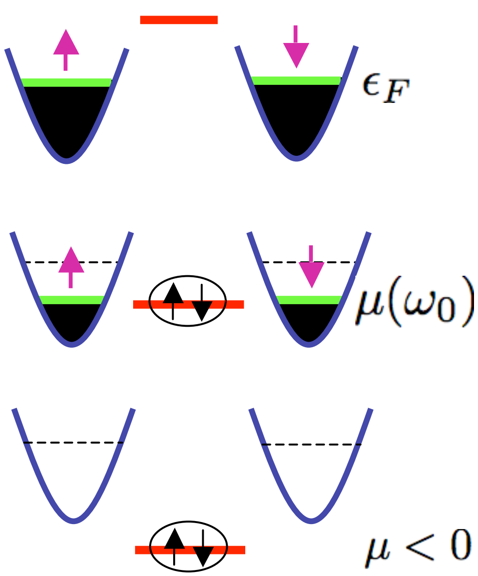
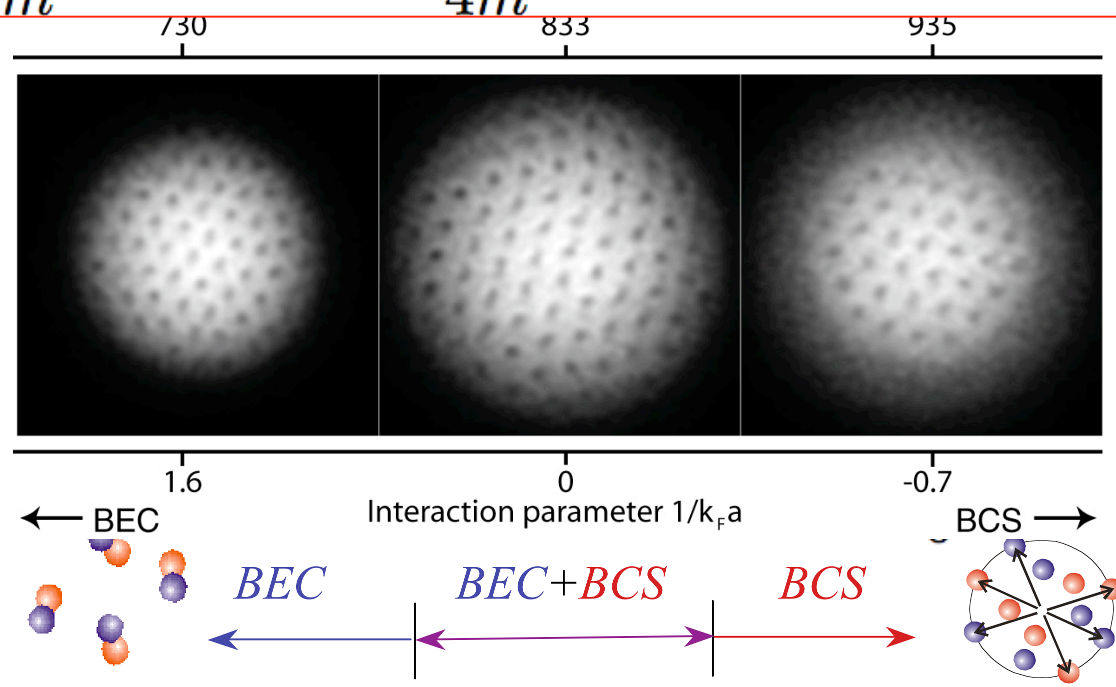
$$\mathcal{H}_{2ch} = \psi_\sigma^\dagger\left(\frac{\hat{p}^2}{2m}-\mu\right)\psi_\sigma + \phi^\dagger\left(\frac{\hat{p}^2}{4m}-2\mu+\epsilon_0\right)\phi - g\phi\psi_\uparrow^\dagger\psi_\downarrow^\dagger$$





S-wave resonant superfluidity

$$\mathcal{H}_{2ch} = \psi_\sigma^\dagger\left(\frac{\hat{p}^2}{2m} - \mu\right)\psi_\sigma + \phi^\dagger\left(\frac{\hat{p}^2}{4m} - 2\mu + \epsilon_0\right)\phi - g\phi\psi_\uparrow^\dagger\psi_\downarrow^\dagger$$

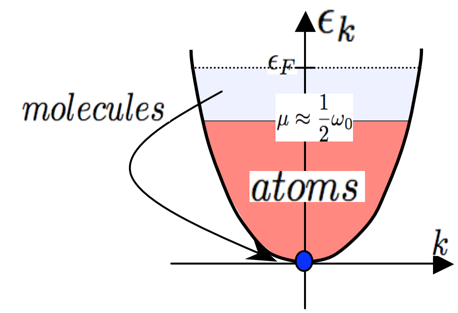
dimensionless coupling: $$\gamma \sim \left(\frac{g\sqrt{n}}{\epsilon_F}\right)^2 \sim \frac{g^2}{\epsilon_F^{1/2}} \sim \frac{1}{r_0 n^{1/3}}$$

$\gamma_{202G}^{^{40}K} \approx 5, \ \Delta B \sim 1G \sim 100\mu K$

$\gamma_{544G}^{^{6}Li} \approx 0.1, \ \Delta B \sim 0.1G \sim 10\mu K$

$\epsilon_F \sim 1\mu K$

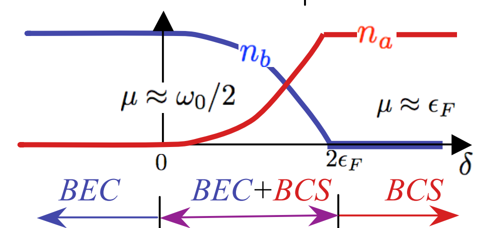
- **narrow** resonance $\gamma \ll 1 \ \rightarrow \ \mathrm{MFT}: \phi(x) = B$

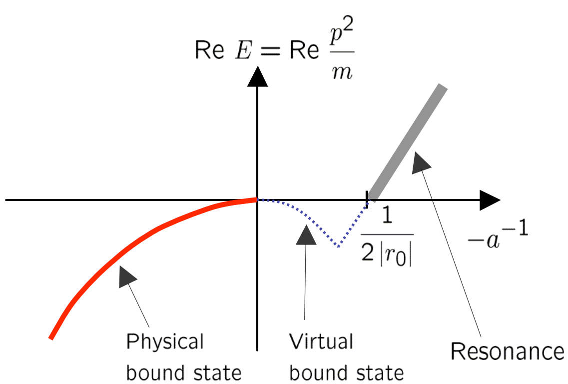


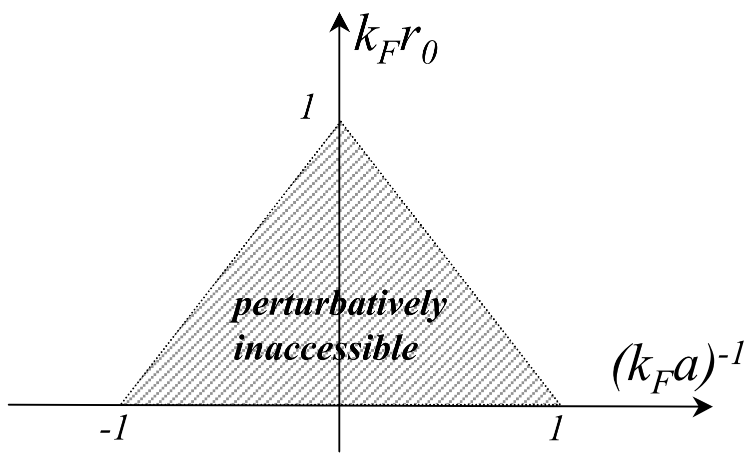
- **broad** resonance $\gamma \gg 1$

Strongly coupled ϕ and ψ

⇨ *MFT quantitatively uncontrolled*



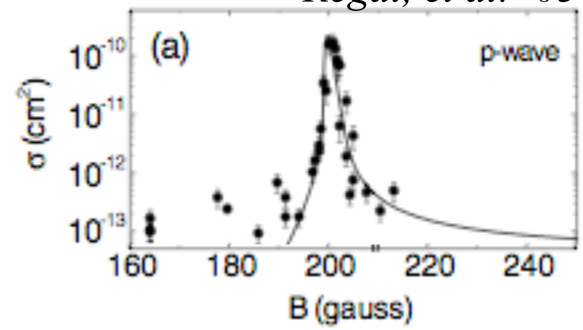




$$\gamma \approx |T_{k_F}| n/\epsilon_F \approx \frac{1}{(k_F a)^{-1} - k_F r_0 + 1}$$

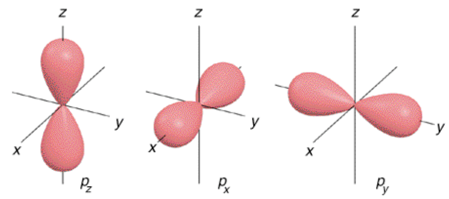
Finite angular momentum superfluidity

Regal, et al. '03

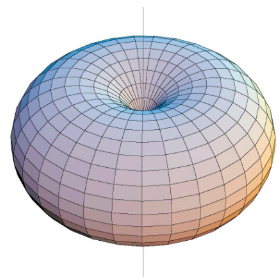


Motivation:

- *p-wave Feshbach resonances exist*
- *examples of ^{3}He and high-Tc superconductors*
- *multiple superfluids phases*

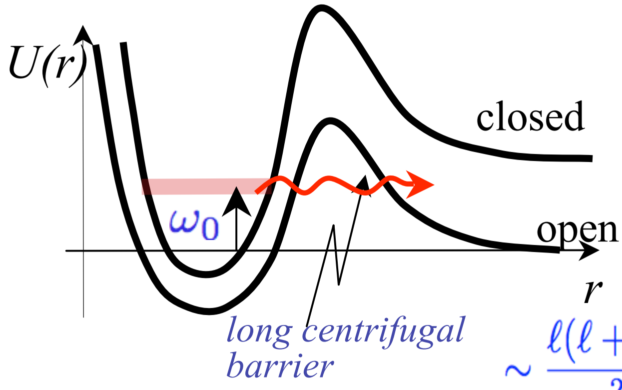


- *anisotropic gap with gapless excitations*

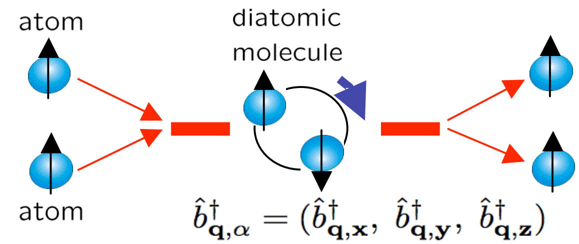


- *conventional (thermal and quantum) and topological phase transitions with detuning*
- *non-Abelian vortex excitations ⇒ topological QC?*

P-wave Feshbach resonant scattering



naturally ***narrow!***



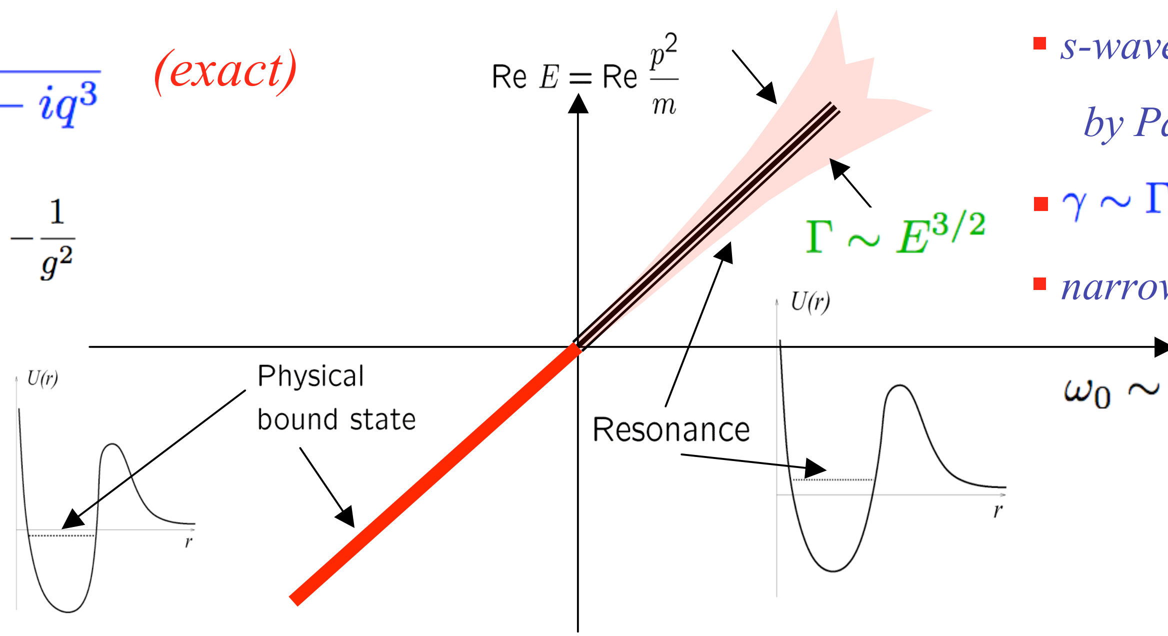
$\sim \frac{\ell(\ell+1)}{r^2}$ ➡ escape (molecular life) time $\tau \sim \Gamma^{-1} \sim E^{-\frac{3}{2}} \gg E^{-1}$, for $E \to 0$

$$\mathcal{H}_{2ch} = \psi^{\dagger}\frac{\hat{p}^2}{2m}\psi + \vec{\phi}^{\dagger}\left(\frac{\hat{p}^2}{4m} + \epsilon_0\right)\vec{\phi} - ig\vec{\phi}\cdot\psi^{\dagger}\nabla\psi^{\dagger}$$

$$f_p = \frac{q^2}{-v^{-1} + \frac{q_0}{2}q^2 - iq^3}$$ *(exact)*

with $v^{-1} \sim -\frac{g^2}{\omega_0}$, $q_0 \sim -\frac{1}{g^2}$

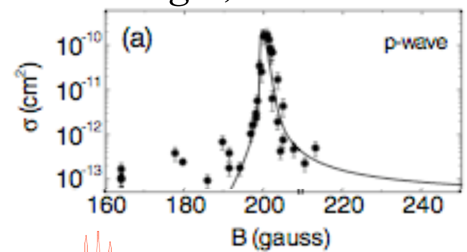
$$f_p(q) = \frac{q^2}{F(q^2) - iq^3}$$



- *s-wave suppressed by Pauli principle*
- $\gamma \sim \Gamma/E \sim E^{\frac{1}{2}} \ll 1$
- *narrows with* $\epsilon_F,\ n$

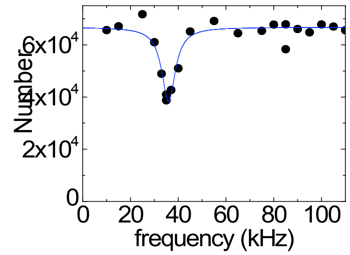
Experimental hopes for p-wave superfluidity

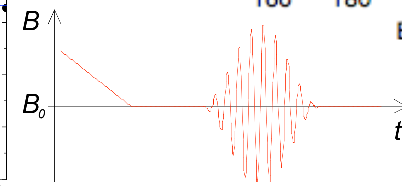
Regal, et al.'03

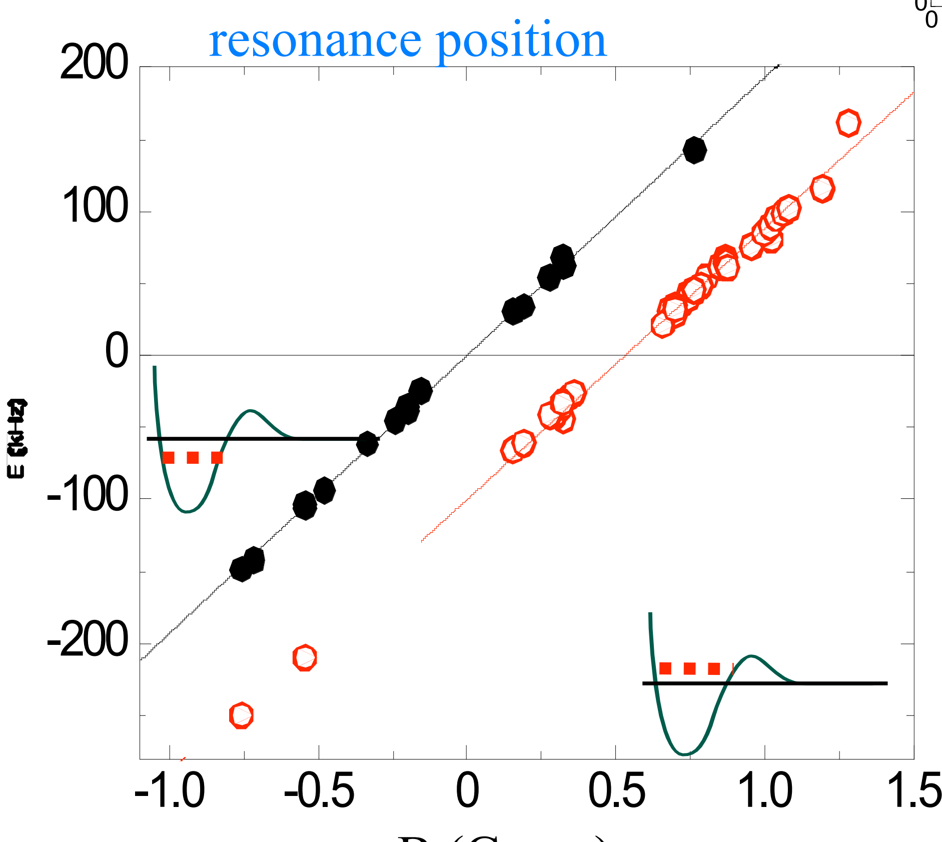


- p-wave Feshbach resonance in ^{40}K, ^{6}Li
- making p-wave molecules:

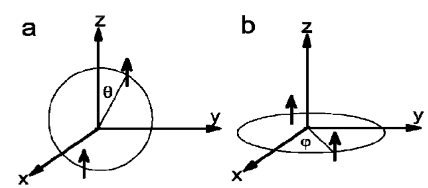
  resonant disappearance of atoms with oscillating B(t)





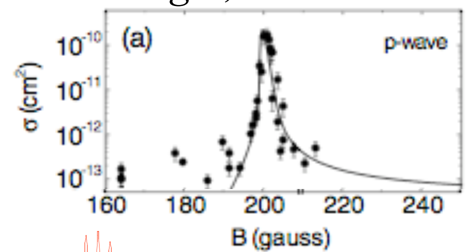


Gaebler, et al.'07



Experimental hopes for p-wave superfluidity

Regal, et al.'03



- p-wave Feshbach resonance in ^{40}K, ^{6}Li
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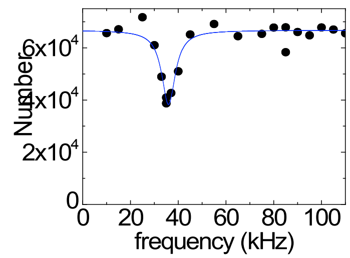
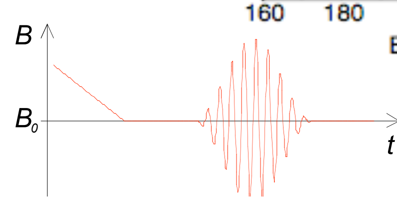
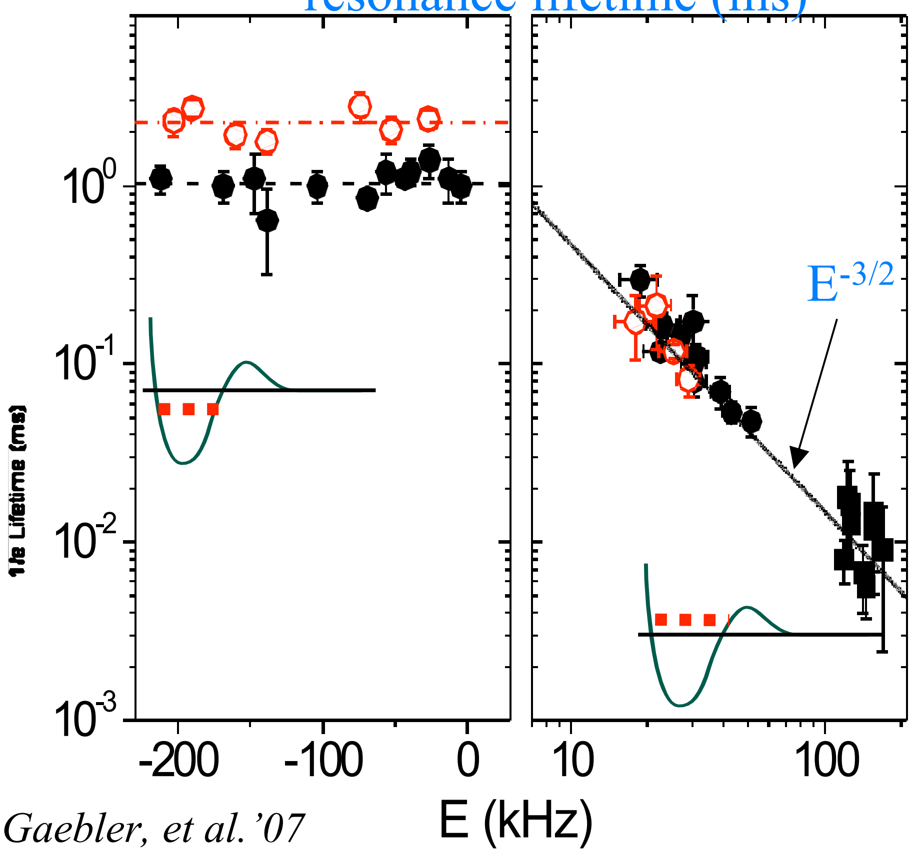
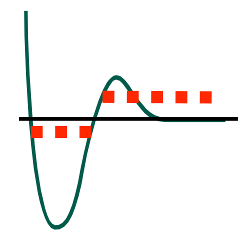




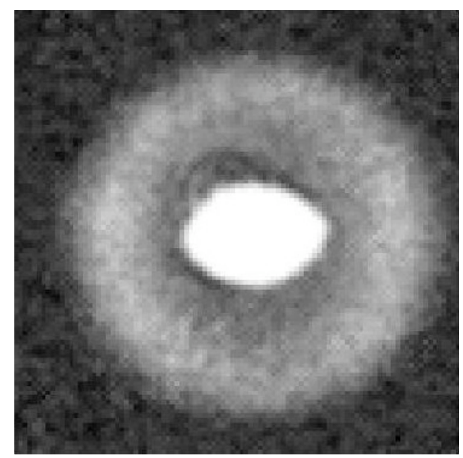
resonance lifetime (ms)



Gaebler, et al.'07



to see molecules:



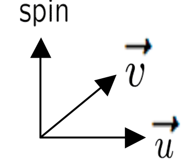
look for energetic atoms

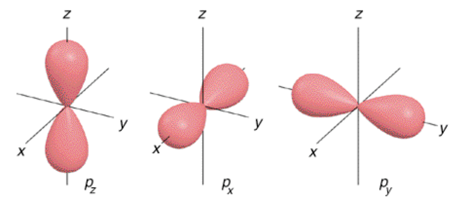
P-wave resonant superfluidity

$$\mathcal{H}_{2ch} = \psi^\dagger \frac{\hat{p}^2}{2m}\psi + \vec{\phi}^\dagger\left(\frac{\hat{p}^2}{4m} + \epsilon_0\right)\vec{\phi} - ig\vec{\phi}\cdot\psi^\dagger\nabla\psi^\dagger$$

dimensionless coupling: $\gamma \sim \left(\frac{g\sqrt{n}k_F}{\epsilon_F}\right)^2 \sim g^2\epsilon_F^{1/2} \sim \frac{n^{1/3}}{q_0}$

- ***narrow*** resonance $\gamma \ll 1 \;\rightarrow\; \mathrm{MFT}: \vec{\phi}(x) = \vec{B}$
- ***complex vector*** order parameter:





$$\vec{B} = \vec{u} + i\,\vec{v} \;\Longleftrightarrow\; \psi_0 = B_z,\; \psi_\pm = \pm(B_x \pm iB_y)$$

$$\vec{B}\cdot\vec{k} = \sum_{m=0,\pm k}\psi_m Y_{1,m}(\hat{k})k$$

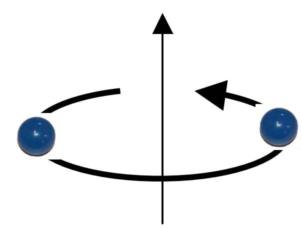
- sample states:

$v = 0 \;\Longleftrightarrow\; |m=0\rangle$ along $\vec{u}$

$(k_x \quad \beta - state\ in\ {}^3He)$

$u = v \;\Longleftrightarrow\; |m=1\rangle$ along $\vec{u}\times\vec{v}$

$(k_x + ik_y \quad "axial"\ Anderson-Morel\ state\ in\ {}^3He)$



$L_z = \pm 1, 0$

Mean-field theory $(\gamma \sim g^2 \epsilon_F^{1/2} \ll 1)$

$$H_p = \sum_{\mathbf{k}} \frac{k^2}{2m} a_{\mathbf{k}}^\dagger a_{\mathbf{k}} + \sum_{\mathbf{q},\alpha} (\frac{q^2}{4m} + \epsilon_{0\alpha}) b_{\mathbf{q},\alpha}^\dagger b_{\mathbf{q},\alpha} + \sum_{\mathbf{k}} [\Delta_{\mathbf{k}}(\vec{B}) a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger + h.c.]$$

- *superfluid ground state:*

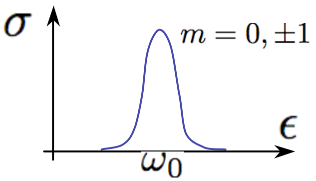
molecular BEC $|\vec{B}\rangle$ + Cooper pairing $|\mathrm{BCS}_{\vec{\mathrm{B}}}\rangle = \Pi_k (u_{\mathbf{k}} + v_{\mathbf{k}} a_{-\mathbf{k}}^\dagger a_{\mathbf{k}}^\dagger)|0\rangle$

(closed) *(open)*

- *excitation spectrum:* $H_{ex} = \sum_{\mathbf{k}} E_{\mathbf{k}}^{(a)} \alpha_{\mathbf{k}}^\dagger \alpha_{\mathbf{k}} + \sum_{\mathbf{k},\alpha} E_{\mathbf{k},\alpha}^{(m)} \beta_{\mathbf{k},\alpha}^\dagger \beta_{\mathbf{k},\alpha}$

$$E_{\mathbf{k}}^{(a)} = \sqrt{\varepsilon_{\mathbf{k}}^2 + \Delta_{\mathbf{k}}^2}, \qquad E_{\mathbf{k},\alpha}^{(m)} = \sqrt{\epsilon_{\mathbf{k}}^2 + \mu_\alpha \epsilon_{\mathbf{k}}} \qquad \textit{with gap}: \quad \Delta_k = 2g|\vec{B} \cdot \vec{k}|$$

- $\vec{B}, \quad n_b, \quad , n_a, \quad \mu \quad$ *determined by :*
  - energy minimization (gap equation) $\rightarrow \dfrac{\partial E(\vec{B})}{\partial B_\alpha} = 0$
  - atom number equation $\rightarrow 2n_b + n_a = n$



Isotropic resonance at T=0 $(\omega_\alpha = \omega_0)$

$$E = (u^2+v^2)\left[\omega_0 - 2\mu + a_1 \ln\{a_0(u+v)\}\right] + a_1 \frac{u^3+v^3}{u+v} + a_2\left[\left(u^2+v^2\right)^2 + \frac{1}{2}\left(u^2-v^2\right)^2\right]$$

BCS

BEC : $(\vec{B}^* \cdot \vec{B})^2 + \frac{1}{2}|\vec{B}\cdot\vec{B}|^2$

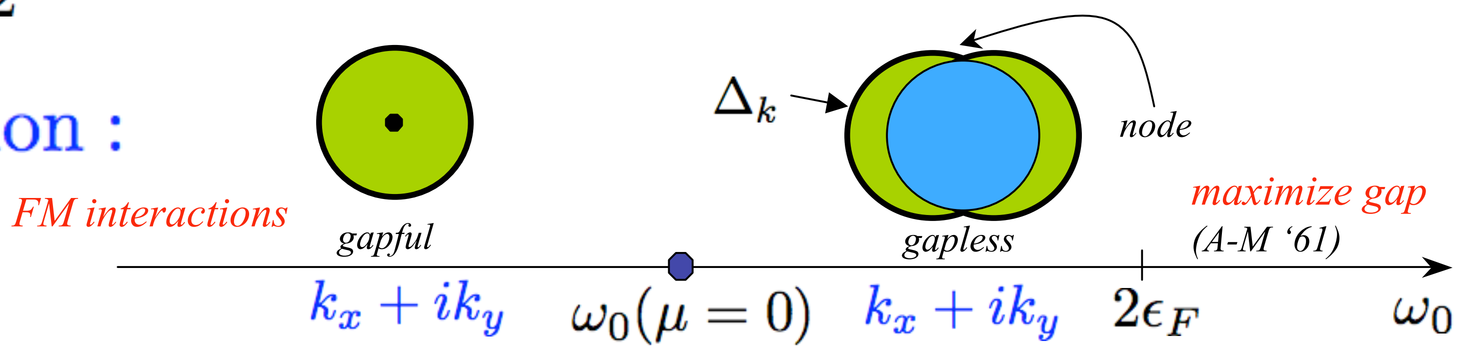
- BCS $(\omega_0 \gg 2\epsilon_F)$:
  - $\mu \approx \epsilon_F + \mathcal{O}(\gamma)$
  - $\dfrac{E_{k_x+ik_y}}{E_{k_x}} = \dfrac{1}{2}e > 1$

  $\Longrightarrow$ $(Anderson - Morel\ A_1\ phase)$

  $u = v \sim e^{-(\omega_0 - 2\epsilon_F)/\gamma\epsilon_F} \Longrightarrow k_x + ik_y\ (m=1)$

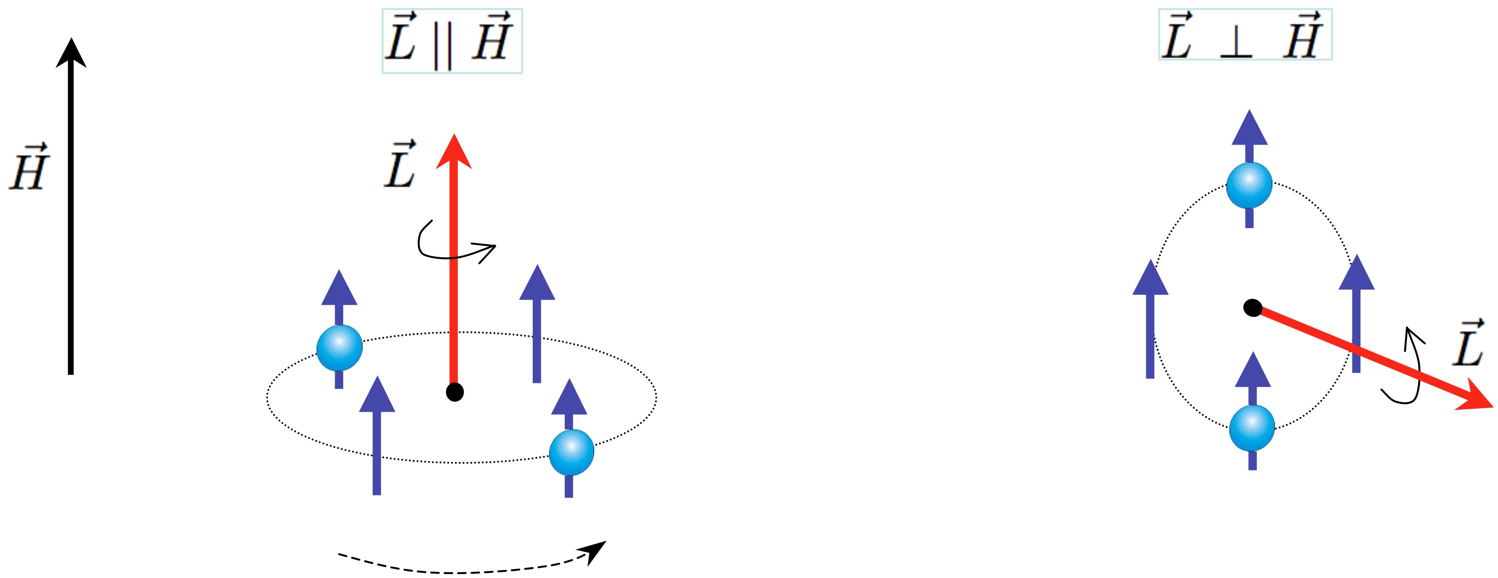
- BEC $(\omega \ll 2\epsilon_F)$:
  - $\mu \approx \frac{1}{2}\omega_0 + \mathcal{O}(\gamma)$ $\Longrightarrow$ $u = v \approx \sqrt{n - n(\omega_0/2)} \Longrightarrow k_x + ik_y\ (m=1)$

- transition :

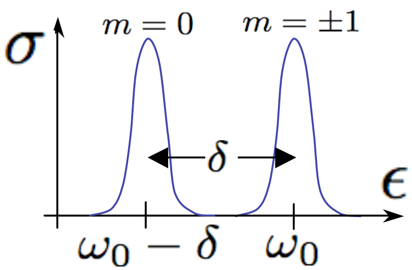


Dipolar-interaction FR splitting

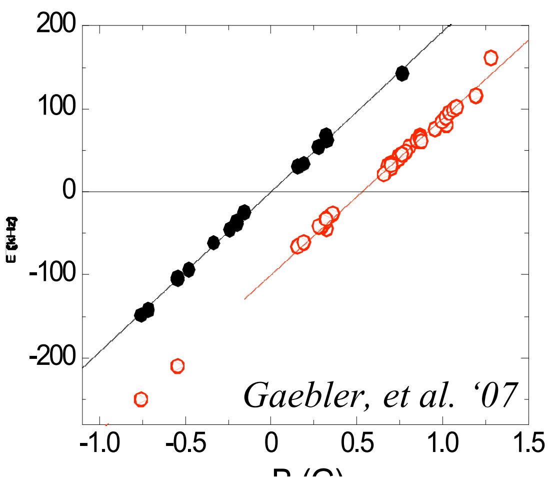
Ticknor, et al '03



$$H_{\mathrm{molecule}} = \omega_{\parallel} b_{\parallel}^{\dagger} b_{\parallel} + \omega_{\perp} \vec{b}_{\perp}^{\dagger} \cdot \vec{b}_{\perp} \qquad \omega_{\parallel} < \omega_{\perp}$$



Ticknor, Regal, et al. '03



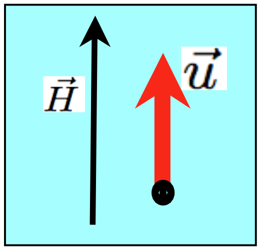
Gaebler, et al. '07

$$E = \sum_{\alpha} (u_{\alpha}^2 + v_{\alpha}^2)[\omega_{\alpha} - 2\mu + \ldots$$

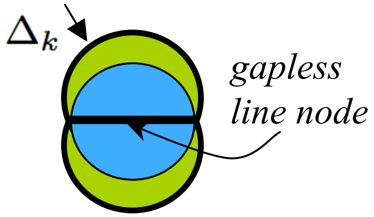
P-wave superfluid phases

$$E=\sum_\alpha (u_\alpha^2+v_\alpha^2)\left[\omega_{0\alpha}-2\mu+a_1\ln\{a_0(u+v)\}\right]+a_1\frac{u^3+v^3}{u+v}+a_2\left[\left(u^2+v^2\right)^2+\frac{1}{2}\left(u^2-v^2\right)^2\right]$$

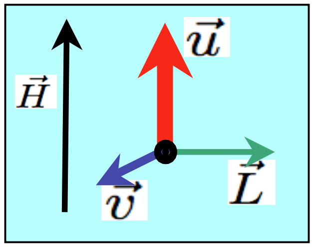
- k_x *- state:* *β-phase of ^{3}He* $(m_{||}=0)$



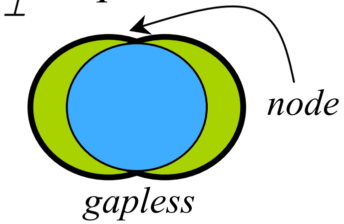
  - $u \approx u_0 e^\delta,\ \ v = 0$
  - equatorial node line for $\mu > 0\ \Rightarrow\ C \sim T^\alpha$, ...
  - fully gapped for $\mu < 0\ \Rightarrow C \sim e^{-|B-B_0|^2/T}$, ...
  - spontaneously broken symmetries: *U(1)*



- $k_x + i\,\sigma\, k_y$*- state:* *"deformed" A_1 -phase of ^{3}He* $(m_\perp=1)$



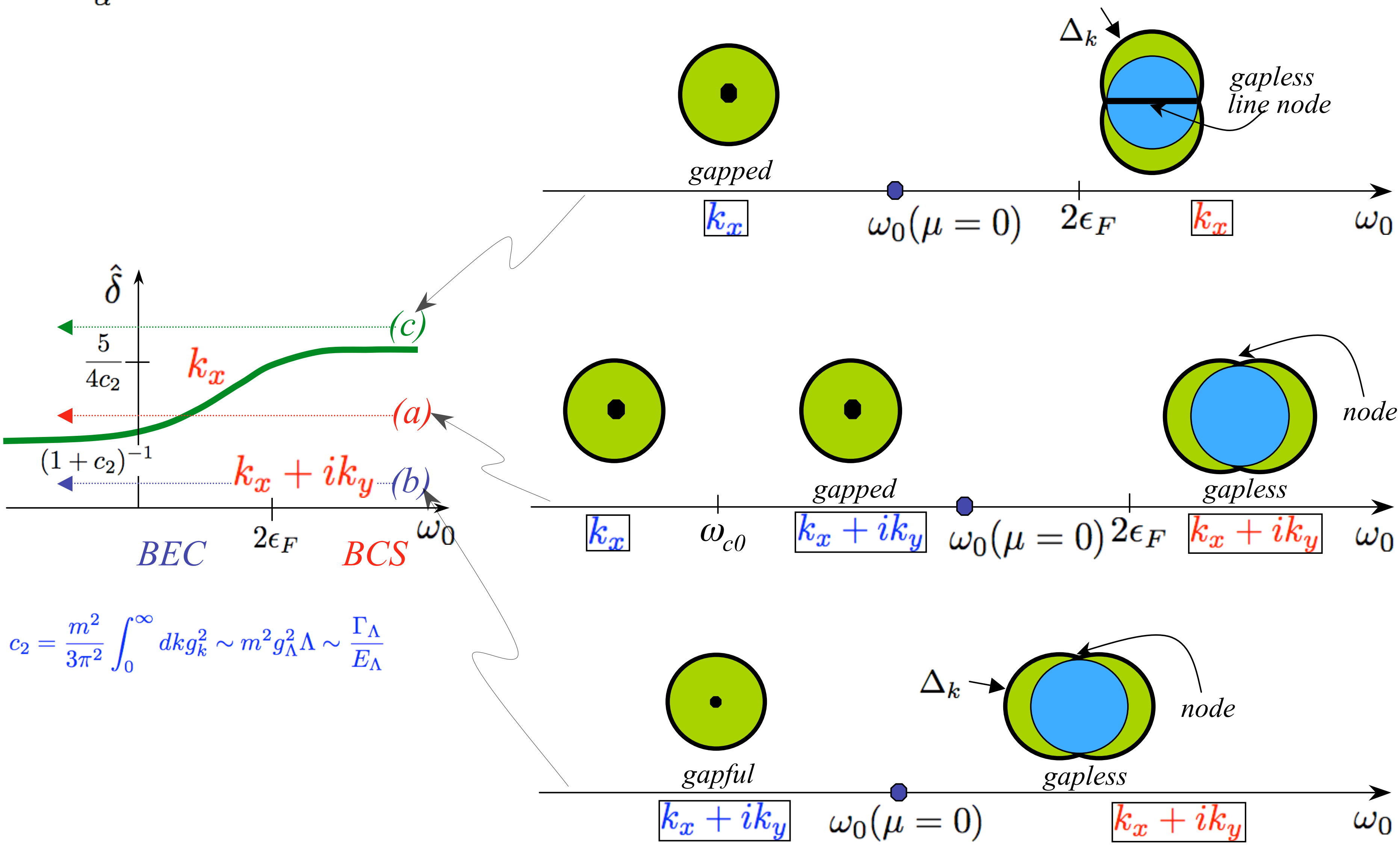
  - $u \approx u_0\,(1+\delta)\,e^{\delta/2},\ v \approx u_0\,(1-\delta)\,e^{\delta/2}\ \Rightarrow\ |m_\perp=1\rangle + \delta\,|m_\perp=-1\rangle$
  - polar point nodes for $\mu > 0 \Rightarrow C \sim T^\alpha$, ...
  - fully gapped for $\mu < 0\ \Rightarrow C \sim e^{-|B-B_0|^2/T}$, ...
  - spontaneously broken symmetries: *U(1), O(2), T*



T=0 phase diagrams

also see Cheng, Yip '05

$$E = \sum_\alpha (u_\alpha^2 + v_\alpha^2)\left[\omega_{0\alpha} - 2\mu + a_1 \ln\{a_0 (u+v)\}\right] + a_1 \frac{u^3+v^3}{u+v} + a_2 \left[(u^2+v^2)^2 + \frac{1}{2}(u^2-v^2)^2 \right]$$

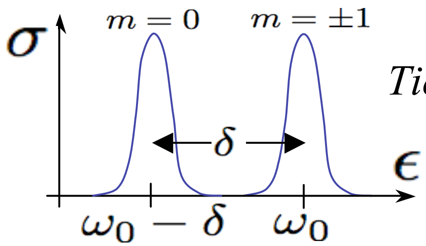


$$c_2 = \frac{m^2}{3\pi^2}\int_0^\infty dk g_k^2 \sim m^2 g_\Lambda^2 \Lambda \sim \frac{\Gamma_\Lambda}{E_\Lambda}$$

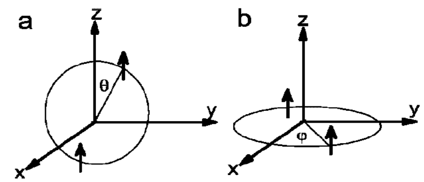
Anisotropic p-wave superfluidity

Gurarie, L.R., Andreev '05
Cheng and Yip '05

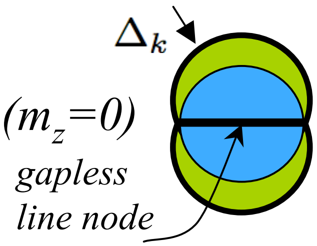
- *resonance splits:*

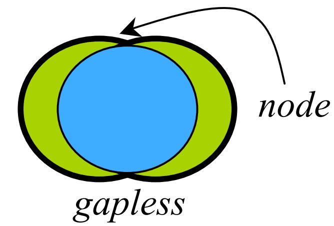


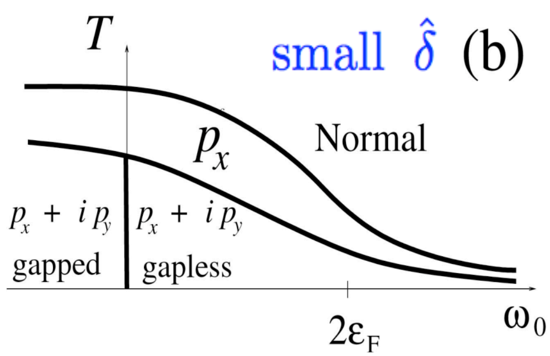
Ticknor, Regal, et al.

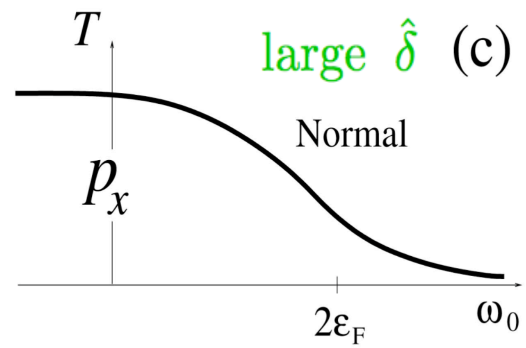


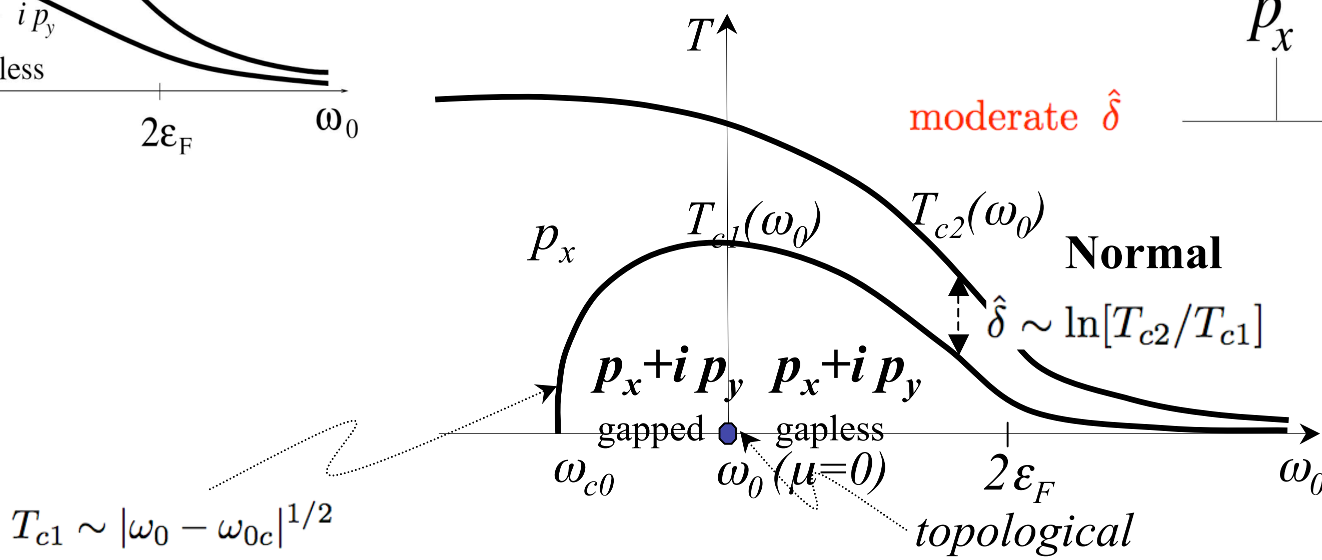
- *two competing states:* p_x $(m_z=0)$ and $p_x+i\,p_y$ $(m_z=\pm1)$





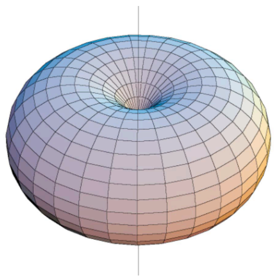


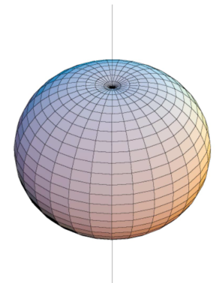




Gapless $\longrightarrow$ gapped superfluid transitions

$p_x + i\, p_y$ $\xrightarrow{\mu \text{ changes sign}}$ $p_x + i\, p_y$





$$E_p = \sqrt{\left(\frac{p^2}{2m} - \mu\right)^2 + g^2\,(p_x^2 + p_y^2)\,|B|^2}$$

$$E_p = \sqrt{\left(\frac{p^2}{2m} + |\mu|\right)^2 + g^2\,(p_x^2 + p_y^2)\,|B|^2}$$

G. E. Volovik, JETP Lett. 80, 343 (2004)

p_x $\xrightarrow{\mu \text{ changes sign}}$ p_x

$$E_p = \sqrt{\left(\frac{p^2}{2m} - \mu\right)^2 + g^2\, p_x^2\,|B|^2}$$

$$E_p = \sqrt{\left(\frac{p^2}{2m} + |\mu|\right)^2 + g^2\, p_x^2\,|B|^2}$$

$p_x + i\, p_y$ superfluid in 2D

- Pfaffian (Moore-Read) state from FQH $\quad |p_x + ip_{y\,BCS}\rangle = \prod_p \left[u_p + v_p a^\dagger_{-p} a^\dagger_p\right] |0\rangle$

$$\Psi\left(z_1, z_2, \ldots, z_{2N}\right) = \sum_P (-1)^P \frac{1}{z_{P_1} - z_{P_2}} \frac{1}{z_{P_3} - z_{P_4}} \cdots \frac{1}{z_{P_{N-1}} - z_{P_N}}$$

- topological classification in terms of u_p and v_p

Anderson's pseudospin

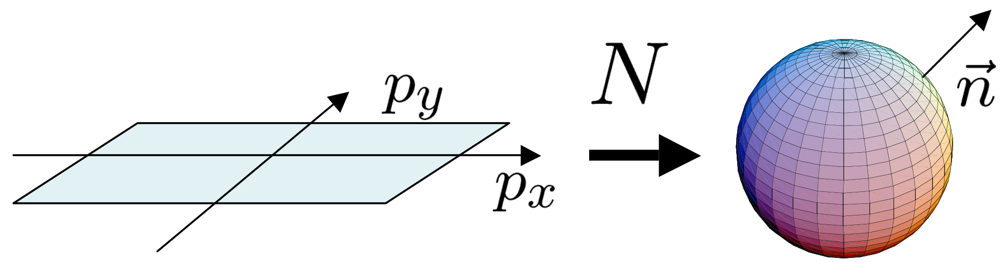
$$\begin{cases} n_x + in_y = 2v^* u \\ n_z = |v|^2 - |u|^2 \end{cases}$$

$$\vec{n} = \frac{1}{\sqrt{\left(\frac{p^2}{2m_a} - \mu\right)^2 + 4B^2\left(p_x^2 + p_y^2\right)}} \begin{pmatrix} 2gBp_x \\ -2gBp_y \\ \frac{p^2}{2m_a} - \mu \end{pmatrix}$$

Explicit calculations show that

N=*0* if $\mu < 0$

N=*1* if $\mu > 0$



$$N = \frac{1}{8\pi} \int d^2p \; \left[\vec{n} \cdot \partial_\alpha \vec{n} \times \partial_\beta \vec{n} \;\; \epsilon_{\alpha\beta}\right]$$

topological invariant

$p_x + i\,p_y$ superfluid in 2D

- Pfaffian (Moore-Read) state from FQH

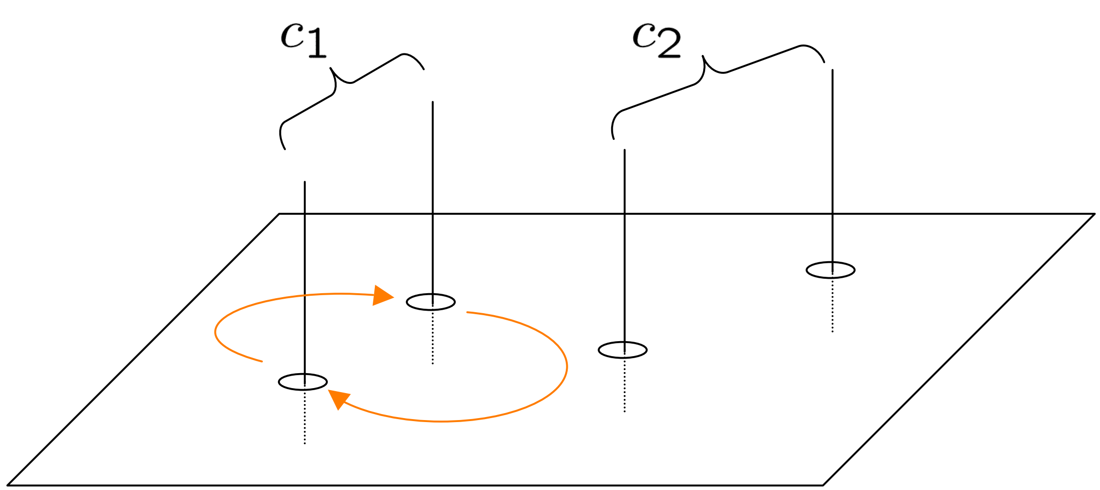
$$|p_x + ip_y{}_{BCS}\rangle = \prod_p \left[u_p + v_p a^\dagger_{-p} a^\dagger_p\right] |0\rangle$$

$$\Psi(z_1, z_2, \ldots, z_{2N}) = \sum_P (-1)^P \frac{1}{z_{P_1} - z_{P_2}} \frac{1}{z_{P_3} - z_{P_4}} \cdots \frac{1}{z_{P_{N-1}} - z_{P_N}}$$

- topological classification in terms of u_p and v_p
- gapped (N=1, BCS) ⇒ gapped (N=0, BEC) superfluid transition at $\mu=0$

Read and Green, PRB 61, 10267 (2000)

- vortex excitations with non-Abelian statistics *Ivanov, PRL (2001)*



one fermion (2 states –either empty or occupied fermion) per two vortices

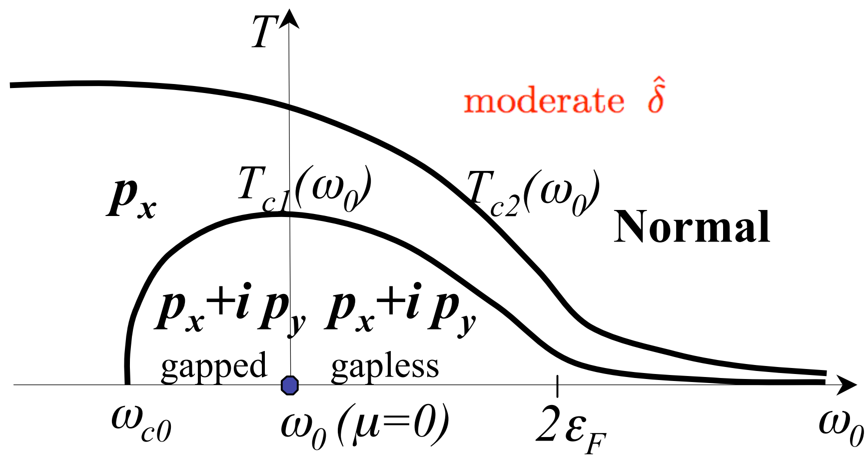
$2^{\frac{n}{2}}$ states per n vortices

- suggested to be used as qubits for quantum computers

Kitaev, Ann. Phys. 303, 2 (2003)

Summary of p-wave superfluidity

- mapped out T, $\omega_0 \propto B$, δ phase diagram for p-wave Feshbach resonant Fermi gas
  - *p_x and $p_x + i\, p_y$ superfluids*
  - *thermal, quantum and topological SF => SF transitions*



- quantitatively accurate description for small $\gamma=\Gamma/\varepsilon_F$ (low n)

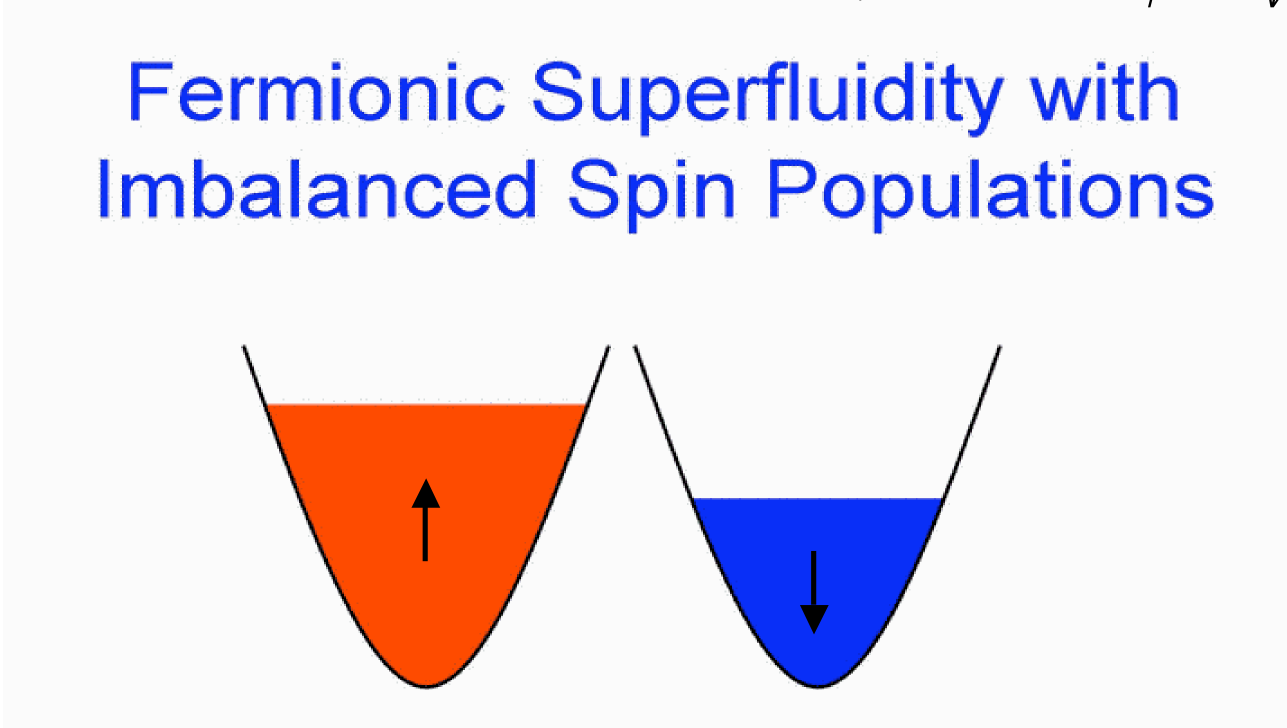
- realization of topological states, majorana zero modes, and non-Abelian statistics of Pfaffian (Moore-Read) state

- p-wave Feshbach molecules observed in K^{40}

- …BUT
  - ❖ *short (msec) molecular lifetime (see Levinson, et al, PRL 2007)*
  - ❖ *what about Li^6*
  - ❖ *need better quantitative understanding of stability*

Imbalanced ("magnetized") BEC-BCS

- motivation: *superconductivity in B field, quarks-gluon plasma,…*
- natural realization in cold atoms: $H_h = H - h(N_\uparrow - N_\downarrow)$



$$\mathcal{H} = \psi_\sigma^\dagger\left(\frac{p^2}{2m} - \mu_\sigma\right)\psi_\sigma + \lambda\psi_\uparrow^\dagger\psi_\downarrow^\dagger\psi_\downarrow\psi_\uparrow$$

$$n = \psi_\uparrow^\dagger\psi_\uparrow + \psi_\downarrow^\dagger\psi_\downarrow, \quad \Delta n = \psi_\uparrow^\dagger\psi_\uparrow - \psi_\downarrow^\dagger\psi_\downarrow$$

$b_q = B_Q\,\delta_{q,Q}$ **Mean-field theory** *(valid for* $\gamma \sim g^2/\epsilon_F^{1/2} \ll 1$ *)*

$$H_{\mu,h} = H - \mu N - h\Delta N$$

$N = N_{a\uparrow} + N_{a\downarrow} + 2\,N_b$

$\Delta N = N_{a\uparrow} - N_{a\downarrow}$

- ***ground state:*** $|gs\rangle = \Pi'_{\mathbf{k}}\left(u_{\mathbf{k},\mathbf{Q}} + v_{\mathbf{k},\mathbf{Q}} a^\dagger_{-\mathbf{k}+\mathbf{Q}/2,\downarrow} a^\dagger_{\mathbf{k}+\mathbf{Q}/2,\uparrow}\right)|0\rangle$

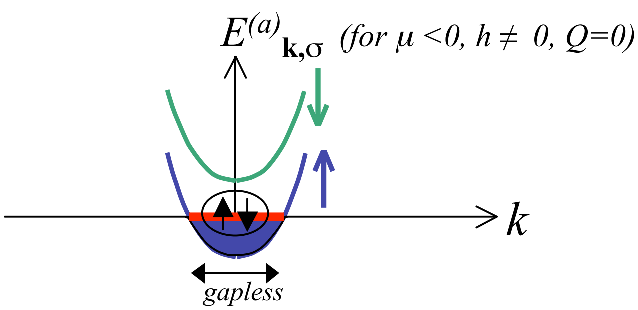
- ***ground state energy:*** $E_{gs} = \left(\frac{Q^2}{4m} + \delta - 2\mu\right)B_Q^2 - \sum_{\mathbf{k}}(E_k - \varepsilon_k) + \sum_{\mathbf{k}}\left[E_{\mathbf{k},\uparrow}\theta(-E_{\mathbf{k},\uparrow}) + E_{\mathbf{k},\downarrow}\theta(-E_{\mathbf{k},\downarrow})\right]$

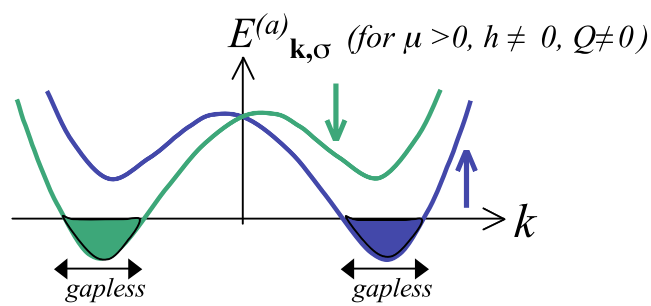
$$E_k = (\varepsilon_k^2 + g^2 B_Q^2)^{1/2}, \quad \varepsilon_k = \frac{k^2}{2m} - \mu + \frac{Q^2}{8m}$$

- ***excitation spectrum:*** $H_{ex} = \sum'_{\mathbf{k},\sigma} E^{(a)}_{\mathbf{k},\sigma}\alpha^\dagger_{\mathbf{k},\sigma}\alpha_{\mathbf{k},\sigma} + \sum'_{\mathbf{k},\sigma} E^{(b)}_{\mathbf{k},\sigma}\beta^\dagger_{\mathbf{k}}\beta_{\mathbf{k}}$

$E^{(a)}_{\mathbf{k},\sigma} = E_k \mp (h + \mathbf{k}\cdot\mathbf{Q}/2m)$, *(gapped and gapless k's)*

$E^{(b)}_k = \sqrt{\epsilon_k^2 + V_0\epsilon_k}$ (for $Q = 0$) *(gapless k's collective; also phonons $Q \neq 0$))*





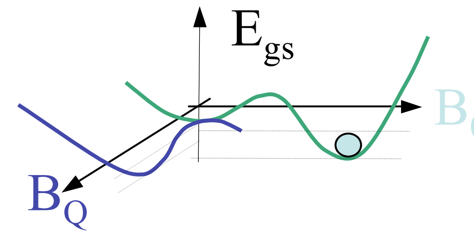
- ***determine*** B_Q, N_b, $N_{a\uparrow}$, $N_{a\downarrow}$ (ΔN_a), Q ***by:***

energy minimization ⇨ $\frac{\partial E_{gs}}{\partial B_Q} = 0$ *(gap equation),* $\frac{\partial E_{gs}}{\partial \mathbf{Q}} = 0$ *(*$P_{total}=0$*)*

μ, h fixed

BSC and crossover regimes $(\delta > 0)$

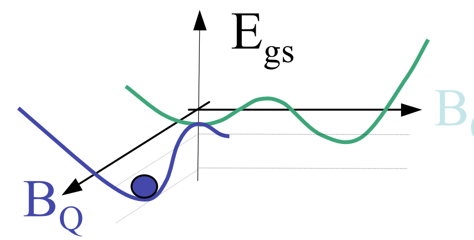
- ***BCS SF*** $B_0 \neq 0$, $B_Q = 0$, $\Delta N = 0$: $0 < h < h_c = \frac{1}{\sqrt{2}}\, g\, B_0(\mu,\delta)$



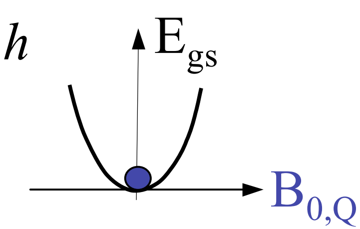
- ***FFLO*** $B_0 = 0$, $B_Q \neq 0$, $\Delta N \neq 0$: $h_c < h < h_{FFLO}(\delta) \xrightarrow{\delta \gg 2\epsilon_F} 1.1\, h_c$ *(Fulde-Ferrell)*

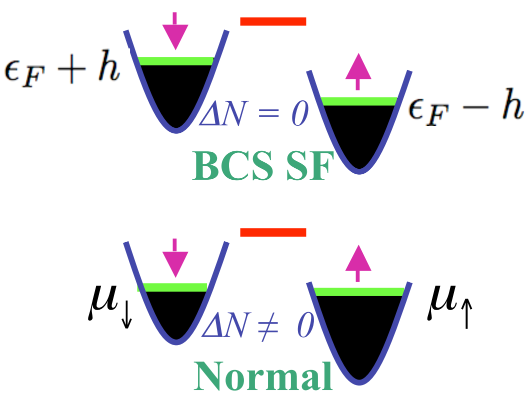
  $\xrightarrow{\delta \approx 2\epsilon_F} h_c$

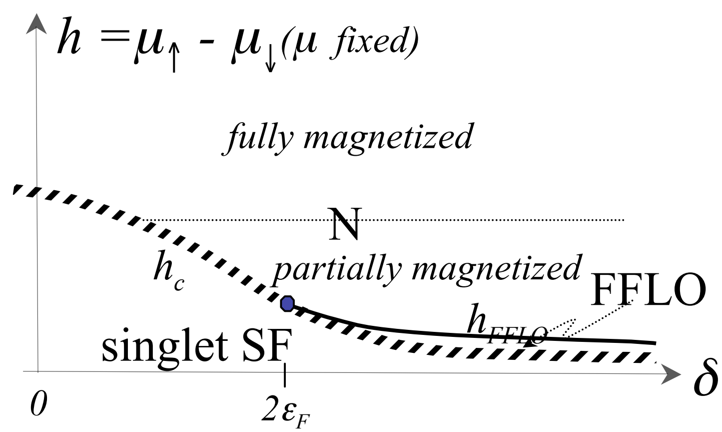
  supersolid: *broken rotational and translational symmetry*



- ***Normal*** $B_0 = 0$, $B_Q = 0$, $\Delta N \neq 0$ *(Pauli "paramagnet")*: $h_{FFLO}(\delta) < h$

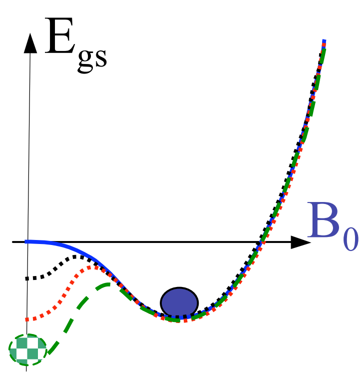






Phase separation ($\delta > 0$)

N, h fixed

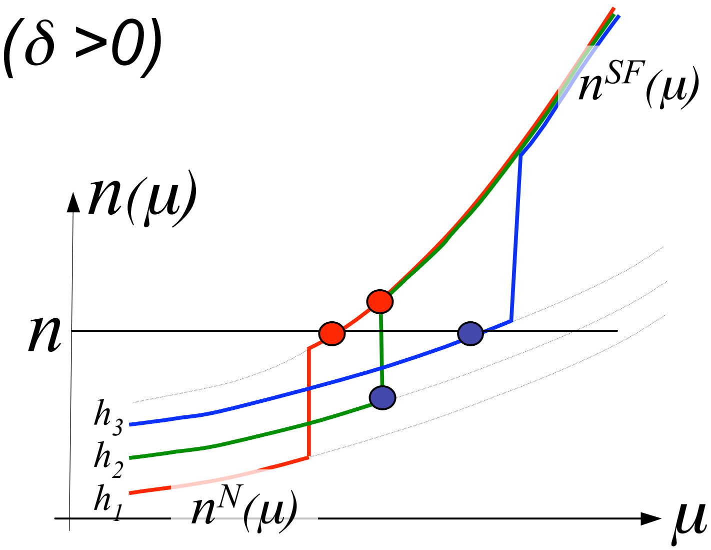


h_c

$h_1 < h_{c1}(\delta) < h_2 < h_{c2}(\delta) < h_3$

SF N

$n = (1-x)n^{SF} + x\, n^{N}$



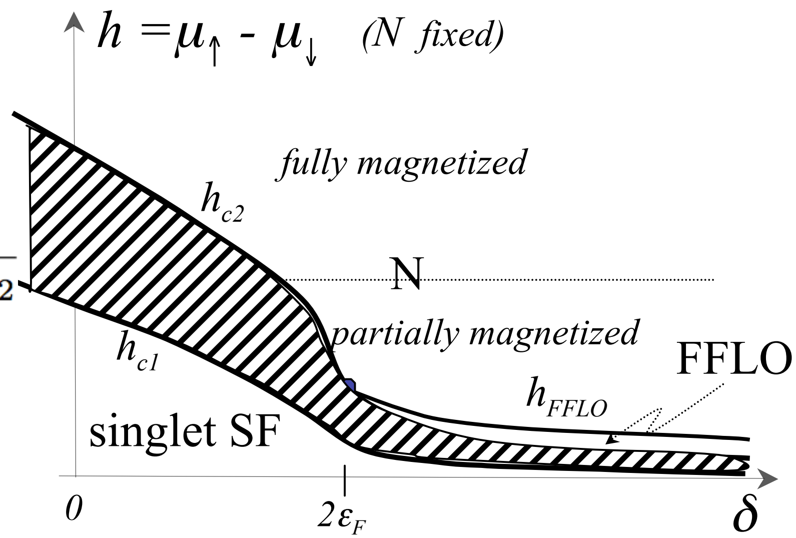
phase separation $x(h, \delta) = (n^{SF} - n)/(n^{SF} - n^{N})$

$$h_{c2}(\delta) = h_c(\mu^{(FFLO)}(N, \delta), \delta) \qquad h_{c1}(\delta) = h_c(\mu^{(SF)}(N, \delta), \delta)$$

$0 < \delta \ll 2\epsilon_F$

$$h_{c2}(\delta) \approx 2^{3/2}\epsilon_F - \delta/2$$

$$h_{c1}(\delta) \approx \epsilon_F \sqrt{\frac{\gamma}{3}} \sqrt{1 - (\delta/2\epsilon_F)^{3/2}}$$



$2\epsilon_F \ll \delta$

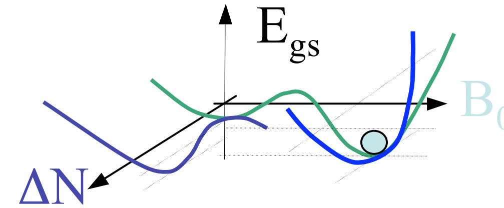
$$h_{FFLO}(\delta) \longrightarrow 1.1 h_c$$

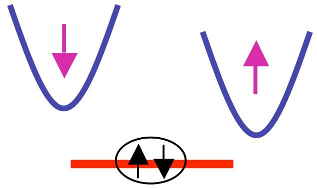
$$\left. \begin{array}{l} h_{c2}(\delta) \\ h_{c1}(\delta) \end{array} \right\} \rightarrow \frac{1}{\sqrt{2}} \Delta_{BCS}$$

BEC regime $(\delta < 0)$

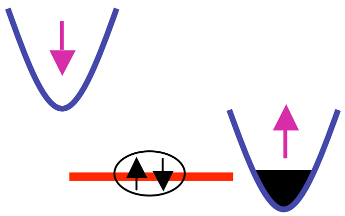
***N*, h fixed**

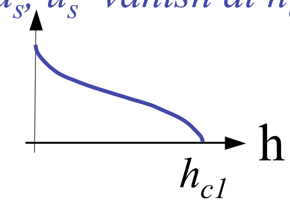
- ***BEC SF*** $B_0 \neq 0$, $B_Q = 0$, $\Delta N = 0$: $0 < h < h_m(\delta) \approx -\delta/2$

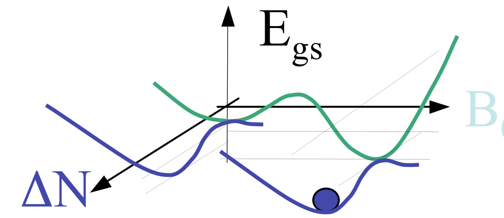




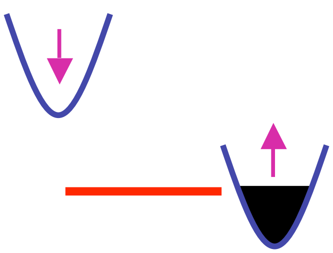
- ***SF_M*** $B_0 \neq 0$, $B_Q = 0$, $\Delta N \neq 0$: $h_m < h < h_{c1}(\delta) \approx -0.65\delta$

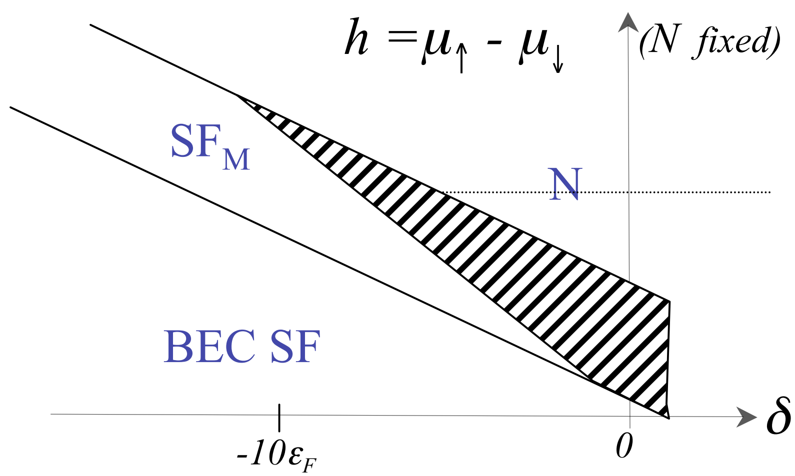


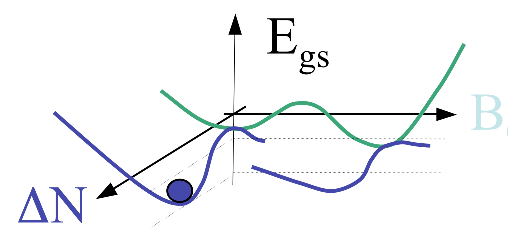




- ***Normal*** $B_0 = 0$, $B_Q = 0$, $\Delta N \neq 0$ *(Pauli "paramagnet")*: $h > h_{c2}(\delta) \approx 2^{3/2}\epsilon_F - \delta/2$

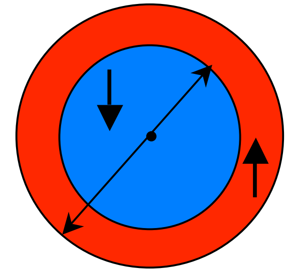




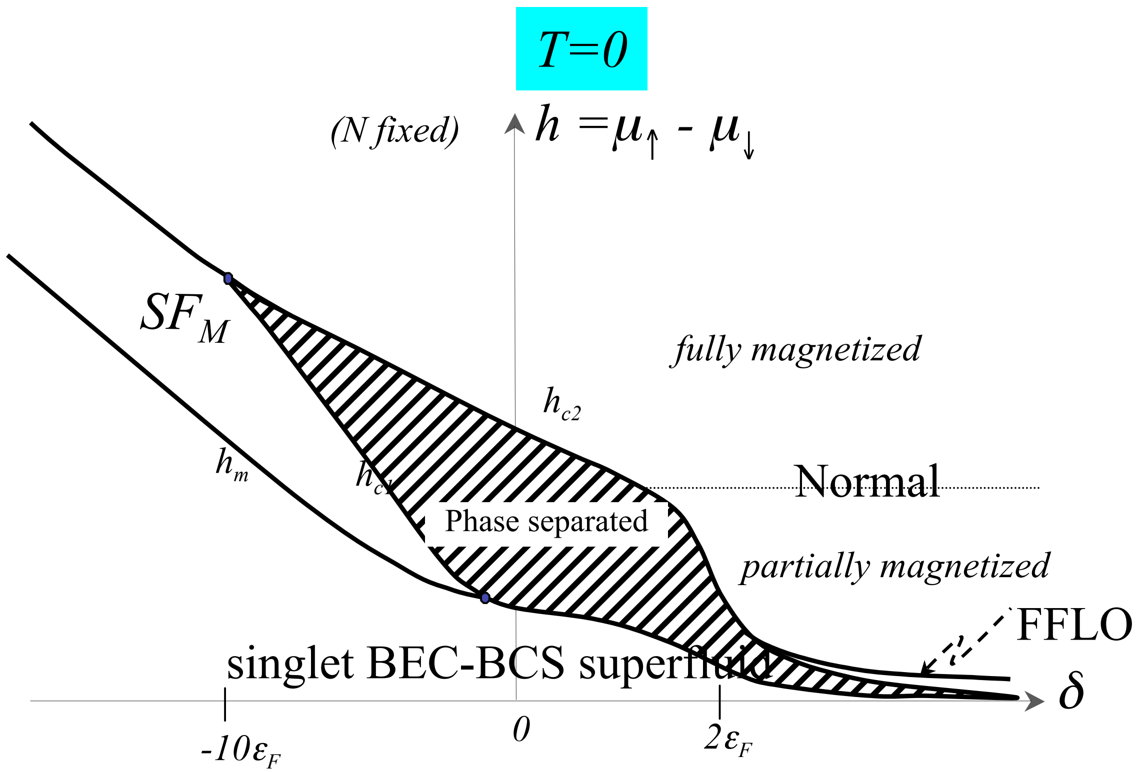


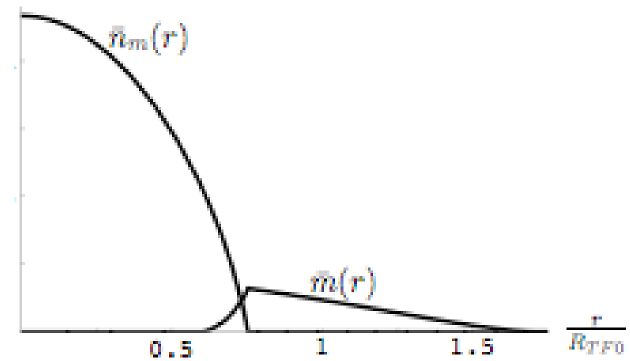
Imbalanced BEC-BCS

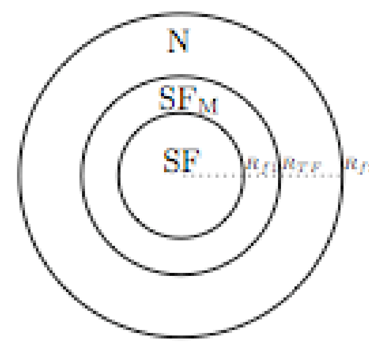
Sheehy, L.R. '05

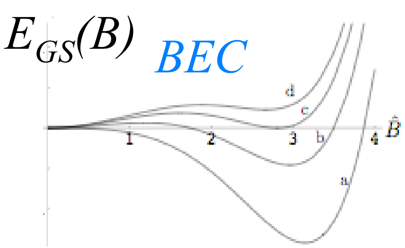


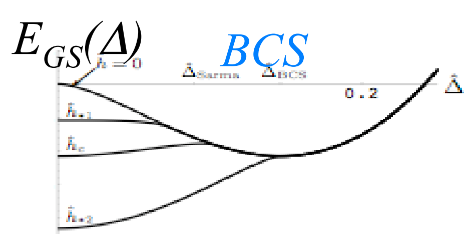
- 1st order transitions and phase separation







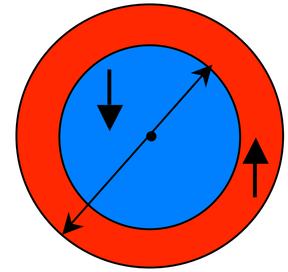


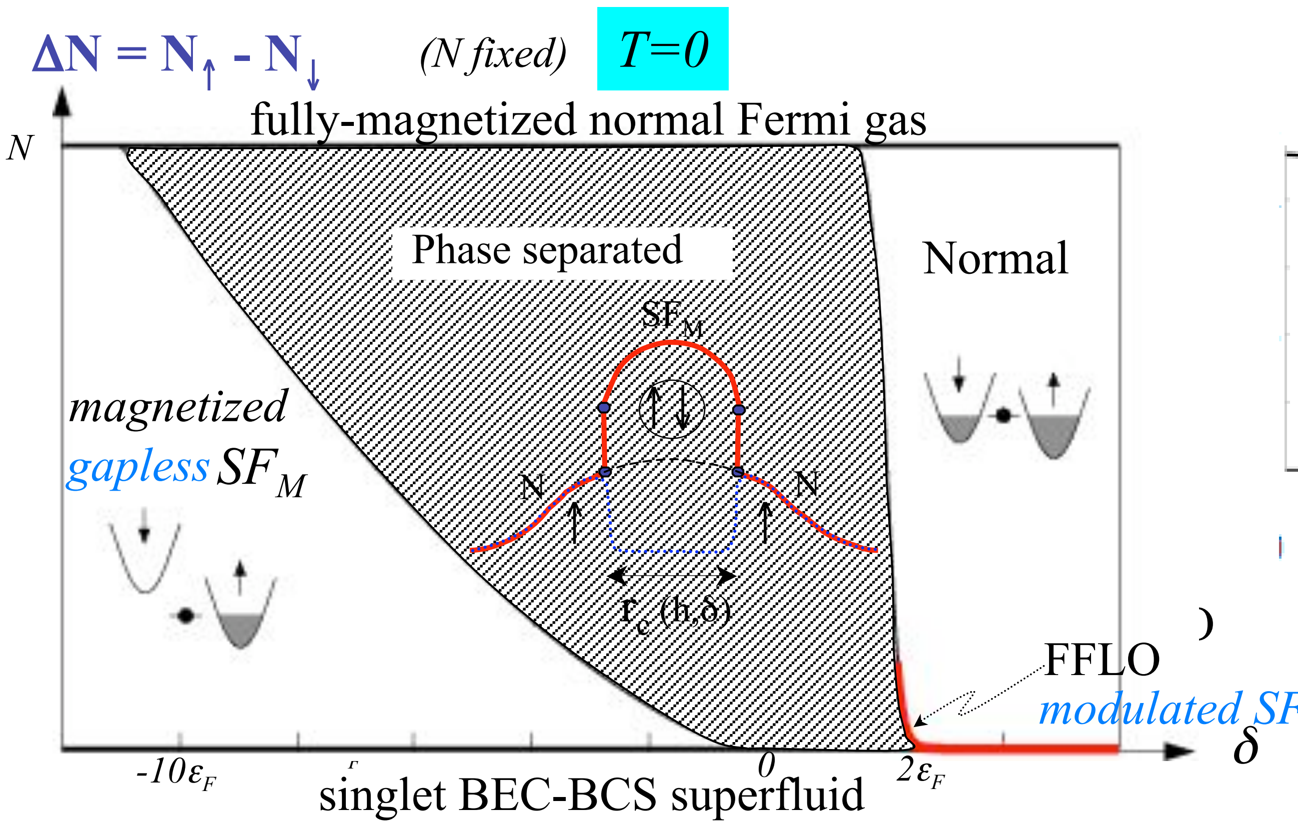


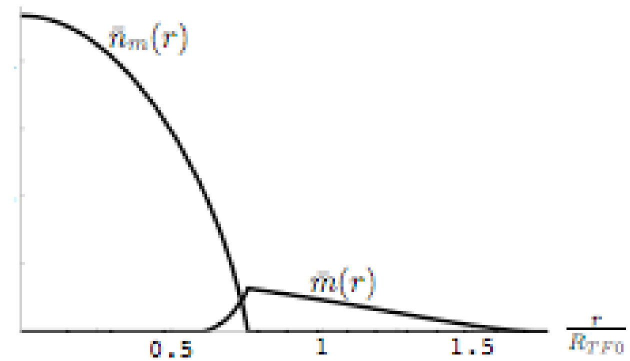
Imbalanced BEC-BCS

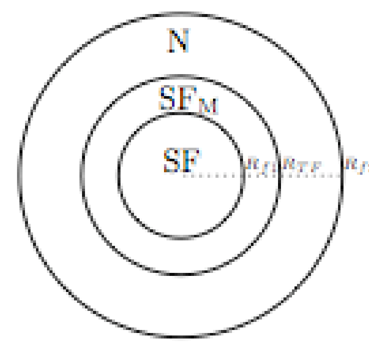
Sheehy, L.R. '05

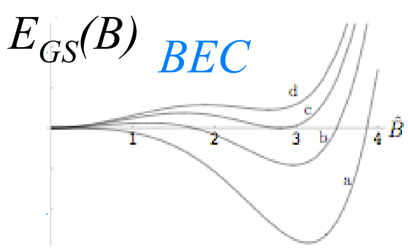
- 1st order transitions and phase separation

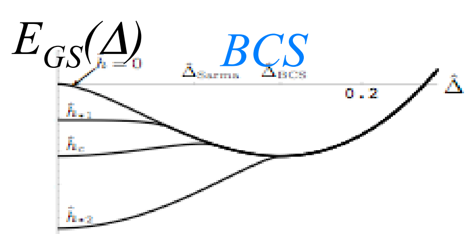






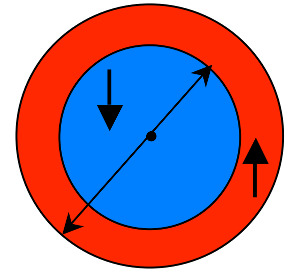






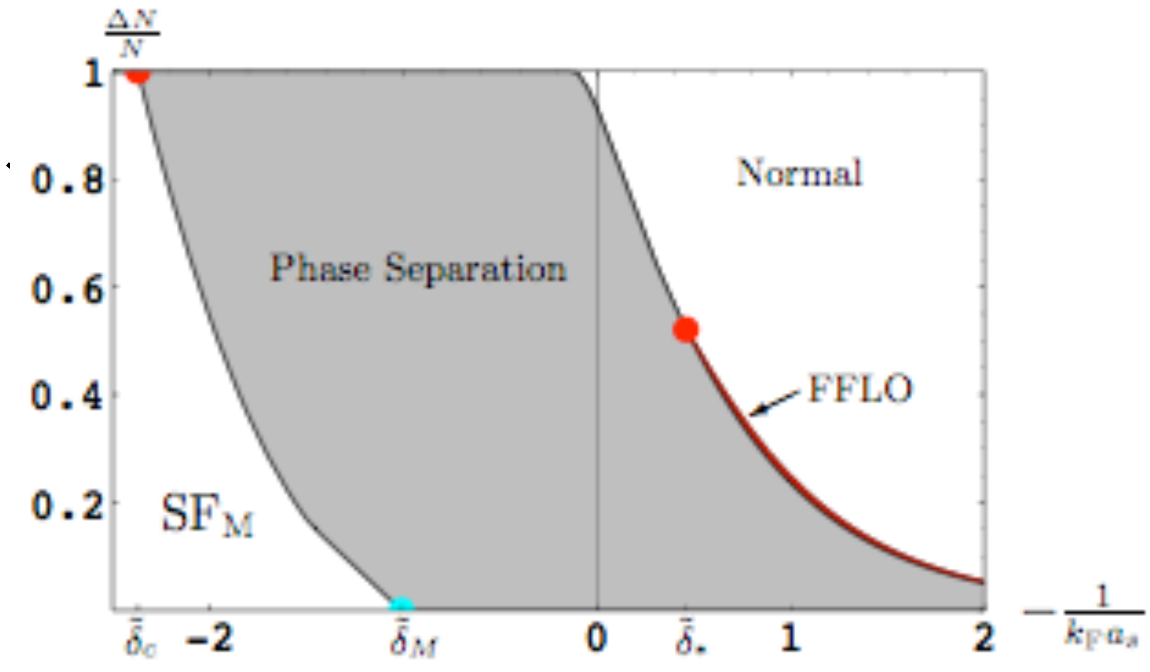
Imbalanced BEC-BCS

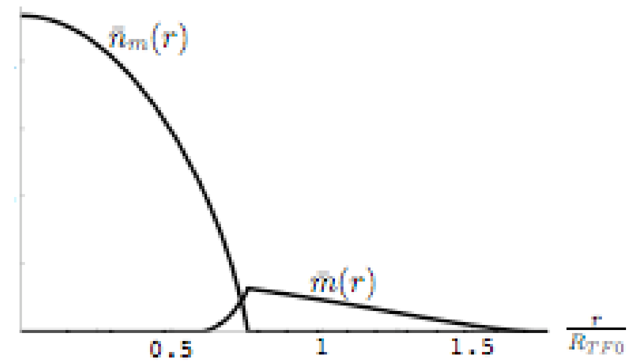
Sheehy, L.R. '05

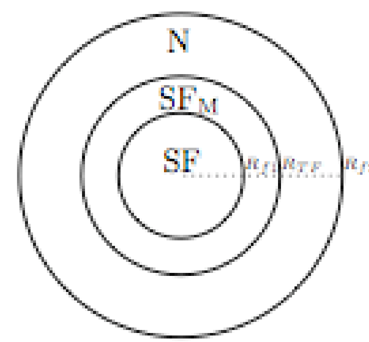


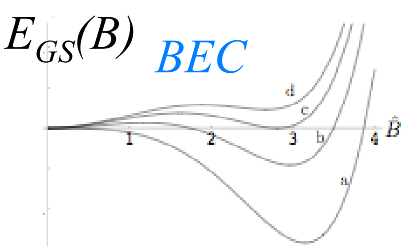
- 1st order transitions and phase separation

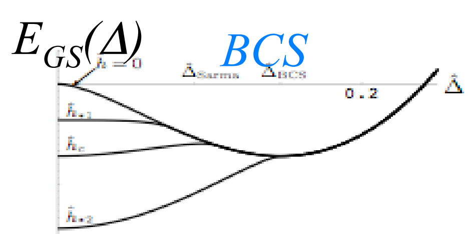
$T=0$







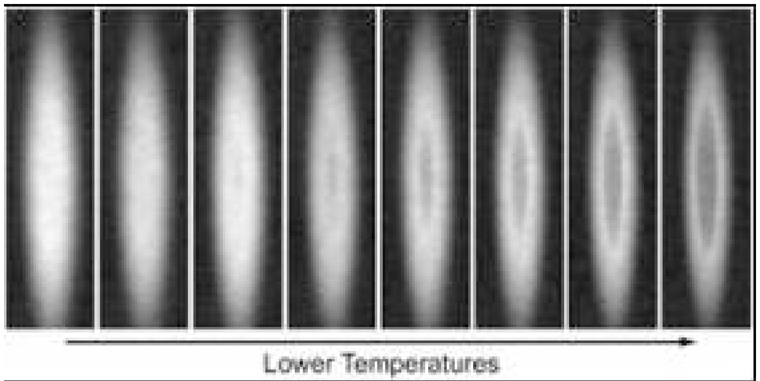


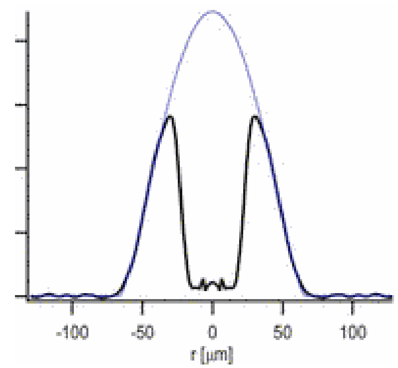


Imbalanced BEC-BCS experiments

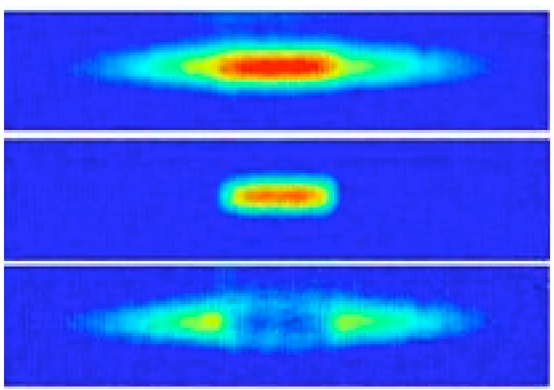
(2006)

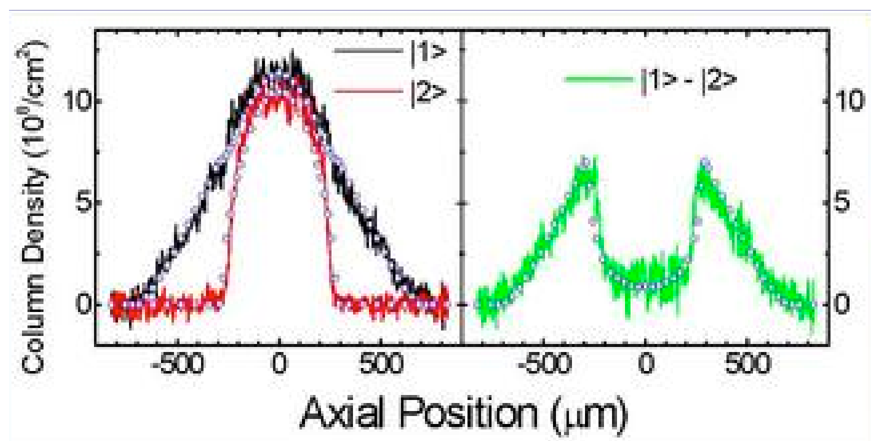
• Ketterle's experiments (vortices, phase separation)





• Hulet's experiments (phase separation, surface tension)





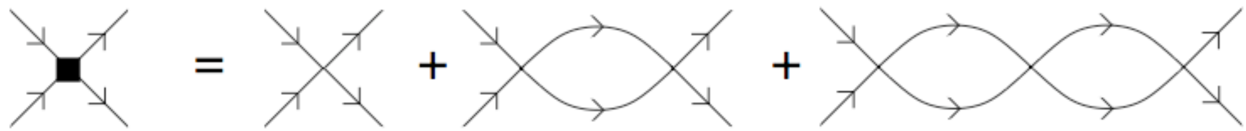
$\gamma >> 1$

Broad resonance scattering

$$\mathcal{H}_{2ch} \longrightarrow \mathcal{H}_{1ch} = \psi_\sigma^\dagger\left(\frac{p^2}{2m} - \mu_\sigma\right)\psi_\sigma + \lambda\psi_\uparrow^\dagger\psi_\downarrow^\dagger\psi_\downarrow\psi_\uparrow$$

• scattering T-matrix relates λ to a:

$T_{k\,k'} =$

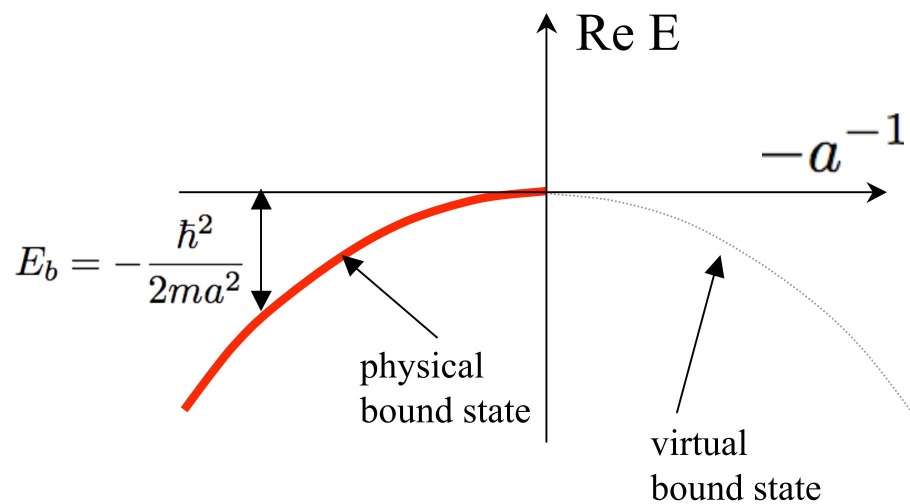


$$+\ldots = \frac{\lambda}{1-\lambda\Pi}$$

$$= -\frac{4\pi\hbar^2}{m} f_{kk'} \approx \frac{4\pi\hbar^2}{m}\frac{1}{a^{-1}+ik}$$

$$\longrightarrow \boxed{a = \frac{m}{4\pi\hbar^2}\frac{\lambda}{1+\lambda/\lambda_c}}$$

$(\lambda_c = \pi d/m)$



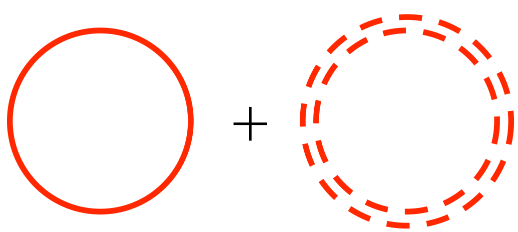
$\gamma >> 1$ **Broad resonance superfluidity: Large N**

- no small parameter for $k_F a \sim n^{1/3} a >> 1$ → *introduce 1/N*

$$\mathcal{H}_{1ch} \xrightarrow{Sp(2N)} \mathcal{H}_N = \psi^\dagger_{\sigma\alpha}\left(\frac{p^2}{2m} - \mu_\sigma\right)\psi_{\sigma\alpha} + \frac{\lambda}{N}\psi^\dagger_{\uparrow\alpha}\psi^\dagger_{\downarrow\alpha}\psi_{\downarrow\beta}\psi_{\uparrow\beta}$$

$$S[\phi] = -\frac{N}{\lambda}\int_0^\beta d\tau d^3r|\phi|^2 - N\mathrm{Tr}\log\left[-G_\phi^{-1}\right] \qquad G_\phi^{-1} = \begin{pmatrix} -\partial_\tau + \frac{\nabla^2}{2m} + \mu_\uparrow & \phi_x \\ \phi_x^* & -\partial_\tau - \frac{\nabla^2}{2m} - \mu_\downarrow \end{pmatrix}$$

$$\begin{aligned} f &= -\frac{1}{\beta V}\log\int D\phi e^{-S[\phi]}, \\ &= N f^{(0)} + f^{(1/N)} + \ldots \end{aligned}$$



MFT

$$f^{(0)} = -\frac{|\Delta|^2}{\lambda} - \int_k (E_k - \xi_k) - \sum_{\sigma=\pm}\int_k \log\left[1 + e^{-\beta(E_k + \sigma h)}\right]$$

Veillette, Sheehy, LR
Nikolic, Sachdev
also Nishida, Son
ε-expansion

$\gamma >> 1$, $k_F a \to \infty$

Universality at unitary point

T.L. Ho '04

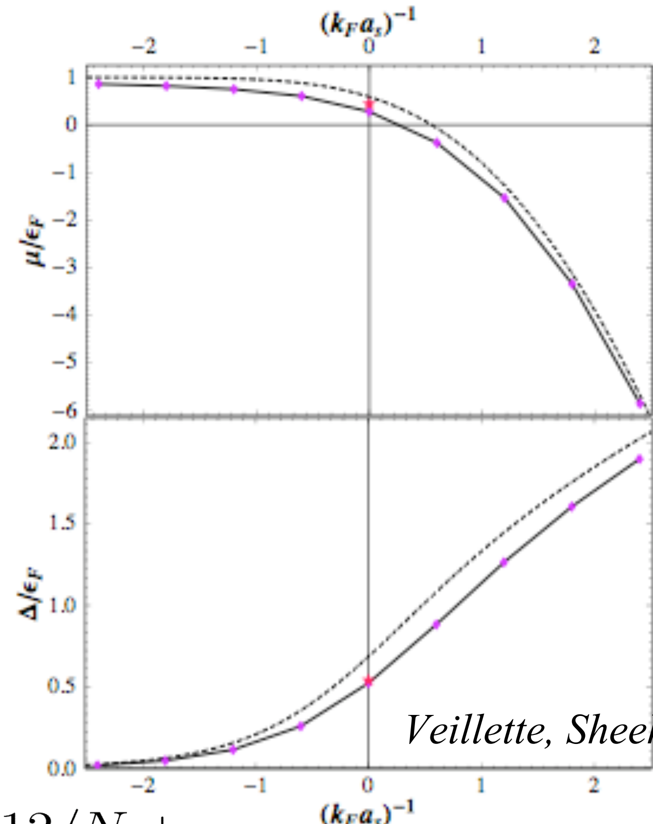
- $f_k = -1/(a^{-1} + i\,k) \to i/k$, k_F is the only scale

check in $N \to \infty$ (BCS) limit:

$$\frac{m}{2\pi\hbar^2 a} \to 0 = \int_k \left(\frac{1}{E_k} - \frac{1}{\epsilon_k}\right)$$

$$f(T,n) = n\epsilon_F \hat{f}(k_B T/\epsilon_F)$$

$$\begin{aligned} \epsilon &= \xi \frac{3}{5}\epsilon_F \\ \mu &= \xi \epsilon_F \\ \Delta &= \alpha \epsilon_F \\ \Delta_{exc} &= \alpha_{exc}\epsilon_F \\ k_B T_c &= \gamma \epsilon_F \\ B &= \xi \frac{2}{3} n \epsilon_F \end{aligned}$$



Veillette, Sheehy, LR

Result from 1/N $\xi = 0.5906 - 0.312/N + \cdots$

Exp with ^{40}K $\xi = 0.46^{+0.05}_{-0.12}$

Questions of current interest

- What are the big fundamental questions?

- Specific questions of current experimental interest:
  - *Unitary Fermi gas (universality)*
  - *Resonant Bose gas (beyond Beliaev)*
  - *Stability to 3-body collisions and other inelastic processes*
  - *Cooling and thermalization*
  - *Experimental probes (development and understanding)*
  - *Phases realizations (e.g., FFLO, p-wave SF, magnetism,…)*
  - *Nonequilibrium quantum dynamics*
  - …

On the horizon

- p-wave superfluidity?
- degenerate molecular gases?
- local many-body lattice models
- multi-site many-body lattice models ⇒ exotic models?
- quantum Hall regime?
- …

…but not before technical hurdles are overcome:

- cooling
- off-site interactions
- stability to inelastic processes near Feshbach resonances
- much larger clouds
- flat traps
- better and wider range of experimental probes
- …