



# The Design, Discovery, Growth and Physical Properties of Novel Intermetallic Compounds

## ***SUPERCONDUCTIVITY***

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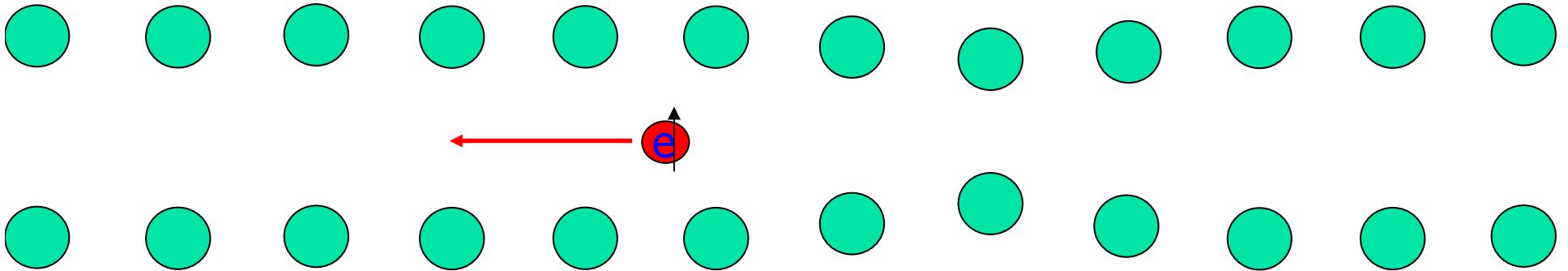
# Physics among light-weights: Superconductivity in $\text{MgB}_2$ and other borides



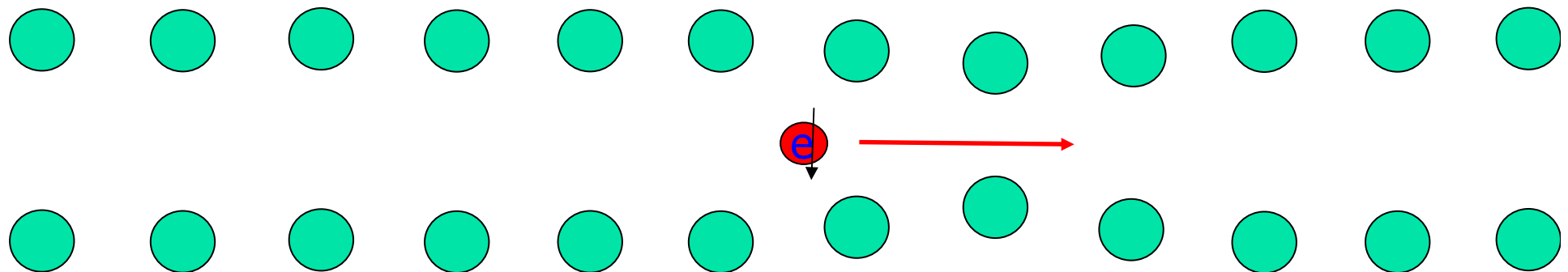
Or, superconductivity from  
 $\text{RNi}_2\text{B}_2\text{C}$  to  $\text{MgB}_2$



The key idea in BCS superconductivity is the interaction between electrons and phonons



An electron moving through the lattice leaves a distortion in its wake. The net, positive charge on this distortion attracts a second electron, briefly forming a Cooper pair.



*This interaction clearly depends on electrons, phonons, and their coupling....*



## BCS Superconductivity: 3 parameters (+ some prejudices)

$$k_B T_C = 1.13 \hbar \omega_D e^{-1/VN(E_F)}$$

- 1) Electrons pair via exchange of a phonon (lattice vibration).  $\omega_D$  is the characteristic frequency of the lattice. For a simple mass/spring system  $\omega = \sqrt{k/m} \rightarrow$  lighter  $m \rightarrow$  higher  $T_C$

Prejudice: higher  $T_C$  values can be found in compounds with light mass elements

- 2)  $V \equiv$  electron-phonon coupling strength.

Higher  $V \rightarrow$  Higher  $T_C$ , BUT too large of a  $V$  leads to distortions (e.g. CDW)

Prejudice: higher  $T_C$  values can be found close to structural phase transitions

- 3)  $N(E_F) =$  Density of states at the Fermi surface. This is basically a caliper of the number of electrons that can participate in the superconducting ground state. (Note that d-levels generally have higher  $N(E_F)$  than p-levels or s-levels.

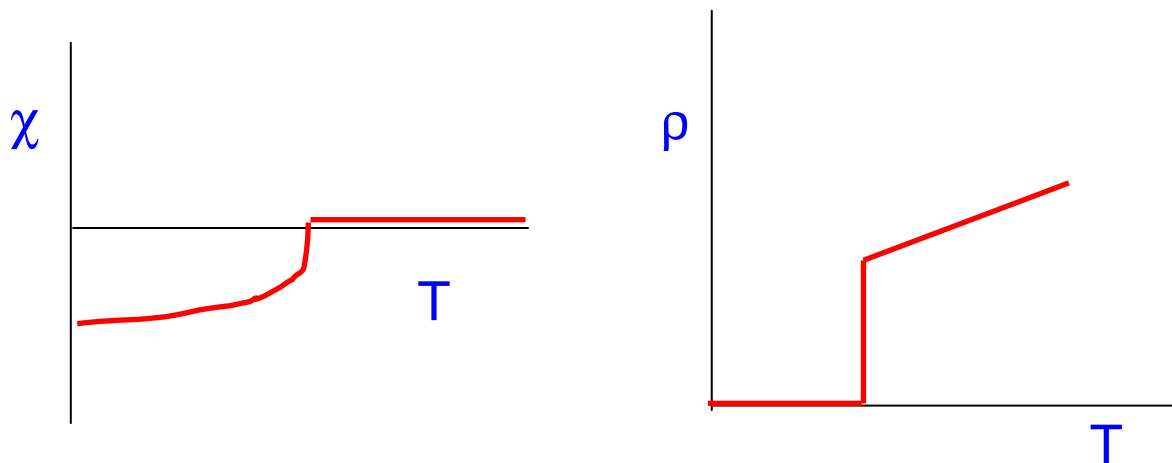
Prejudice: For high  $T_C$  values there has to be a transition metal element (with its d-levels) so as to pump up the  $N(E_F)$



## *Superconductivity: a recent history.*

In the 1950's, 1960's and even 1970's there was a large-scale hunt for higher  $T_c$  superconductors (intermetallics since they were *the only* place to expect such phenomena). In the U.S., these took place in Bell Labs, GE, UCSD, etc.

Binary compounds were made (often by arc melting) and tested for superconductivity, usually by quick  $\rho(T)$  or  $\chi(T)$  measurements at low-T.



These were exciting times! You could make a sample in the morning and discover a new superconductor by afternoon coffee.



For decades  $T_c$  was stuck at  $\sim 20$  K. (Theories were written why B.C.S. could not support  $T_c > 30$  K while others predicted  $T_c$  up to  $\sim 100$  K.)

Then came the copper oxides.... $T_c > 120$  K

Pleasure

*Pain*

Complex materials problems

?Mechanisms?

*Hard to work with*

But even during these exciting times, intermetallics remained interesting....



The past thirteen years have been an exciting time for boride based superconductors..

1994:  $\text{RNi}_2\text{B}_2\text{C}$  and  $\text{YPd}_2\text{B}_2\text{C}$

2001:  $\text{MgB}_2$

*These two classes of compounds have pushed our understanding of superconductivity in intermetallic compounds, extending the range over which superconductivity is known to exist: to higher temperatures and new extremes of interaction with local moments as well as with the underlying lattice.*



Based the three prejudices, what you want is the following:

Light atoms / high characteristic frequencies

Strong electron phonon coupling, but no structural transition

Good  $N(E_F)$ , probably coming (in part) from transition metal

$RNi_2B_2C$  series as well as  $YPd_2B_2C$  are perfect!!

$LuNi_2B_2C$ :  $T_C \sim 17$  K

$N(E_F)$  large with significant Ni contribution

Very close to a structural phase transition

Large Debye temperature (high characteristic frequency)

$YPd_2B_2C$ :  $T_C \sim 23$  K

Metastable....If annealed loses structure---basically “just beyond” a structural phase transition, but, like diamond, can be trapped into structurally metastable state.

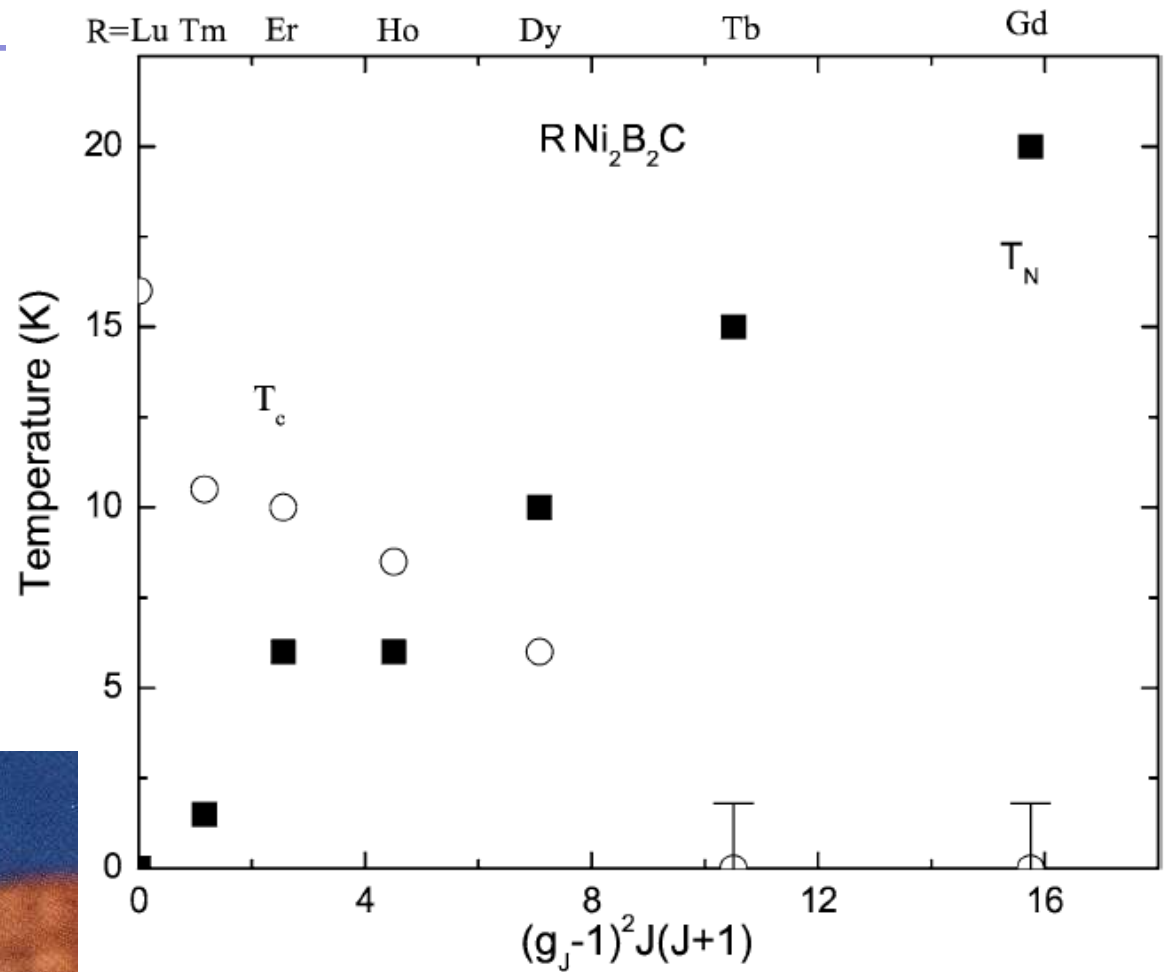




# RNi<sub>2</sub>B<sub>2</sub>C family

R = Gd – Lu, Y

*P. C. Canfield et al.*  
*Physics Today Oct. 1998*



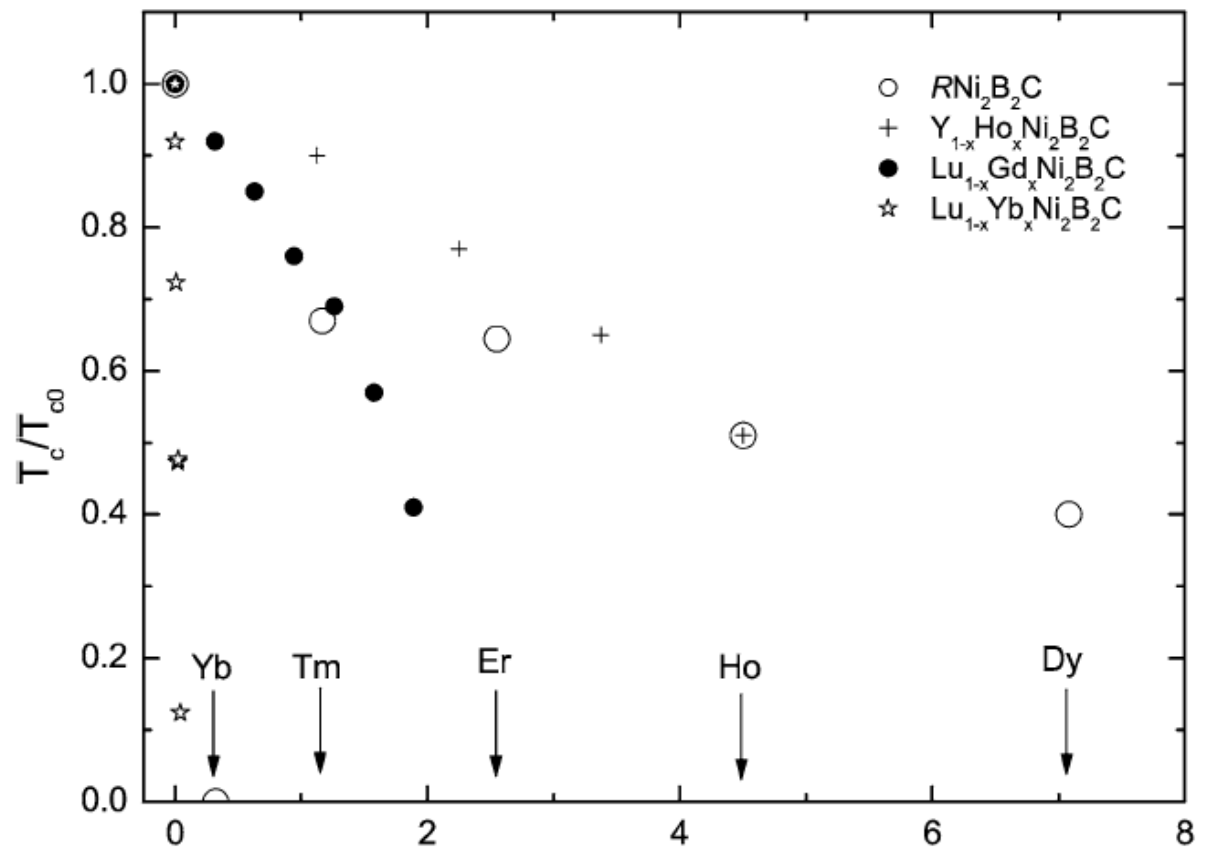
Superconducting for R = Dy, Ho, Er, Tm, Lu, Y  
 $T_c$  values ranging from 17 K – 6 K

Magnetic order for R = Gd, Tb, Dy, Ho, Er, Tm  
 $T_N$  values ranging from 20 K to 1.5 K



The suppression of  $T_c$  with increasingly magnetic rare earth is due to the fact that the Cooper pairs are spin up, spin down pairs. The interaction between the electrons and local moments involves a spin flip, thereby acting as a pair breaker.

This pair breaking is one of the reasons why magnetic superconductors are so interesting: they allow the study of competing effects.



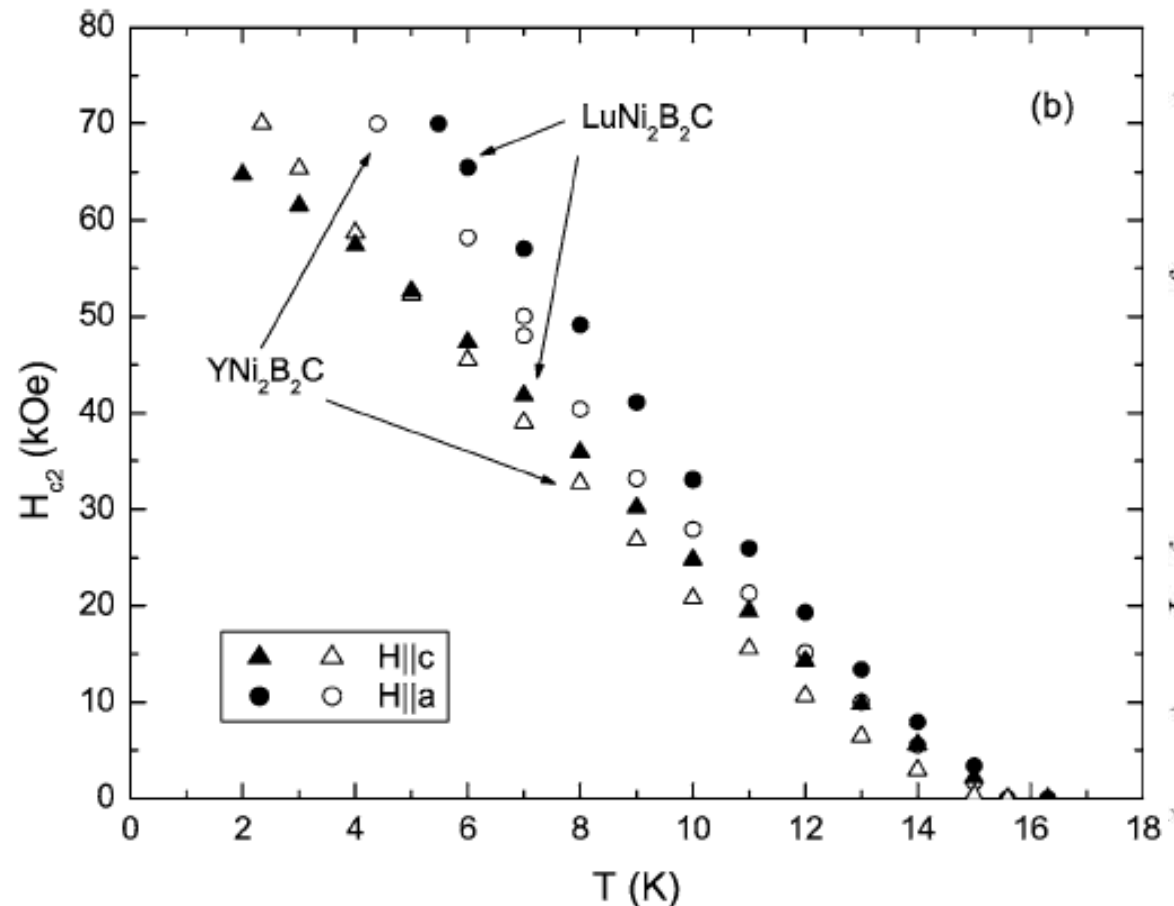


This interplay between superconductivity and local moment magnetism manifests itself most clearly as anomalies in  $T_c$  and  $H_{c2}(T)$ .

*Let's start by looking at  $H_{c2}(T)$*

Since Lu and Y are non-moment bearing rare earths, both  $\text{LuNi}_2\text{B}_2\text{C}$  and  $\text{YNi}_2\text{B}_2\text{C}$  serve as examples of what superconductivity in these compounds looks like when there is no 4f-based, local moment magnetism present.

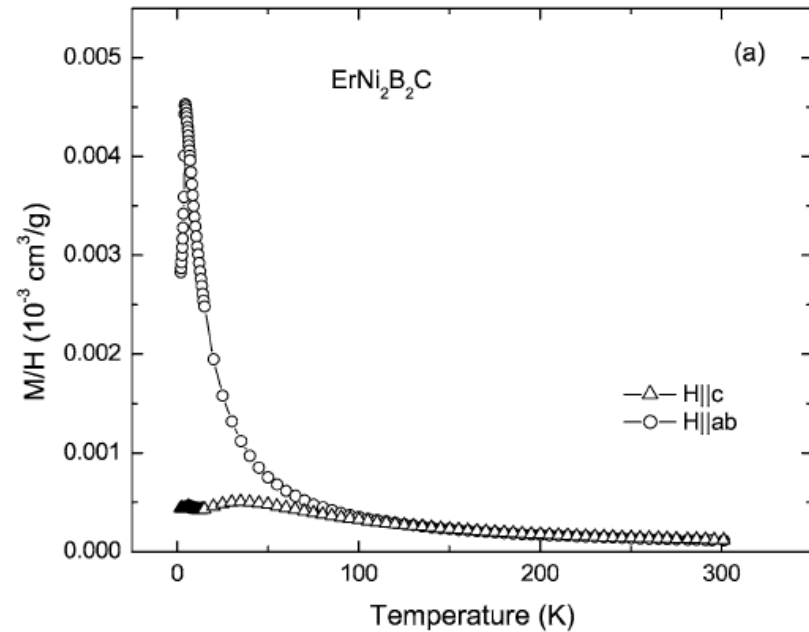
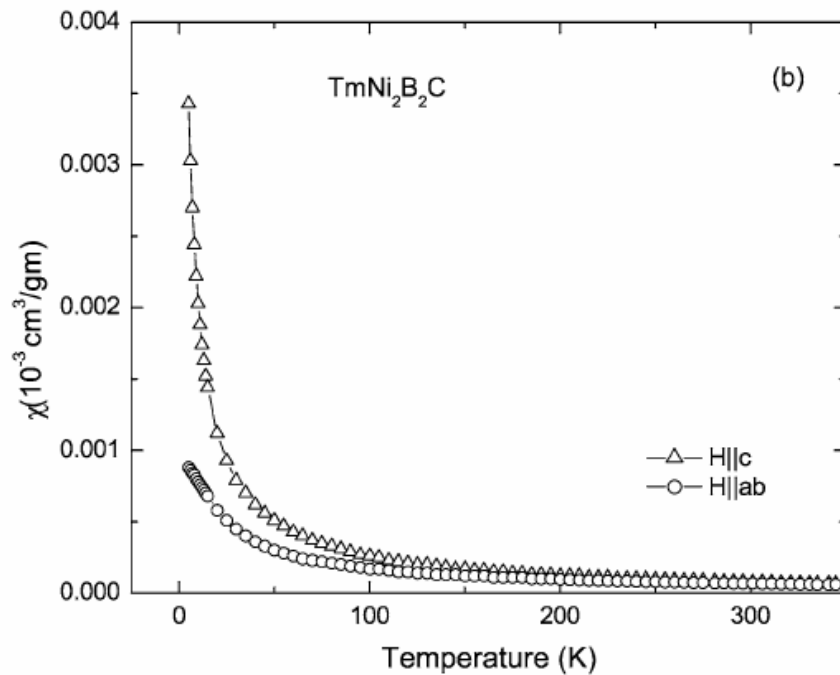
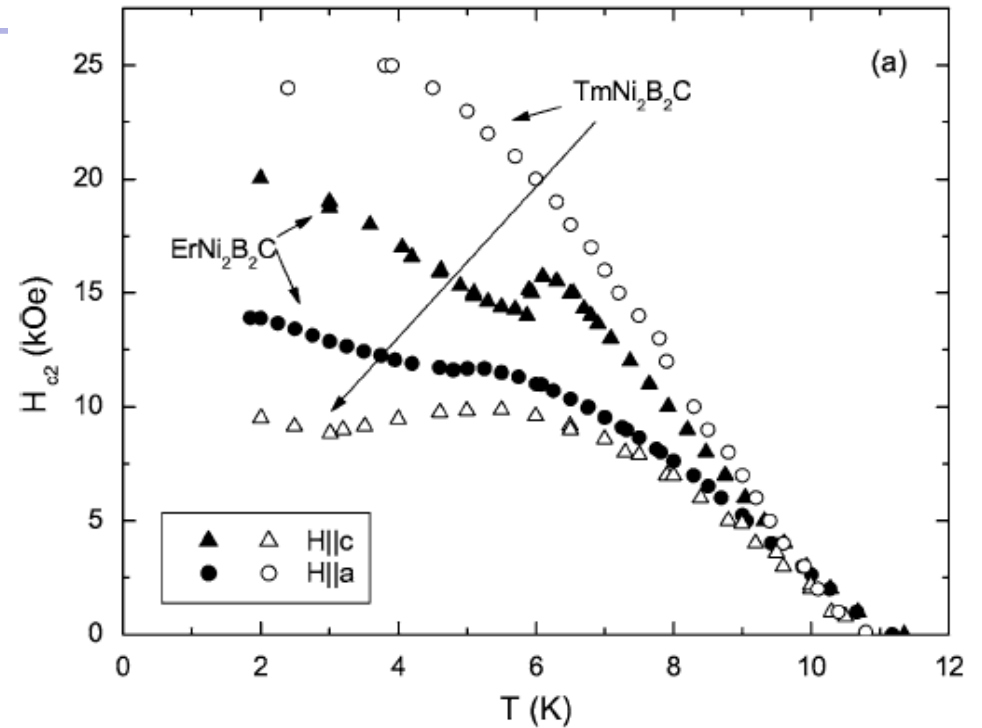
Both manifest relatively high  $T_c$  values and minor  $H_{c2}(T)$  anisotropy.



One effect of local moment magnetism is immediately clear from the correlations between the anisotropy of  $\chi$  and  $H_{c2}(T)$ : larger  $\chi$  leads to smaller  $H_{c2}(T)$ . This is due to the internal field suppressing superconductivity.

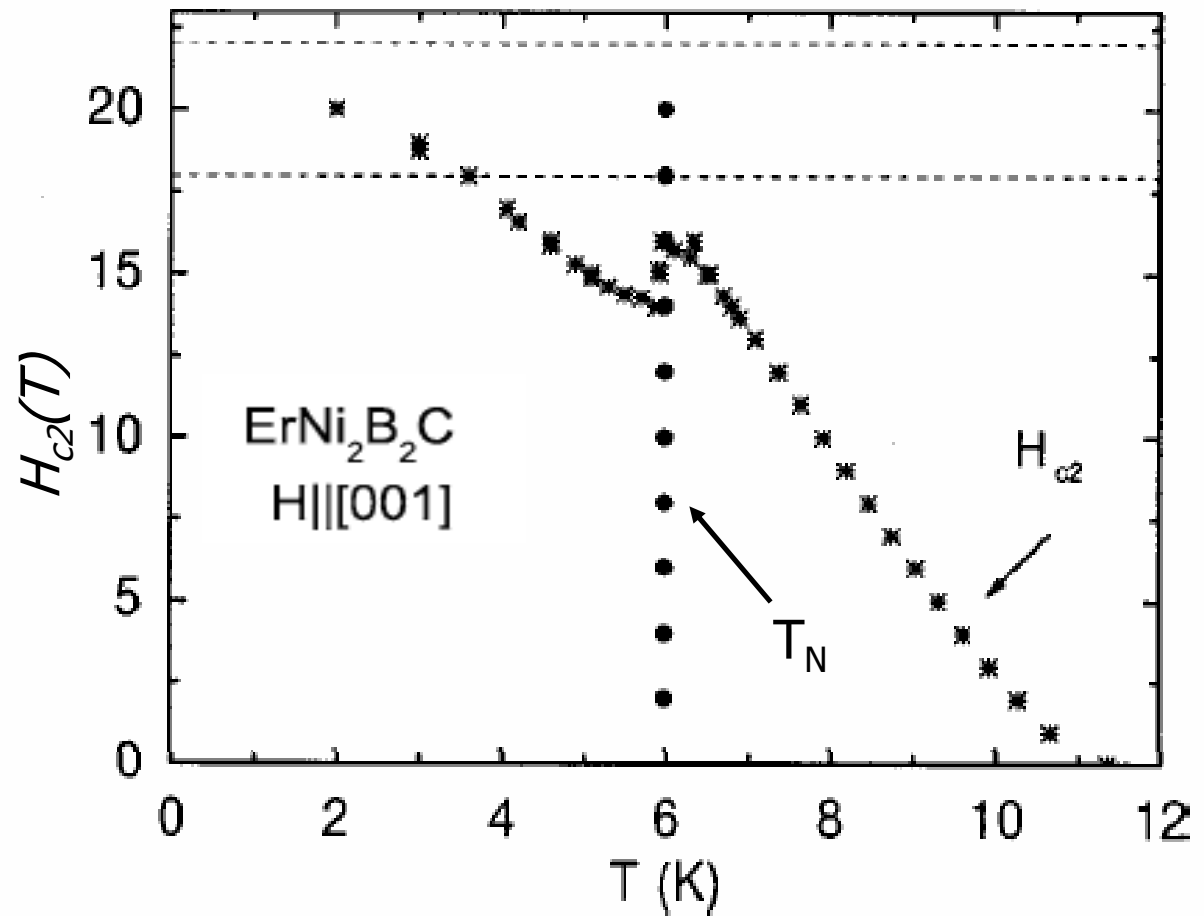
*P. C. Canfield et al. Physics Today Oct. 1998*

*S.L. Bud'ko, P.C. Canfield / C. R. Physique 7 (2006) 56–67*





*In addition:  
sharp features  
in  $H_{c2}(T)$  are  
due to the  
local moment  
ordering:*

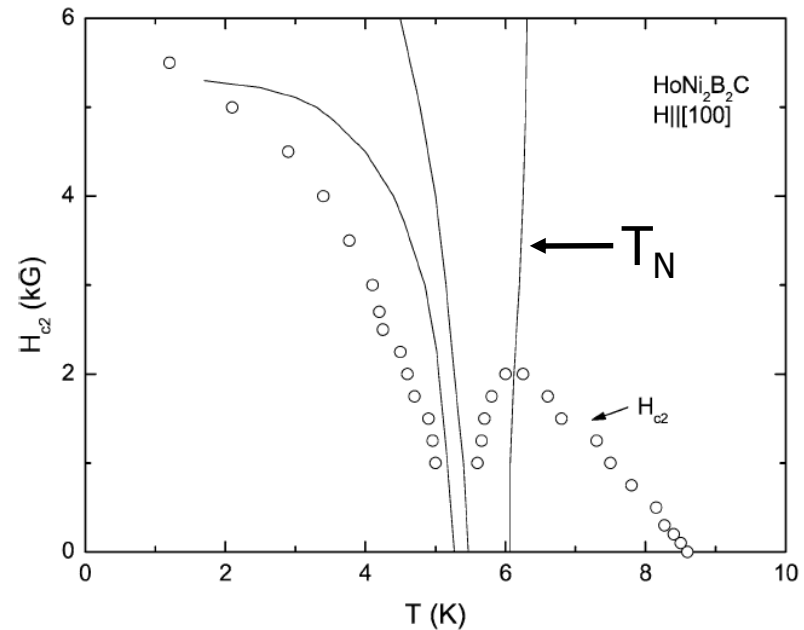
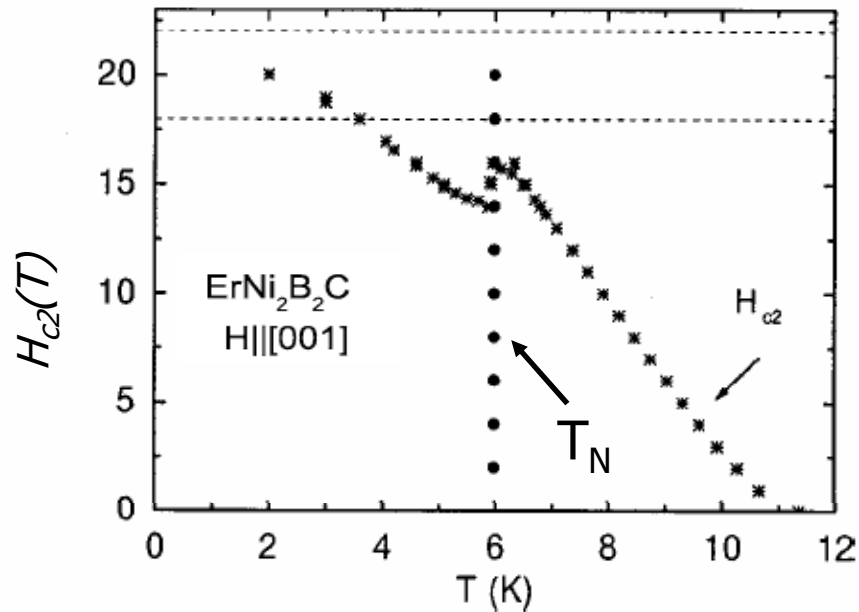


The other conspicuous feature in the  $H_{c2}(T)$  data is the sharp decrease in  $H_{c2}(T)$  when the sample crosses from the paramagnetic to antiferromagnetically ordered state. This can be seen most clearly by creating a composite H-T phase diagram that shows phase lines for both superconductivity as well as local moment order.



This suppression in  $H_{c2}(T)$  has been the focus of intense research and debate. It has been associated with the onset of long range magnetic order that has an ordering wavevector that is determined by Fermi surface nesting (i.e. details of the conduction electron bands).

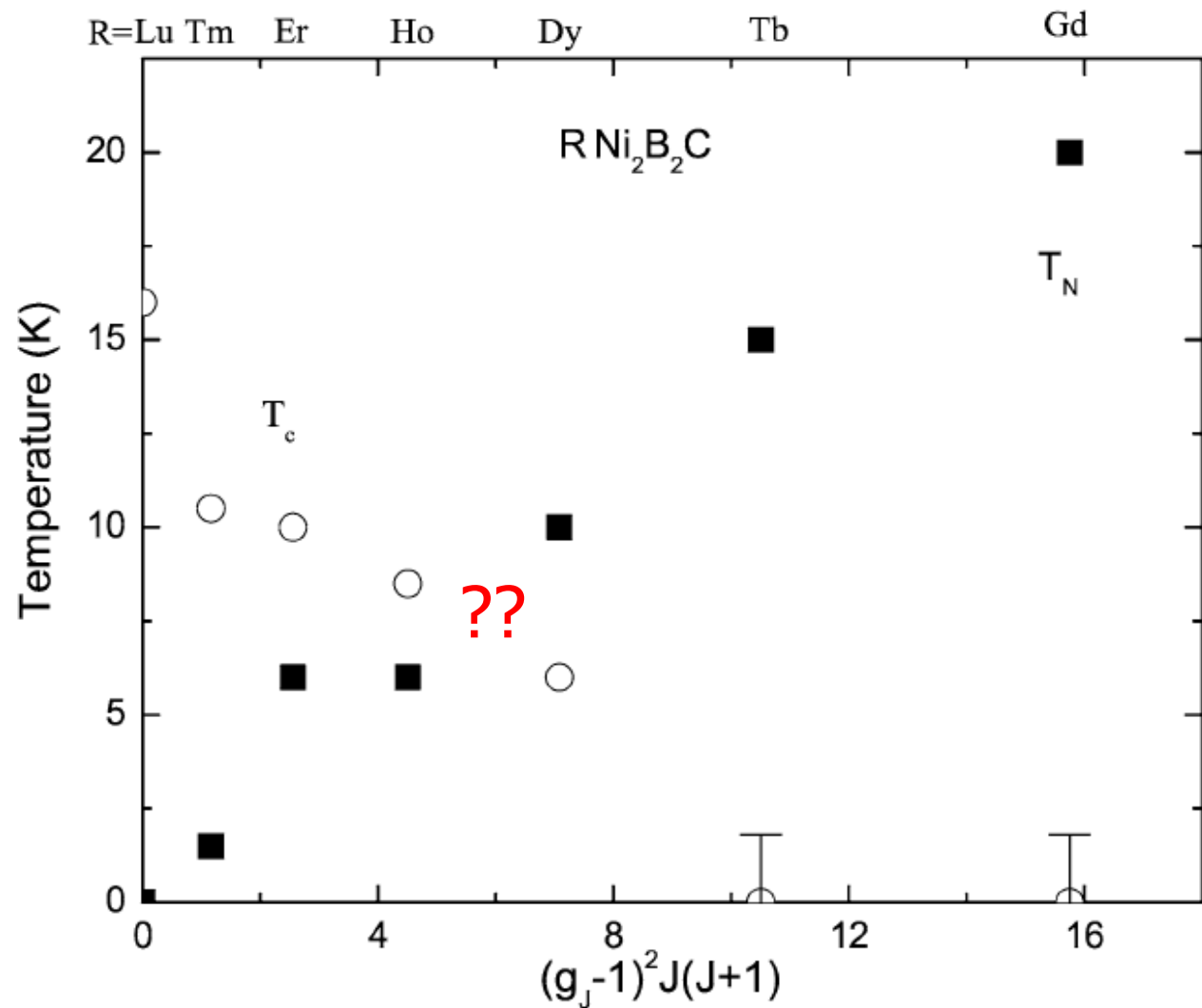
This goes well beyond the scope of an introductory lecture, but is clearly revealed by the simple construction of composite H-T phase diagrams.

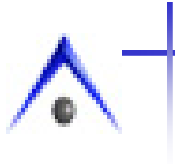




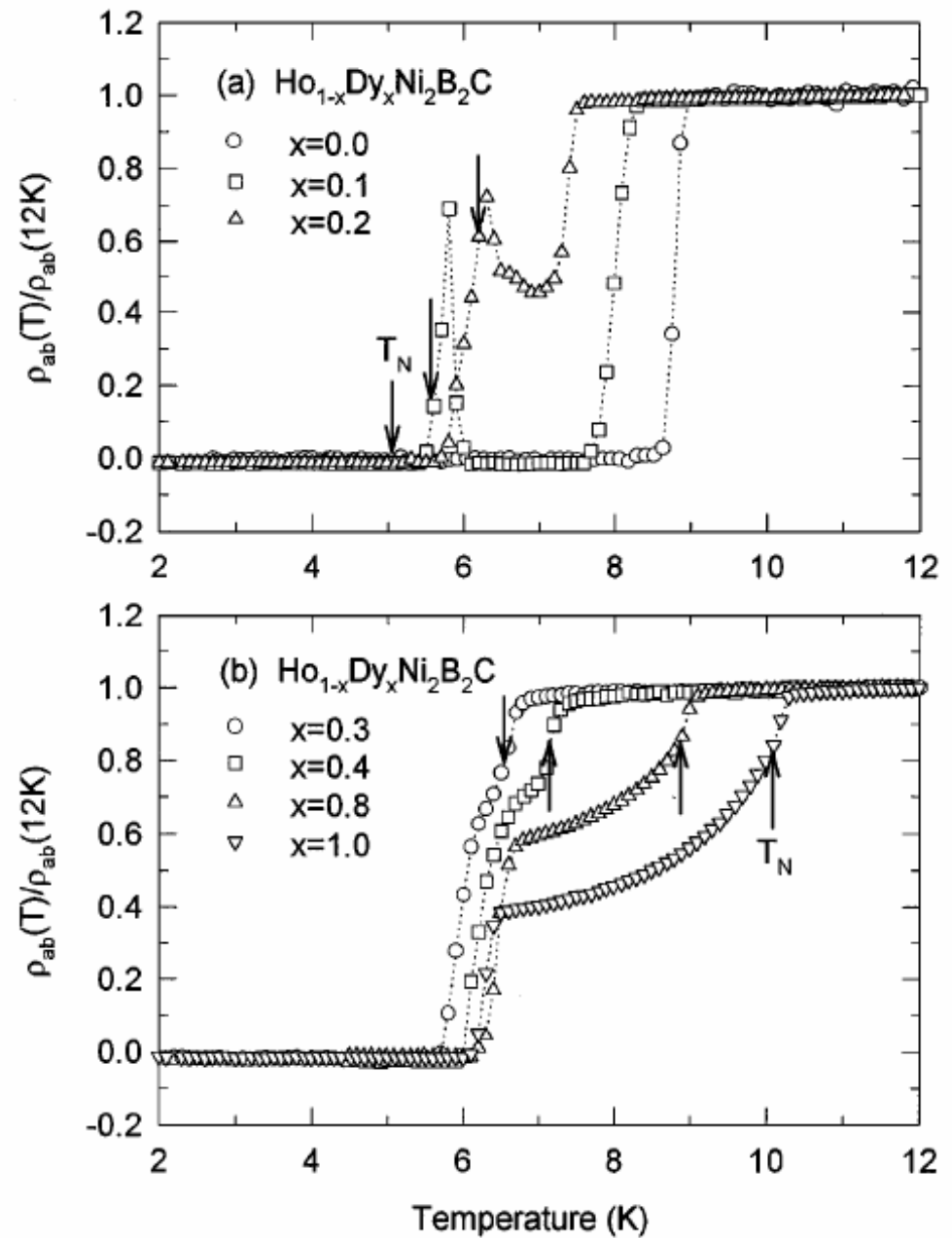
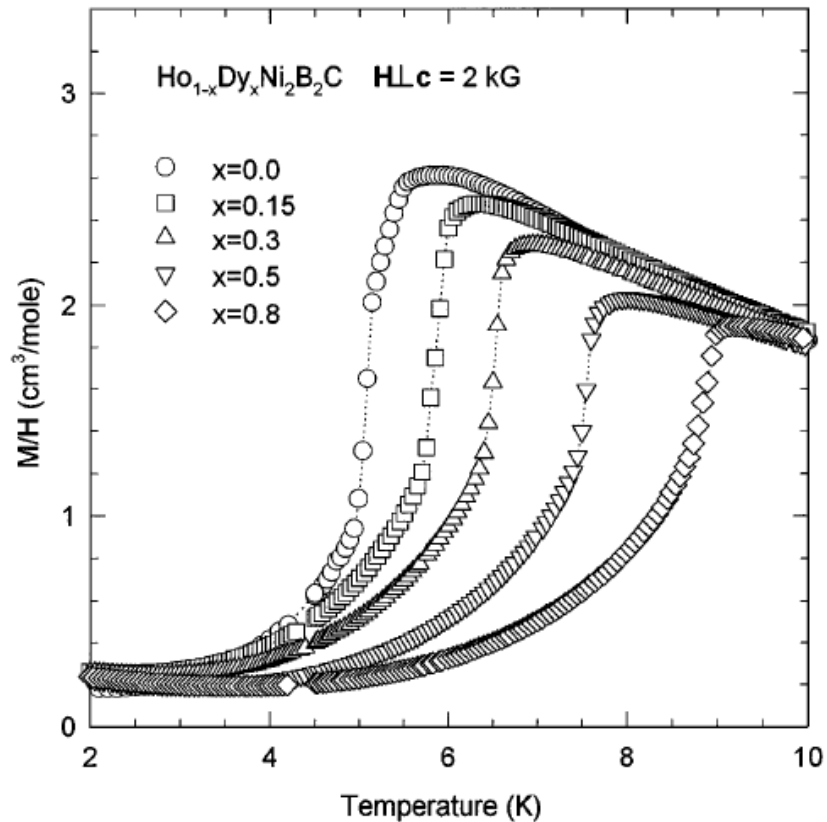
The  $T_c$  versus Rare Earth (or de Gennes factor) plot also brings up what can be called “a clever monkey” question: What happens as we make  $T_c$  cross  $T_N$ ?

In order to address this question we need to grow a specific series of compounds:



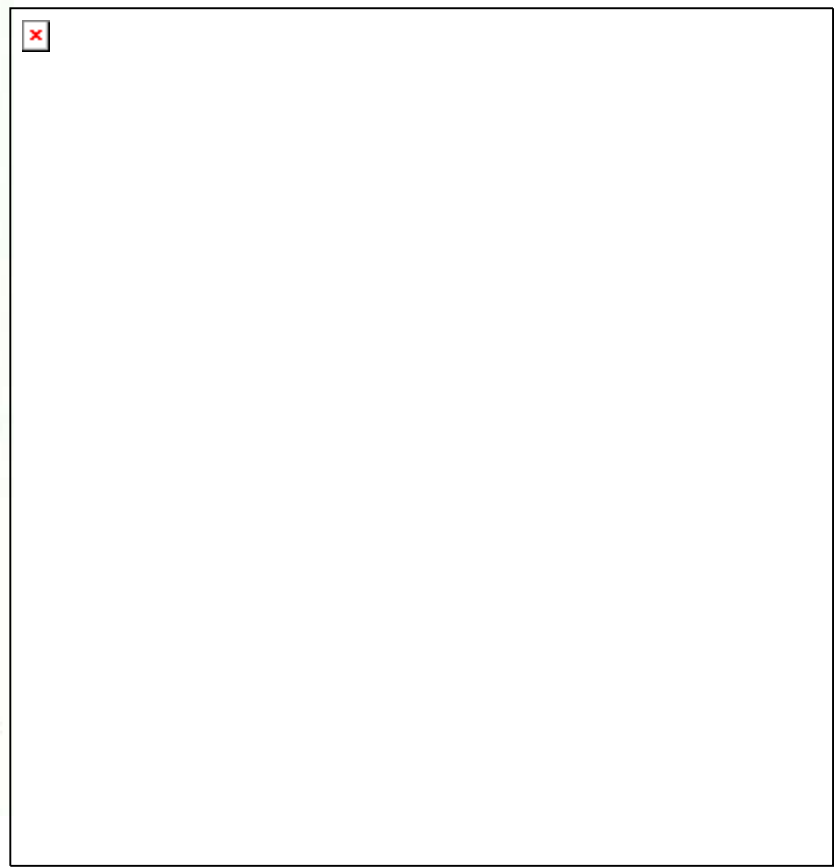
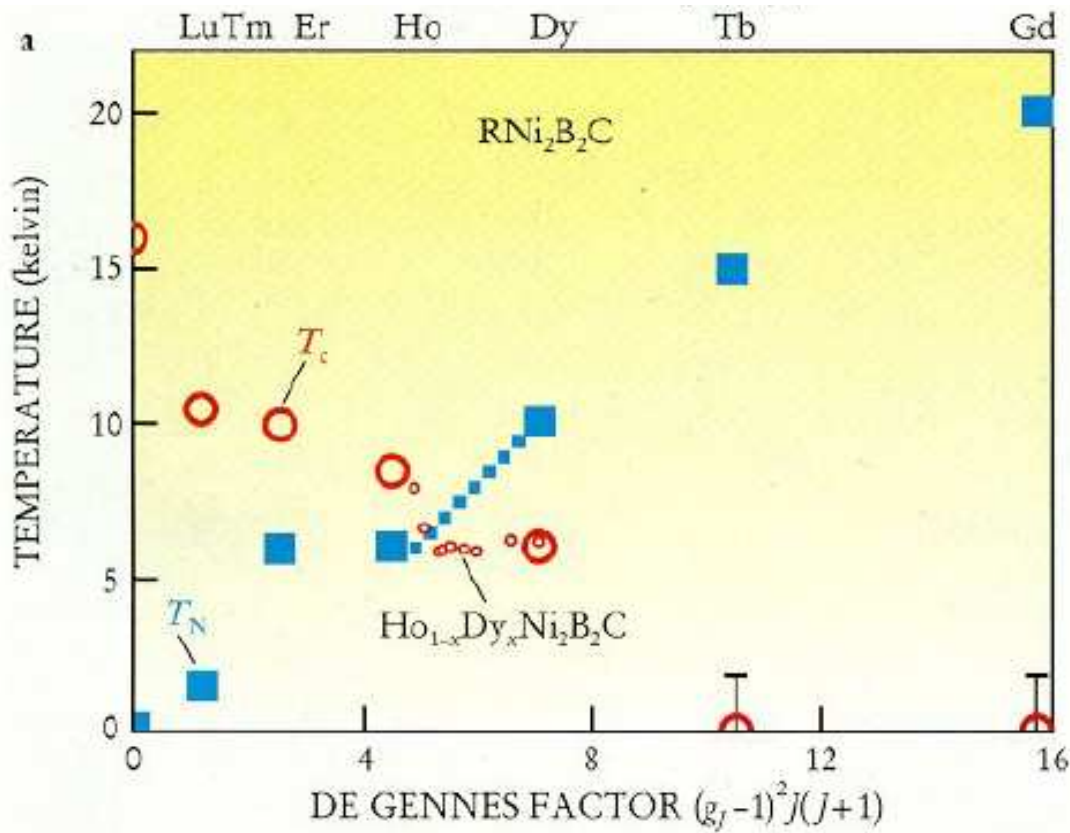


As we discussed last lecture, both  $T_c$  and  $T_N$  can be tracked via  $\chi$  as well as  $\rho$ .





We find that something remarkable happens: Whereas  $T_N$  continues to change in a linear manner,  $T_C$  essentially becomes independent of dG factor when  $T_C < T_N$ .



*P. C. Canfield et al. Physics Today Oct. 1998*

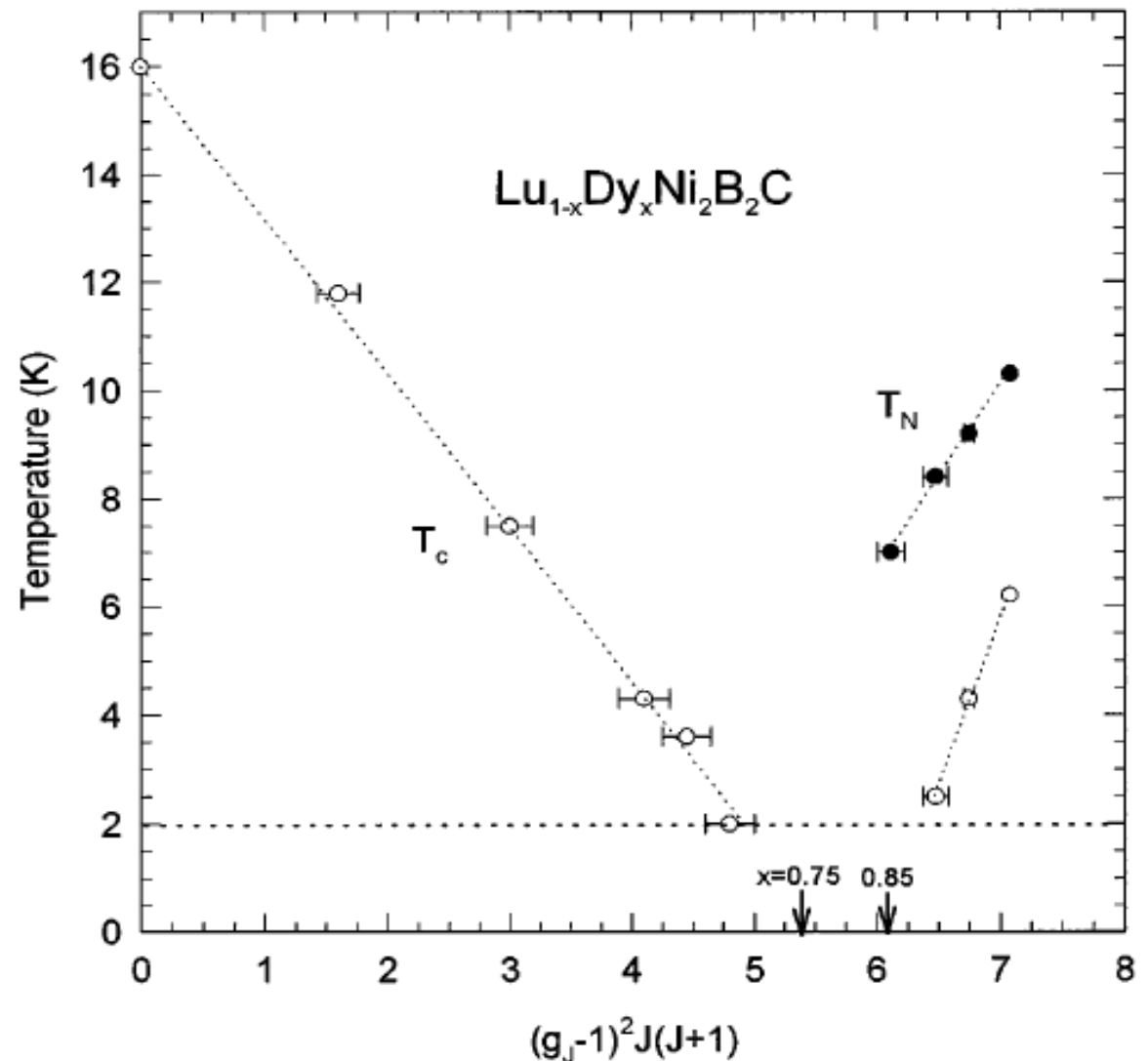
B. K. Cho,\* P. C. Canfield, and D. C. Johnston PHYSICAL REVIEW LETTERS VOLUME 77; 1996 163



In order to get a better idea of what was going on, we retreated to what we thought would be a simpler system:  
 $\text{Lu}_{1-x}\text{Dy}_x\text{Ni}_2\text{B}_2\text{C}$ .

We found that  $\text{DyNi}_2\text{B}_2\text{C}$  was actually a very rare beast indeed: a system that is superconducting only because there is a higher temperature antiferromagnetic transition.

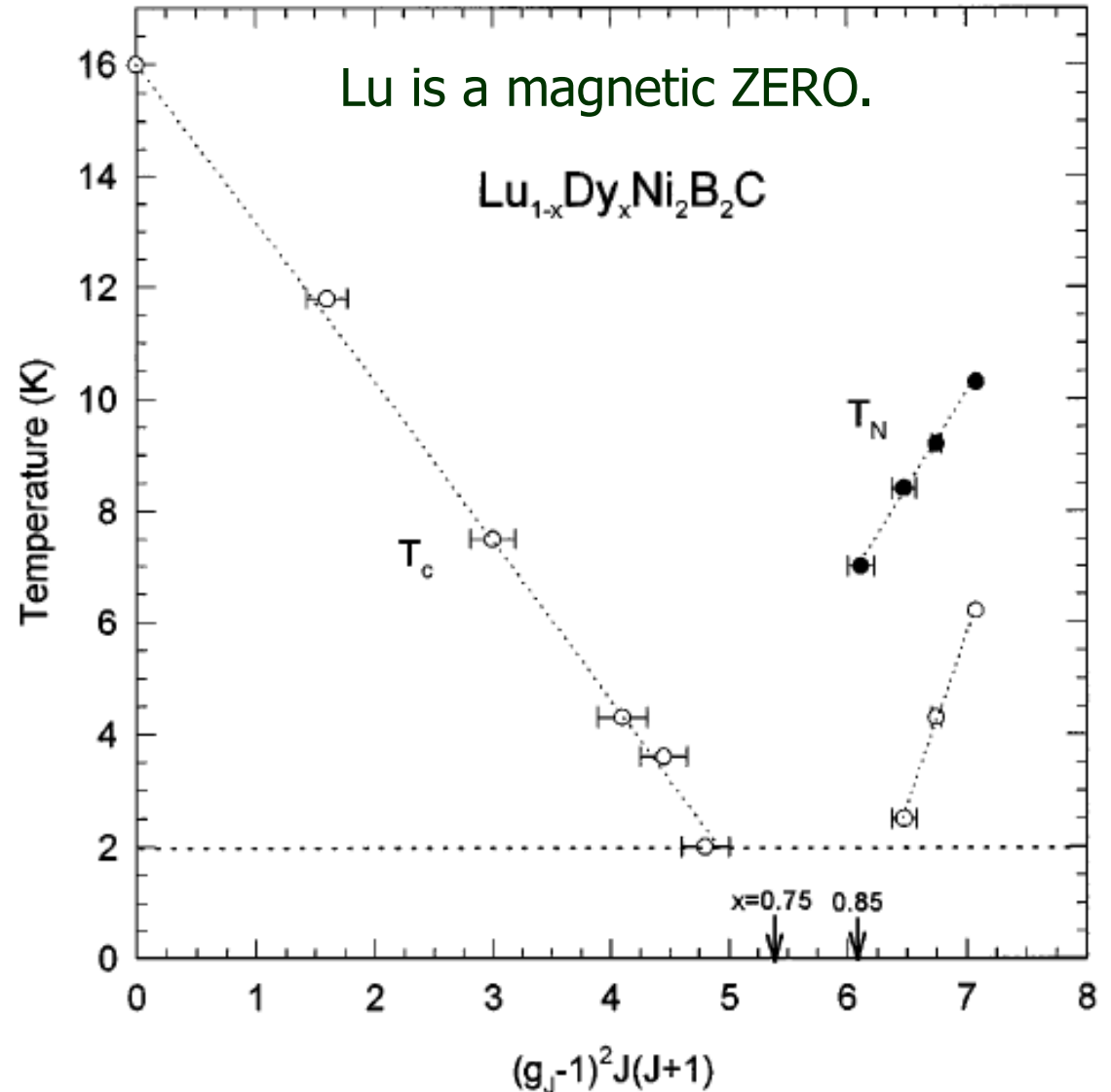
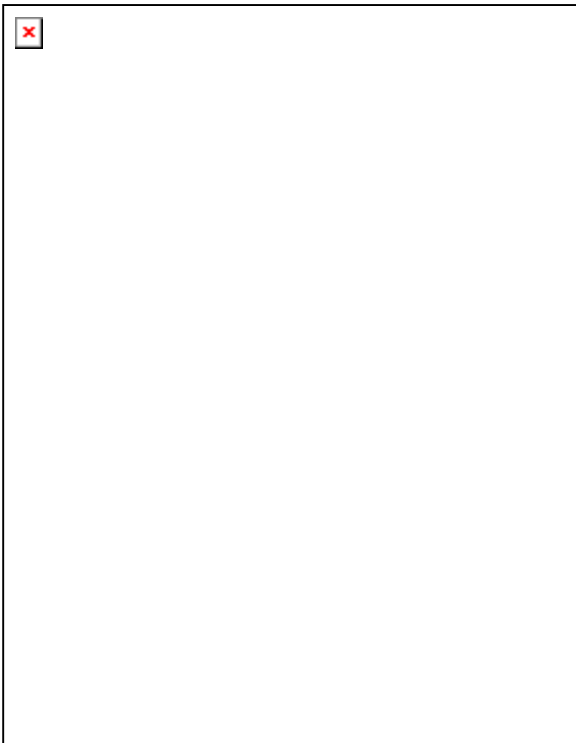
If  $\text{DyNi}_2\text{B}_2\text{C}$  remained paramagnetic, it would never superconduct.





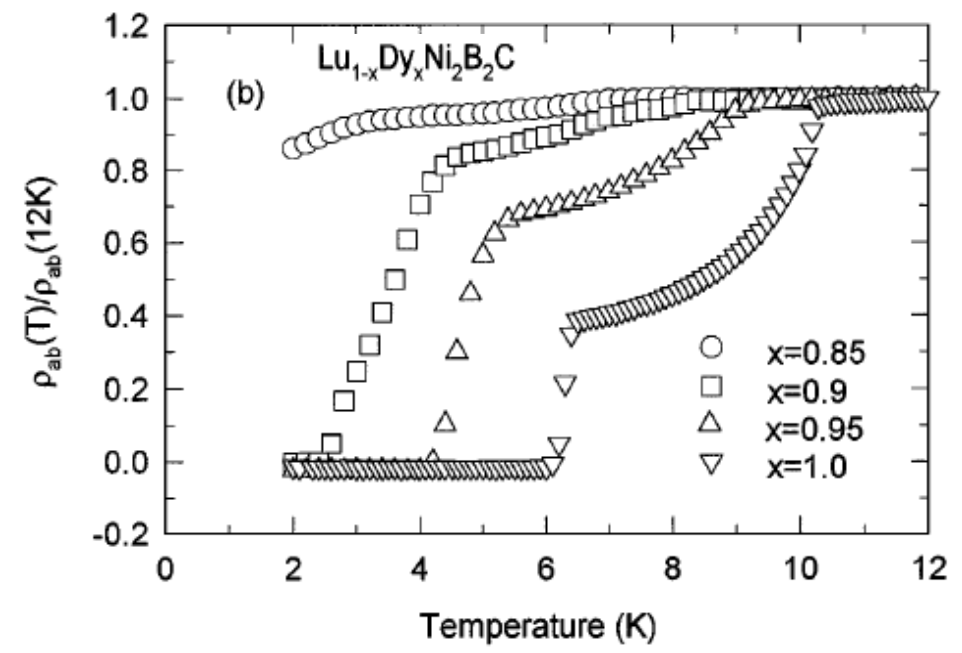
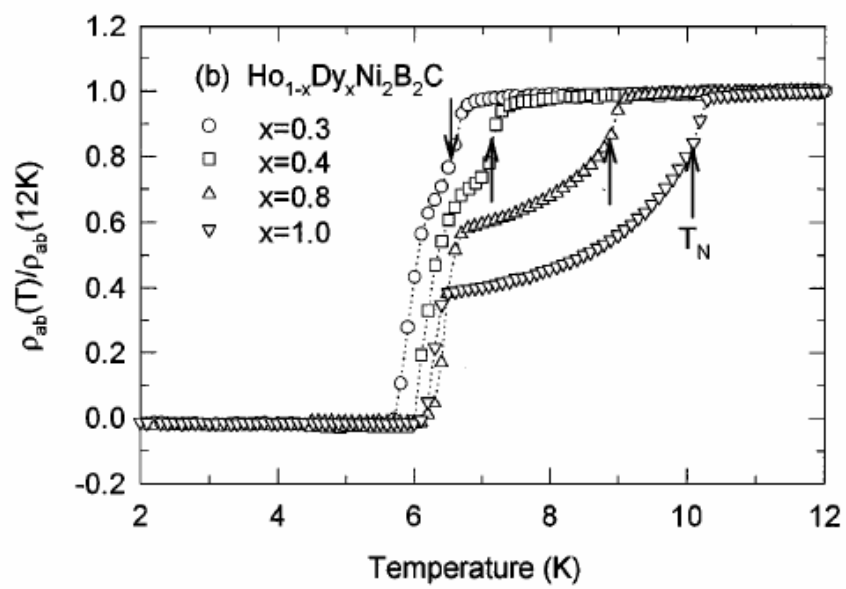
In order to understand this conclusion we need to note that in the ordered state pair breaking is *NOT* from isolated spin flip scattering. The Cooper pair has to interact with magnetic excitations.

Ho and Dy are very similar, especially in the  $RNi_2B_2C$  compounds.



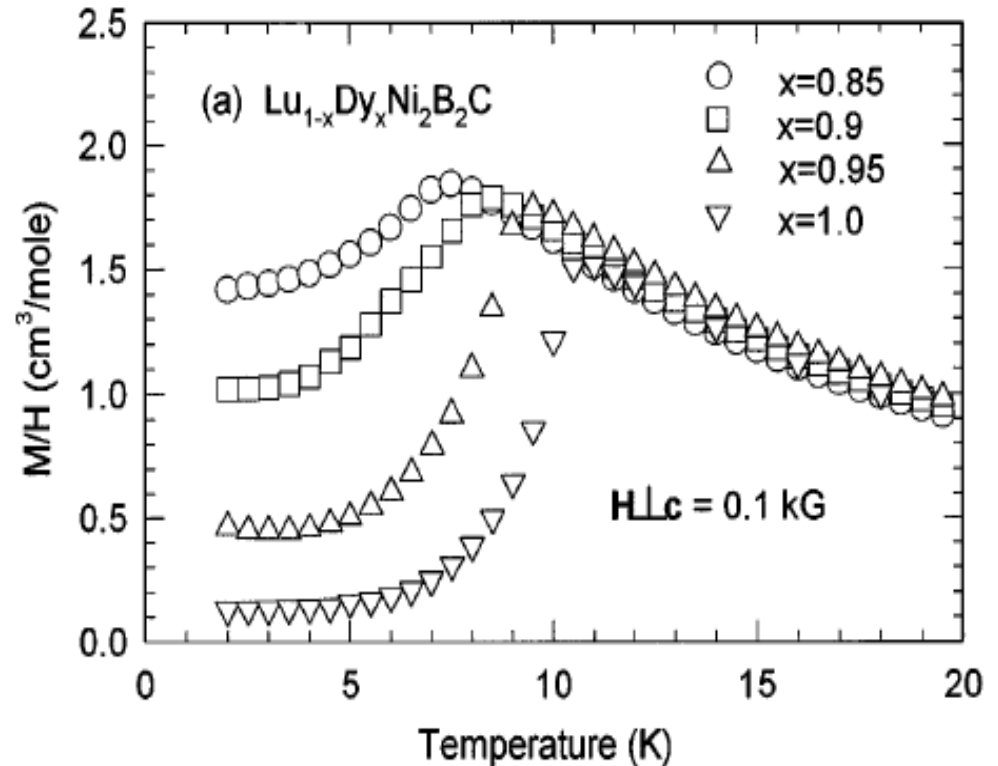
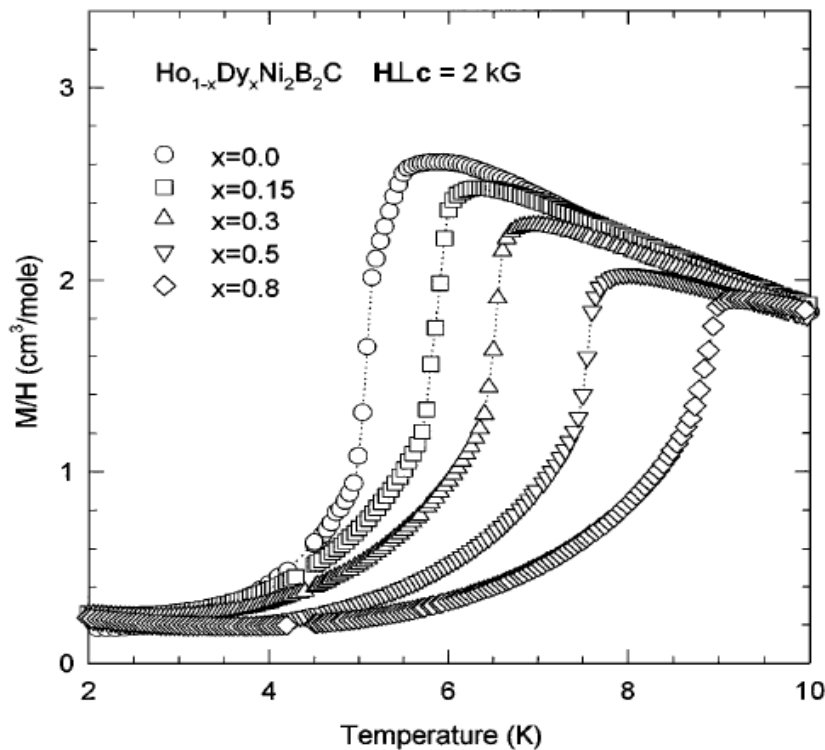


The effects of this can be seen in the  $T < T_N$  resistivity. As Lu is added to  $(\text{Lu}_{1-x}\text{Dy}_x)\text{Ni}_2\text{B}_2\text{C}$  the “loss of spin-disorder scattering” decreases dramatically and rapidly. This is consistent with having the Lu not only suppress  $T_N$ , but also soften the magnetic lattice, allowing many more, low temperature excitations.





The increase of magnetic excitations with the addition of Lu is further confirmed by comparisons of the magnetic susceptibilities of  $(\text{Lu}_{1-x}\text{Dy}_x)\text{Ni}_2\text{B}_2\text{C}$  and  $(\text{Ho}_{1-x}\text{Dy}_x)\text{Ni}_2\text{B}_2\text{C}$ . Whereas the  $\chi(T < T_N)$  remains low for the Ho substitution, it rises rapidly for the Lu substitution.

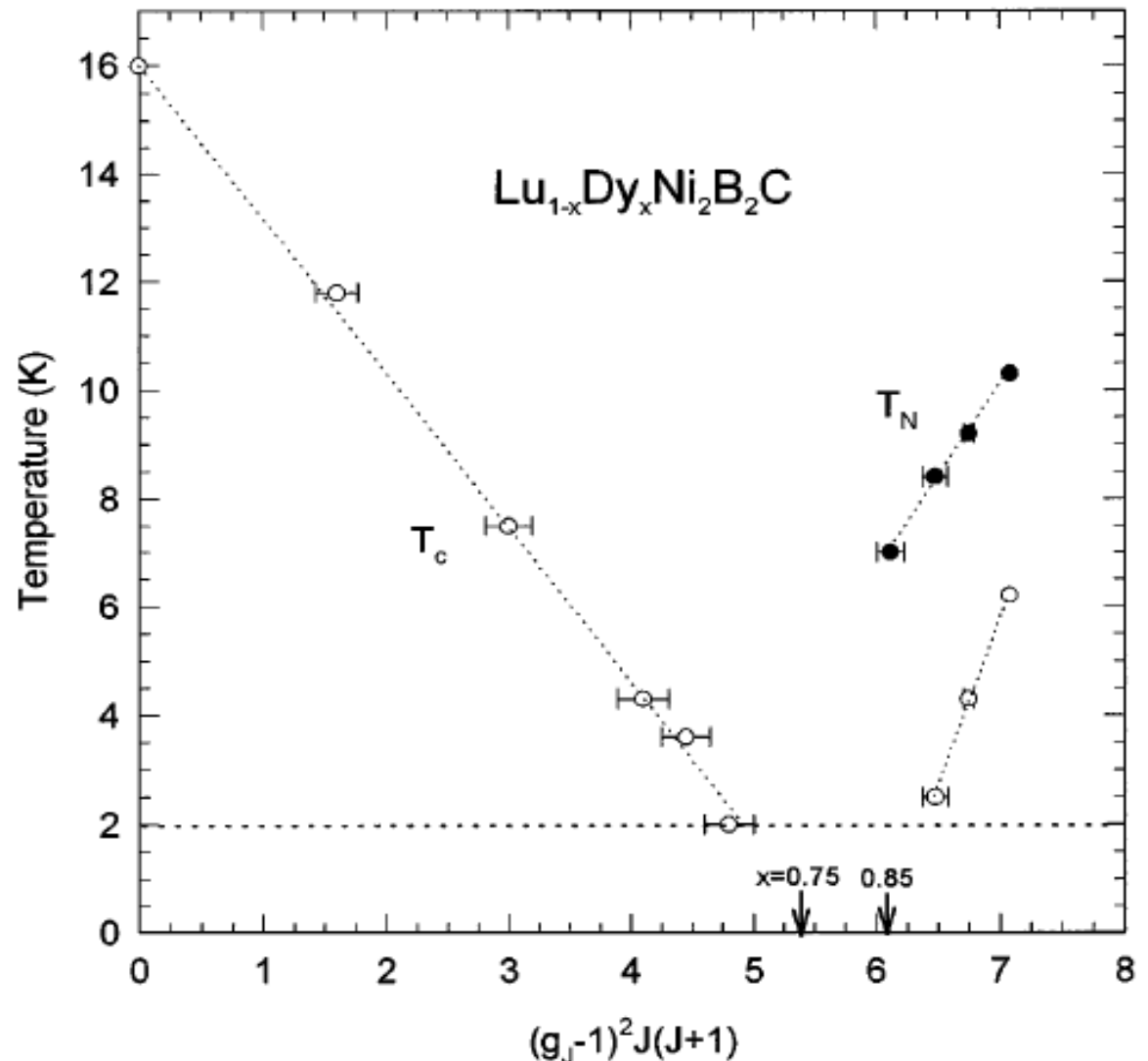




We can look at the T-x phase diagram of  $\text{Lu}_{1-x}\text{Dy}_x\text{Ni}_2\text{B}_2\text{C}$  from both ends. As we add Dy to  $\text{LuNi}_2\text{B}_2\text{C}$  we suppress  $T_c$  via the paramagnetic Dy breaking Cooper pairs.

On the other hand, as we add Lu to  $\text{DyNi}_2\text{B}_2\text{C}$   $T_N$  is suppressed rapidly and pair breaking from low-energy excitations is enhanced dramatically.

If  $\text{DyNi}_2\text{B}_2\text{C}$  only superconducts because it is antiferromagnetic, then any degradation of the antiferromagnetic state is very bad for the superconductivity.





In addition to the interplay between superconductivity and local moment magnetism, the  $\text{RNi}_2\text{B}_2\text{C}$  series has:

Heavy Fermion for  $\text{YbNi}_2\text{B}_2\text{C}$

$T_K \sim 10 \text{ K}$ ;  $T_{\Delta\text{CEF}} \sim 100 \text{ K}$ ;  $T_C, T_N < 100 \text{ mK}$  (possibly 0 K)

Flux Line Lattice phase transitions

Great tunability:  $T_N, H_{C2}, \ell, \xi_0, H_2 \dots$

This series offers a wonderful playground for the study of the interactions between conduction electrons (normal or superconducting) and 4f-electrons (localized or itinerant).

---

But for this lecture, let's return to some of the old prejudices:  
remember #2, highest  $T_C$  near a structural phase transition.



LuNi<sub>2</sub>B<sub>2</sub>C (and YNi<sub>2</sub>B<sub>2</sub>C) are about as close to a CDW transition as they can be....Electron-phonon coupling is nearly optimized.

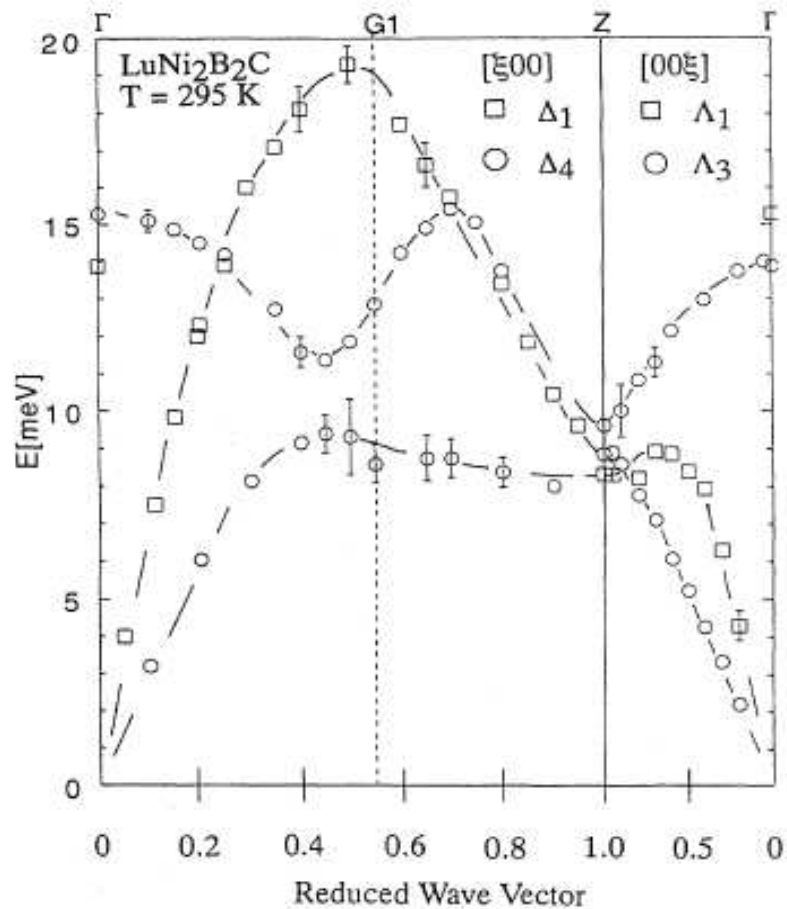


FIG. 1. Room-temperature acoustic and lowest-lying optical phonon dispersion curves of LuNi<sub>2</sub>B<sub>2</sub>C along the  $[\xi 00]$  and  $[00\xi]$  symmetry directions. The lines are intended as guides to the eye. The size of the symbols is a measure of the estimated uncertainties in the measured frequencies.

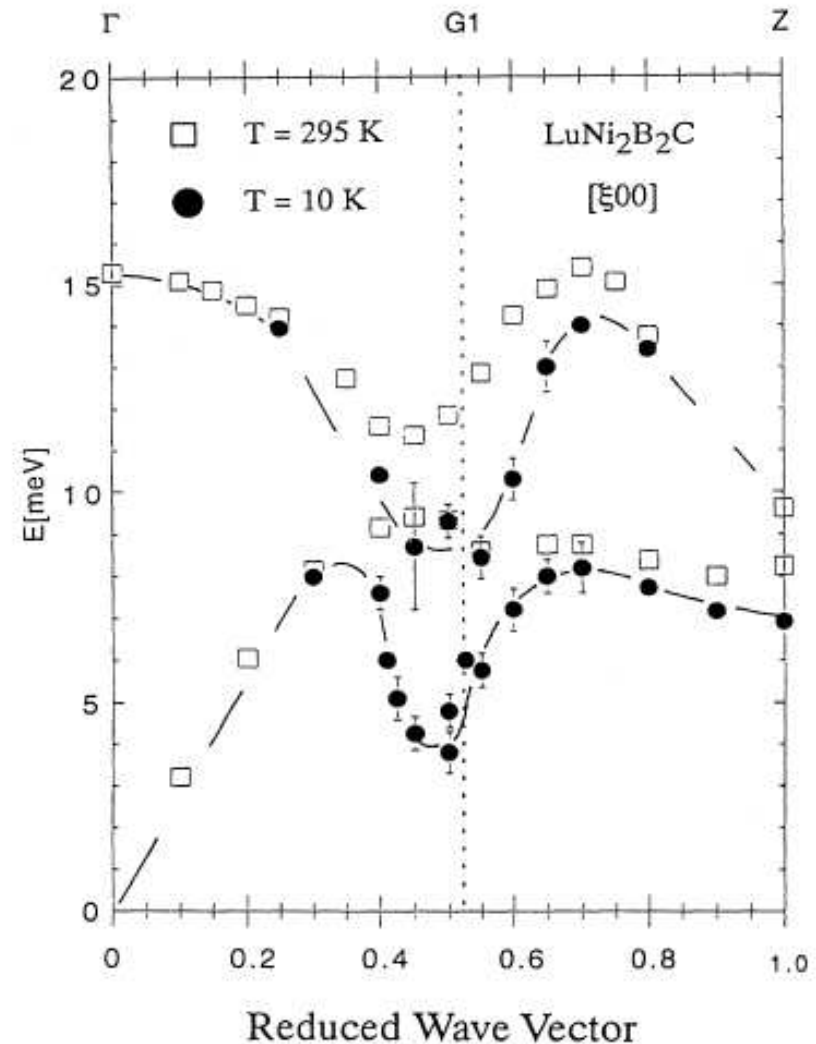
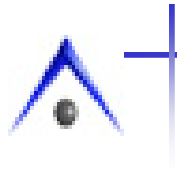


FIG. 2. The  $\Delta_4$   $[\xi 00]$  branches at 295 and 10 K. The lines through the 10 K points are intended as guides to the eye.

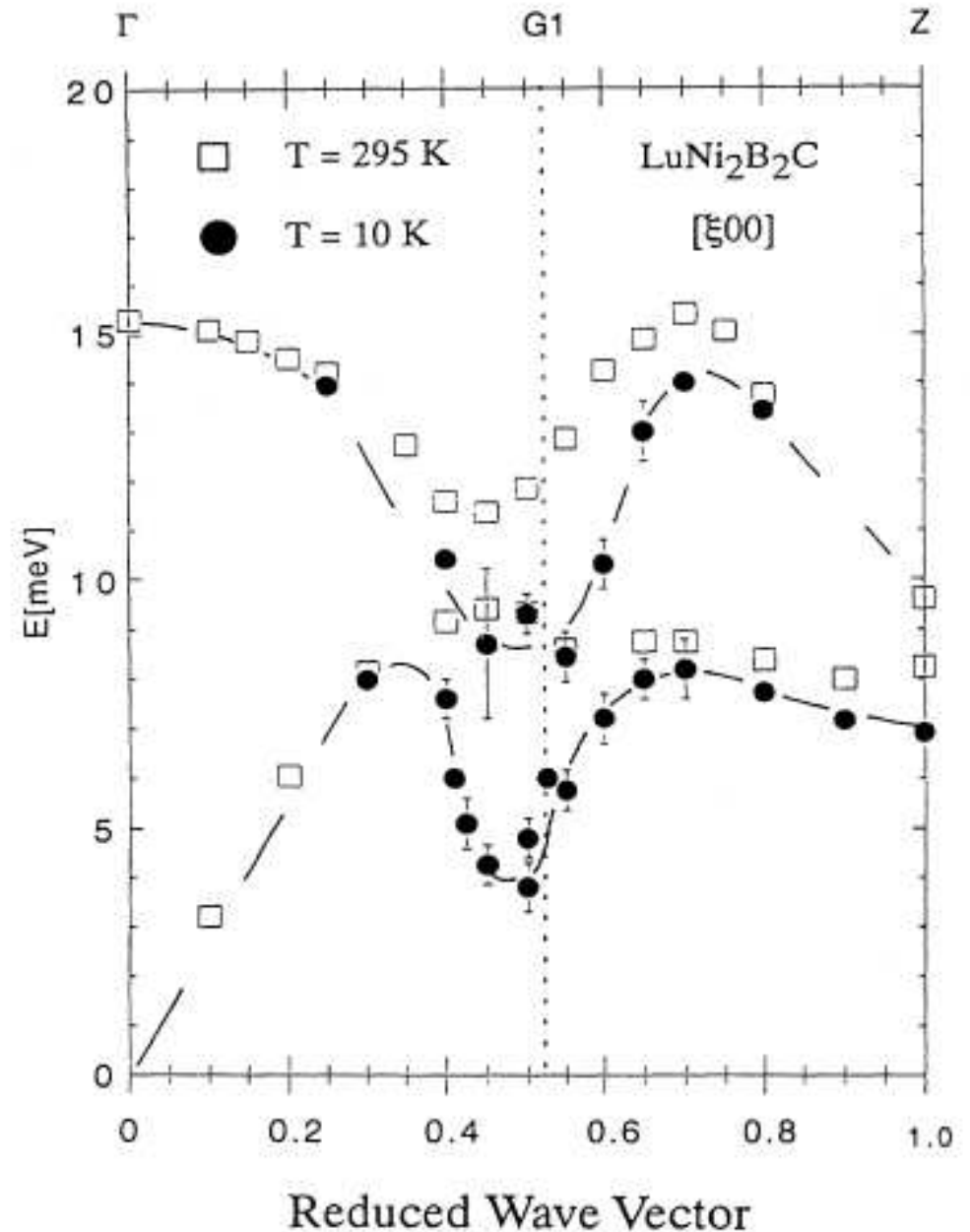


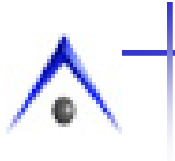
Near the end of what

$$k_B T_C = 1.13 \hbar \omega_D e^{-1/VN(E_F)}$$

has to offer.

If the electron phonon coupling were increased further, then the phonon would soften further and there would be a structural phase transition.... This is consistent with the higher  $T_C$  and metastability of  $YPd_2B_2C$ .





*In the late 1990's the question arose:*


## ***So, what to do next?***

Based on the wealth of physics, as well as the high  $T_C$  values, found in the  $RNi_2B_2C$  family many groups decided to look for other intermetallics with light elements and see if similar (or even higher)  $T_C$  values could be found.

Several groups (including my own) were examining compounds with combinations of Li, Be, B, C, Mg, Al, Si and other (often transition metal) elements.

In late 2000 the group lead by Prof. J. Akimitsu examined the Ti-Mg-B ternary (*Ti because we got to have those 3-d electrons*) and found...a binary:  $MgB_2$ .

In mid-January, 2001 Prof. Akimitsu announced an  $\sim 40$  K  $T_C$  in  $MgB_2$  as part of a passing reference in a talk at a meeting....



In mid-January, 2001 my group heard rumors of this announcement: i.e. that there may be superconductivity in MgB<sub>2</sub>.

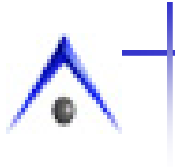
Immediately there were four basic questions.

Can we make it?

Can we confirm  $T_c$ ?

Can we address the mechanism of superconductivity?

Can we delineate its basic properties?

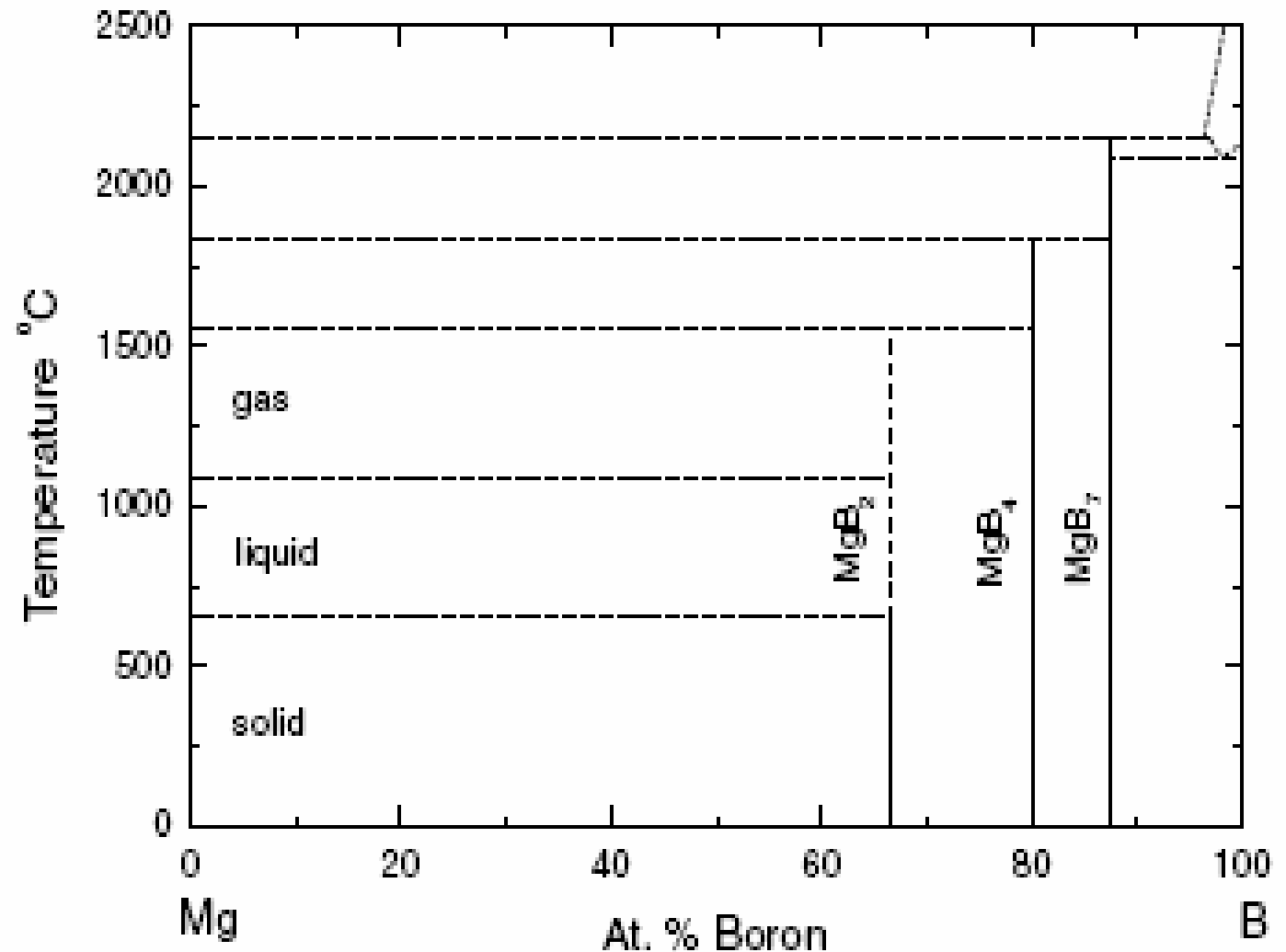


# MgB<sub>2</sub>:

Can we make it?

This is a very inauspicious phase diagram. There is no exposed liquid–solidus line at all.

This is not promising for crystal growth.





Since the Department of Physics at Iowa State University is in the College of Liberal Arts and Sciences, we teach the classics of philosophy: Socrates and Aristotle, Freud and Jung, etc. In the words of Mick and Keith,



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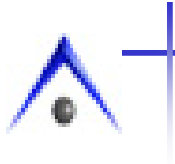


*You can't always get what you want*

*But sometimes you get what you need*



Even though, as a group, we specialize in making measurements on single crystals, we found a lot of good work could be done of very high purity polycrystalline samples....

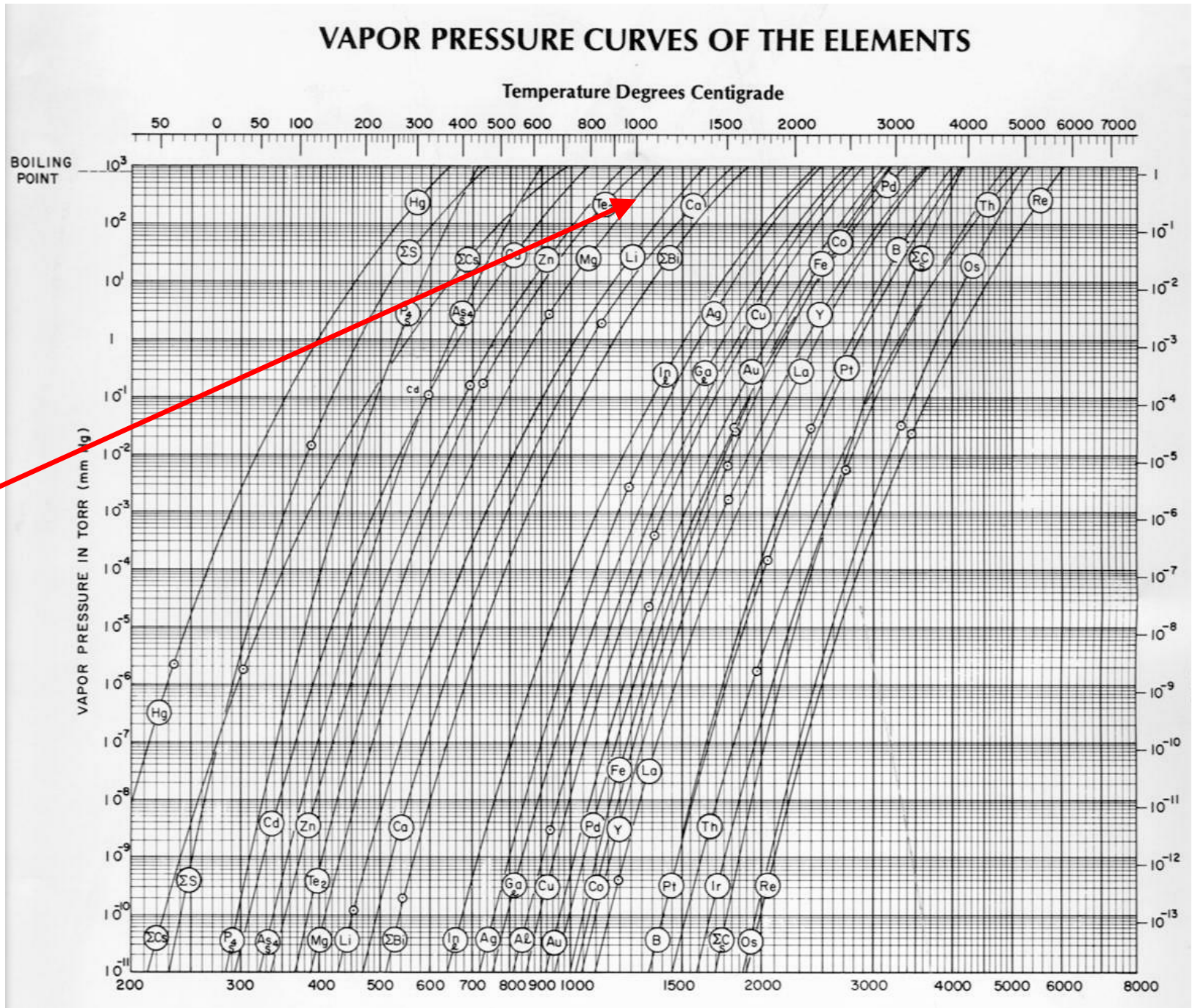


# VAPOR PRESSURE CURVES OF THE ELEMENTS

Often vapor pressure is a problem

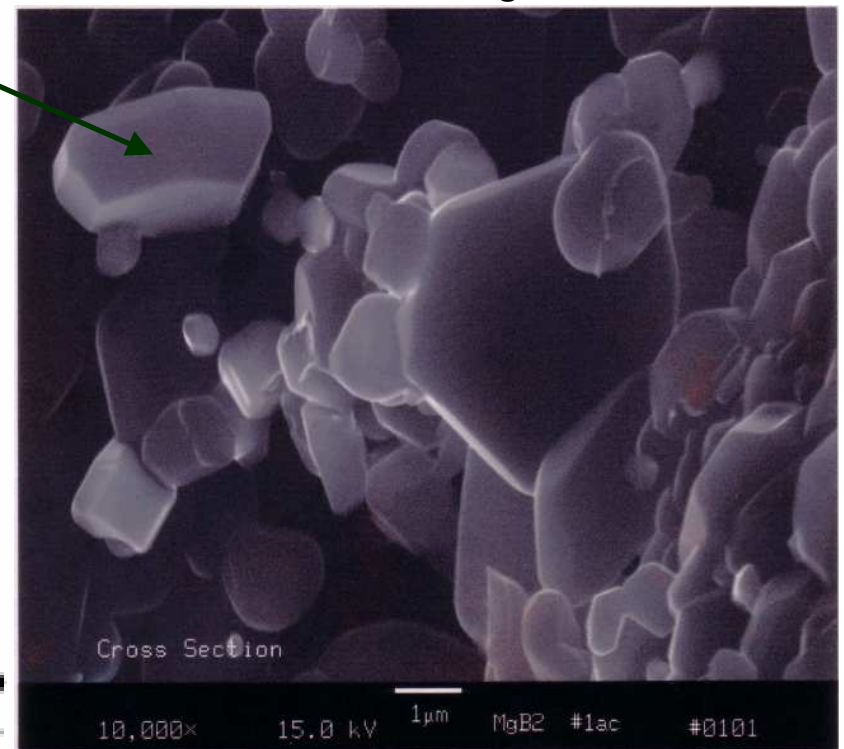
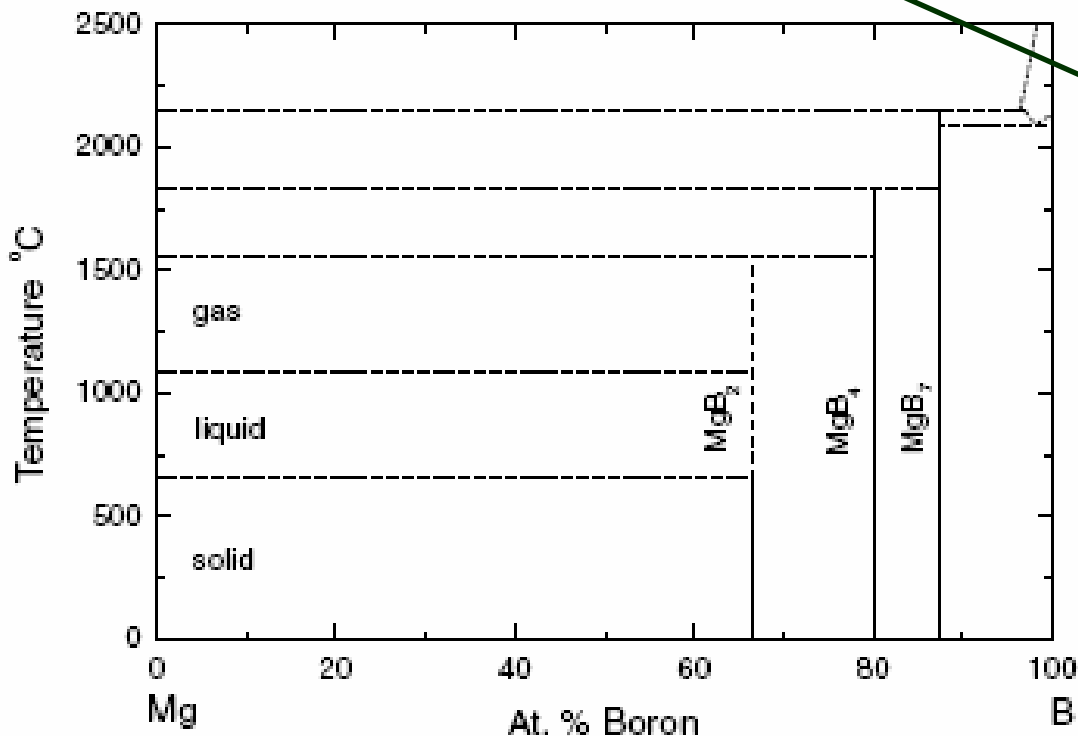
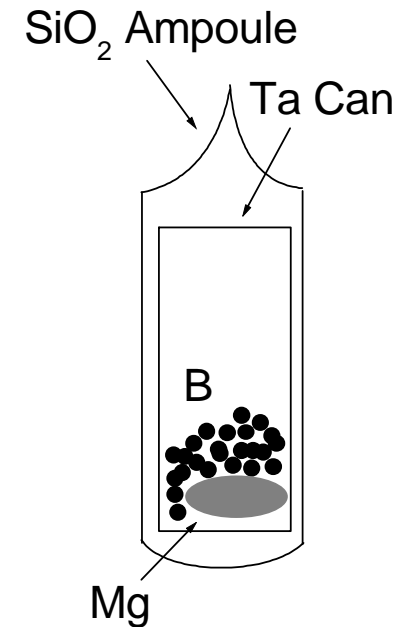
At 950 C Mg has ~ 1/3 atm vapor pressure

*This time we use it to our advantage!*



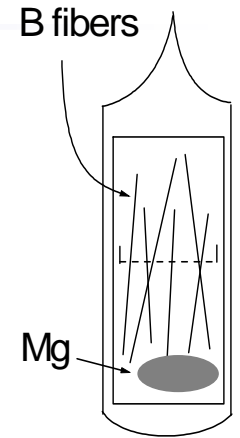


At 950 C the B powder is nowhere close to melting, but we found that it will react with Mg vapor (~1/3 atm. at 950 C) to form  $\text{MgB}_2$  within as little as 2 hours.

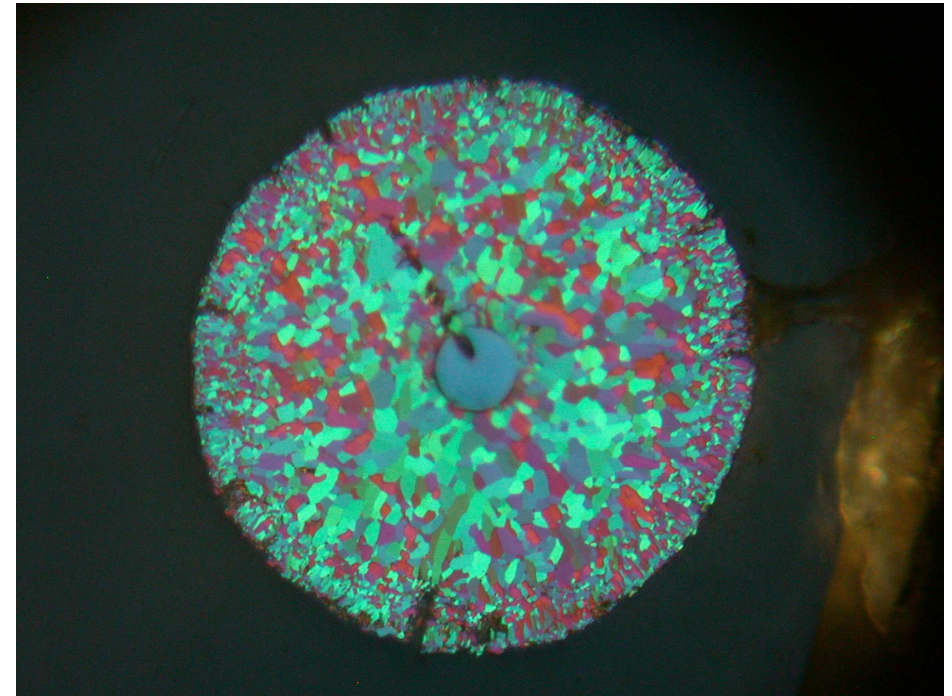
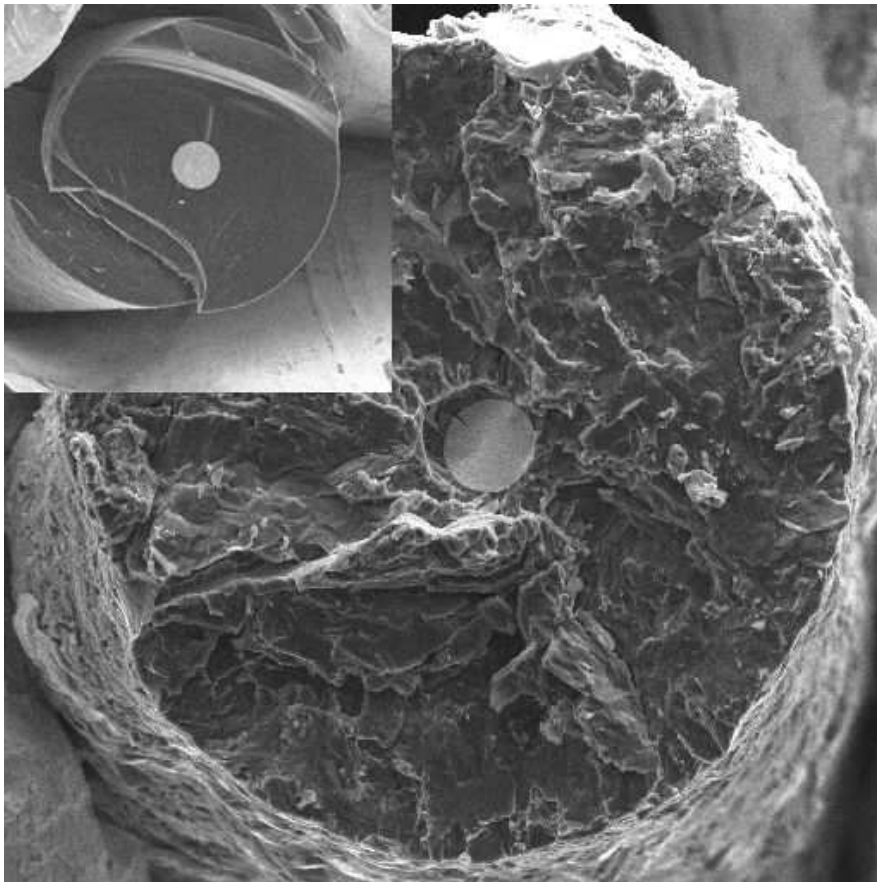




Boron comes in many forms.  $\text{MgB}_2$  can be made as sintered pellets, thin films and...



## Dense $\text{MgB}_2$ Wires



$\text{MgB}_2$  from boron filament

Canfield et al., PRL 86 (2001) 2423

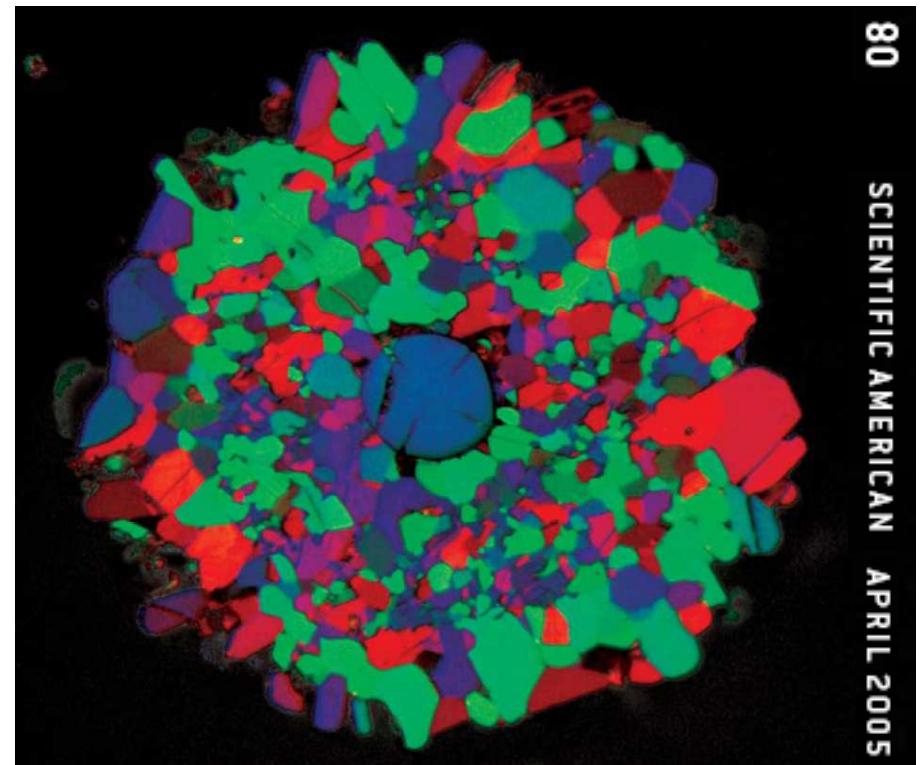
$\text{MgB}_2$  grains illuminate under polarized light.



# Sometimes you can do a lot with polycrystalline samples....

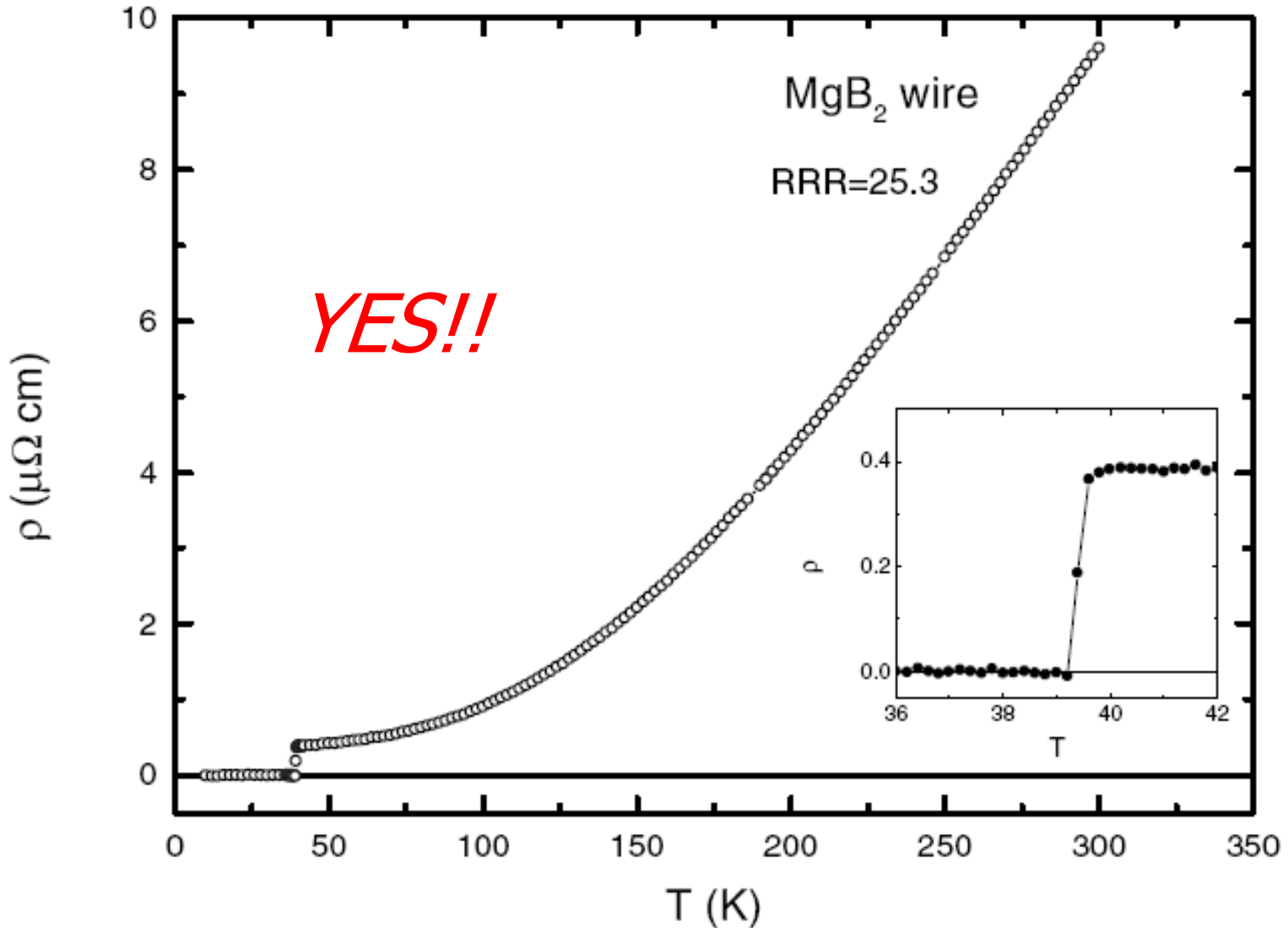


VOLUME 86, NUMBER 9	PHYSICAL REVIEW LETTERS	26 FEBRUARY 2001
<b>Boron Isotope Effect in Superconducting MgB<sub>2</sub></b>		
VOLUME 86, NUMBER 11	PHYSICAL REVIEW LETTERS	12 MARCH 2001
<b>Thermodynamic and Transport Properties of Superconducting Mg<sup>10</sup>B<sub>2</sub></b>		
VOLUME 86, NUMBER 11	PHYSICAL REVIEW LETTERS	12 MARCH 2001
<b>Superconductivity in Dense MgB<sub>2</sub> Wires</b>		
VOLUME 87, NUMBER 4	PHYSICAL REVIEW LETTERS	23 JULY 2001
<b>Anisotropy of Superconducting MgB<sub>2</sub> as Seen in Electron Spin Resonance and Magnetization Data</b>		





# Can we confirm $T_C$ ?

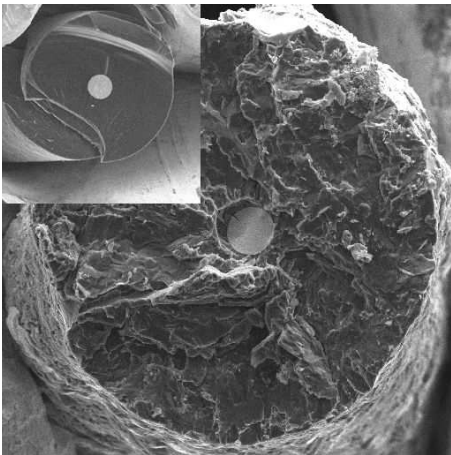




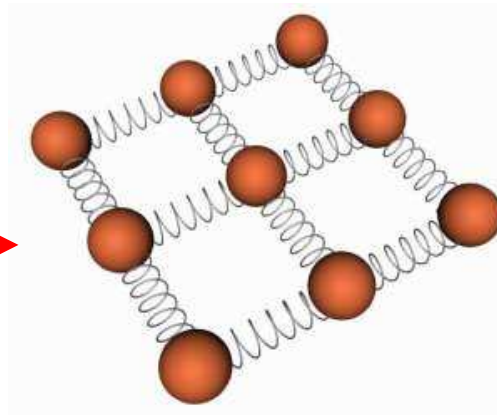
## Can we address the mechanism of superconductivity?

If we start with the BCS equation for  $T_C$ :  $k_B T_C = 1.13 \hbar \omega_D e^{-1/VN(E_F)}$   
we can see that  $T_C$  is directly proportional to  $\omega_D$ .

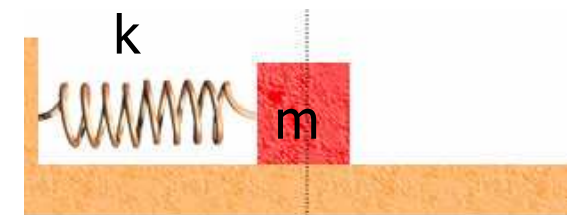
Real sample



Ball-spring model  
of lattice



Single mass  
and spring



$$\omega = \sqrt{k/m}$$

So, within this cascade of gross assumptions,  
 $T_C$  should vary as  $1/\sqrt{m}$



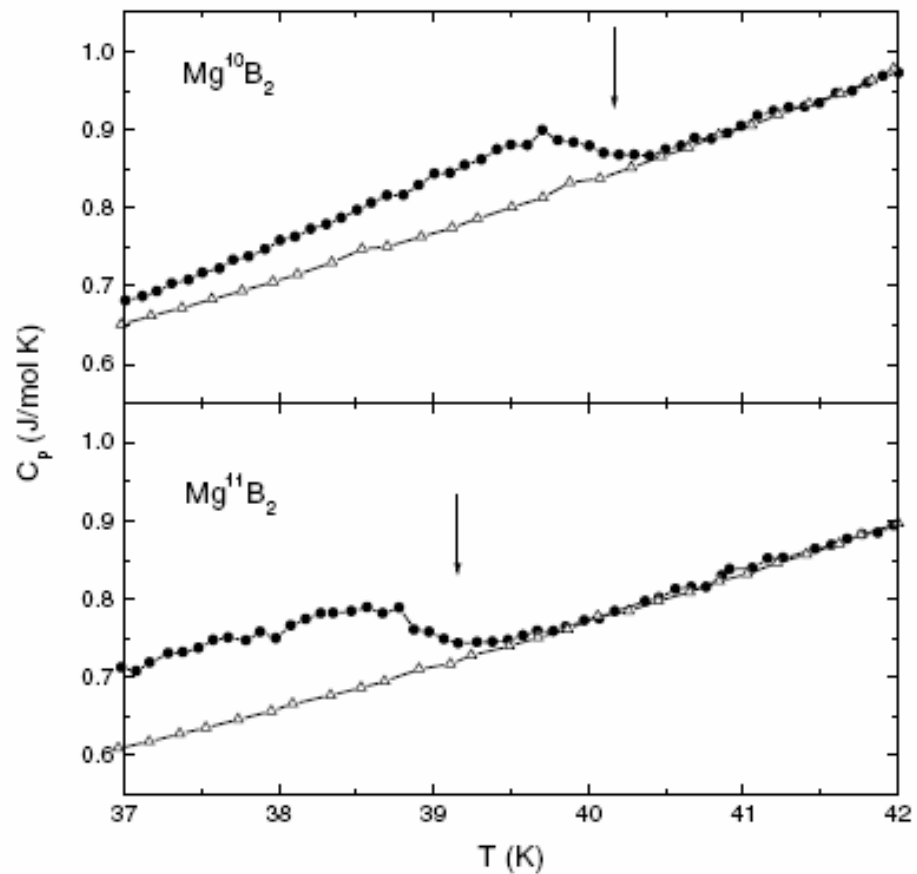
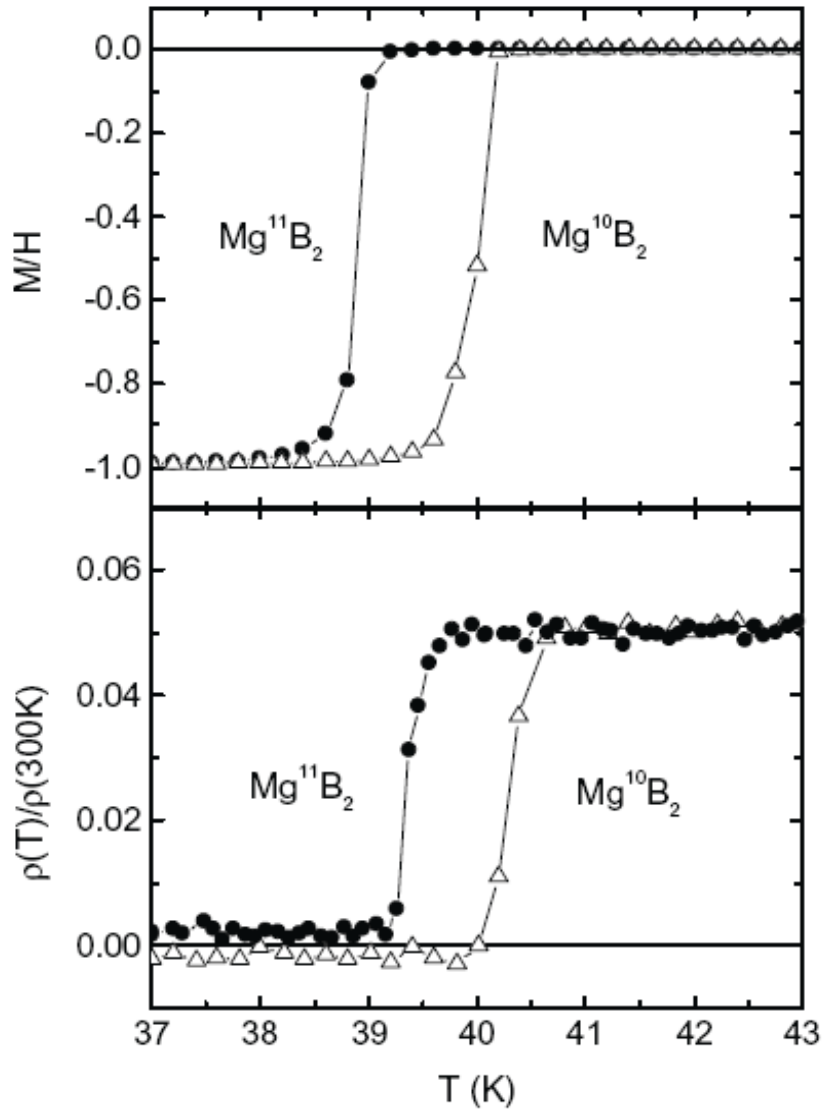
# Boron Isotope Experiment

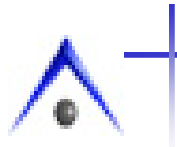
$$\Delta T = 1.0 \text{ K}$$

(for simplest model expected  $\Delta T_C \sim 0.85$ )

$$T_C \sim M^{\alpha_B}, \alpha_B = 0.26$$

Consistent with  
phonon-mediated BCS.





From  $C_p(T)$  data we learn:

For  $\text{MgB}_2$   $\theta_D \sim 750$  K

For Mg  $\sim 320$  K

For Si  $\sim 625$  K

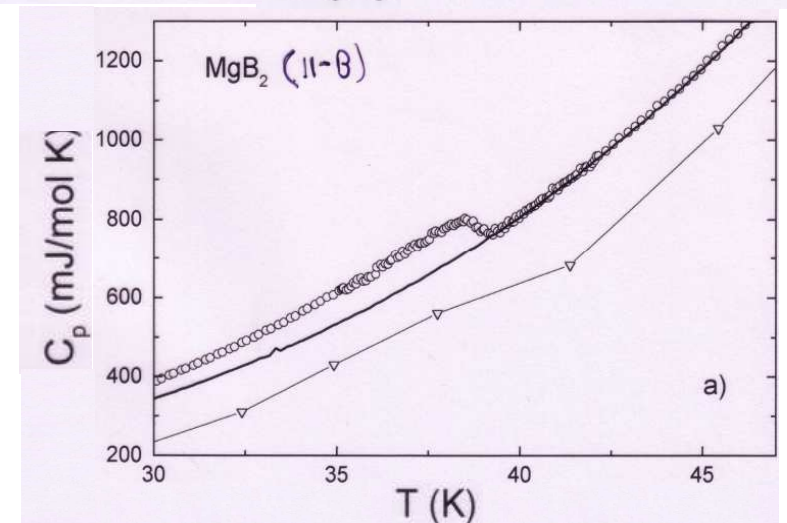
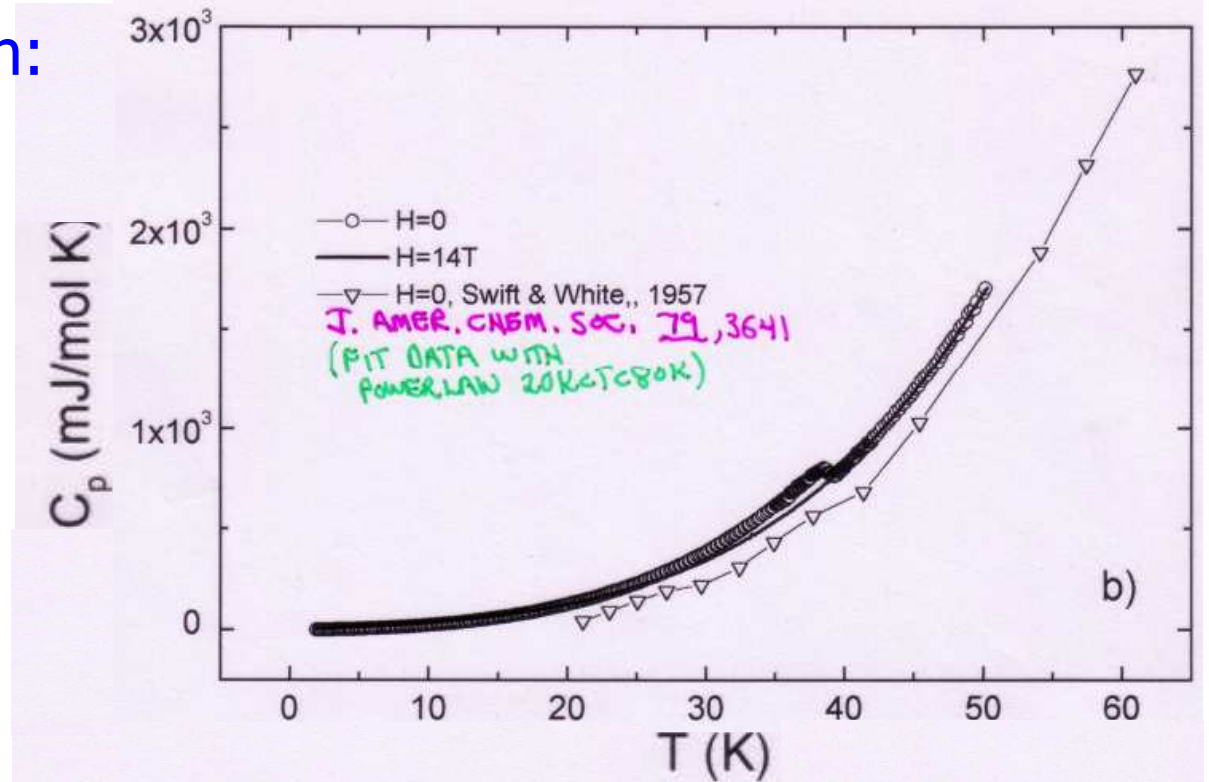
For Diamond  $\sim 1860$  K

(high  $\omega_D$  for  $\text{MgB}_2$ )

$\gamma \sim 2.5 \pm 0.75$  mJ/mole- $\text{K}^2$

(small  $\gamma$  means small  $N(E_F)$ )

**But we should not forget the triangles....**





How was superconductivity in  $\text{MgB}_2$  missed?

Hard to make by simple methods (cannot arc-melt it)

Did not fit prejudice (no d-shell electrons to boost  $N(E_F)$ )

....(Old data did not show superconductivity)....

*$\text{MgB}_2$  forces a shift in emphasis when looking for other higher  $T_C$  compounds*

$$k_B T_C = 1.13 \hbar \omega_D e^{-1/VN(E_F)}$$

Old prejudice: Need to have large  $\omega_D$  and large  $N(E_F)$  and hope for good  $V$

But  $\text{MgB}_2$  has a very small  $N(E_F)$  ( $\gamma \sim 2.5$  mJ/mole-K<sup>2</sup>)

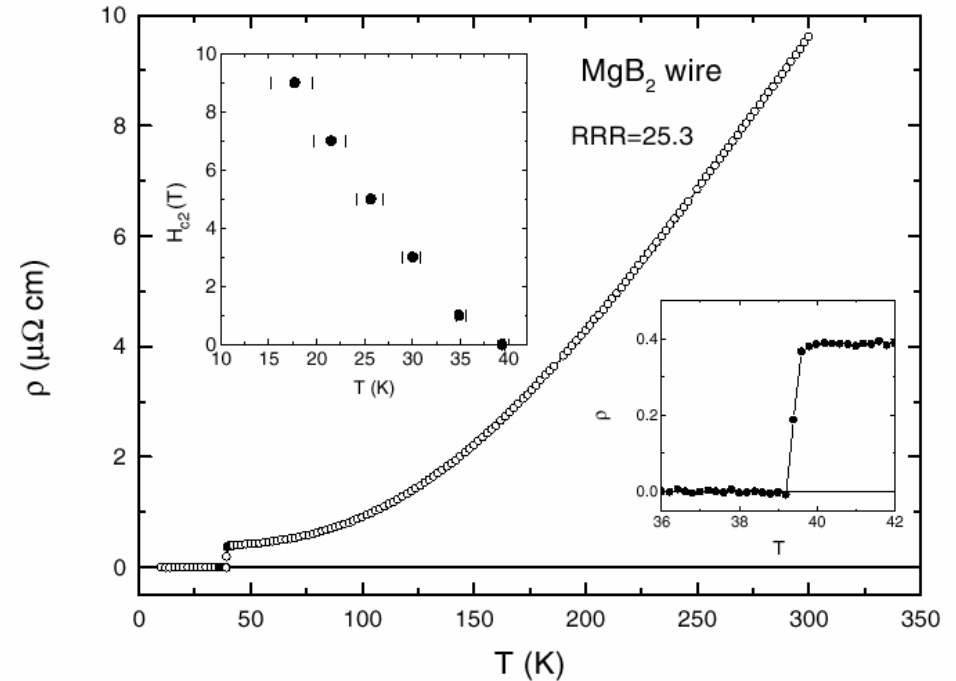
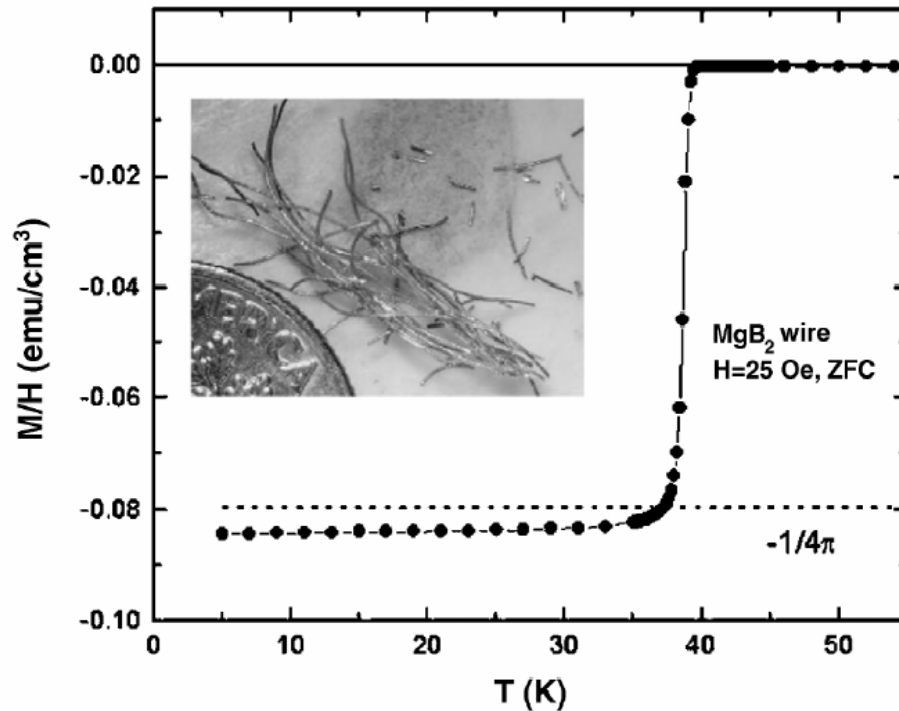
So,  $\text{MgB}_2$  is a low  $N(E_F)$  superconductor

For searches for  $\text{MgB}_2$  like compounds we should look for large  $\omega_D$  and large  $V$  and not obsess about  $N(E_F)$  so much....

*This is a much harder search algorithm.*



# Superconducting Properties



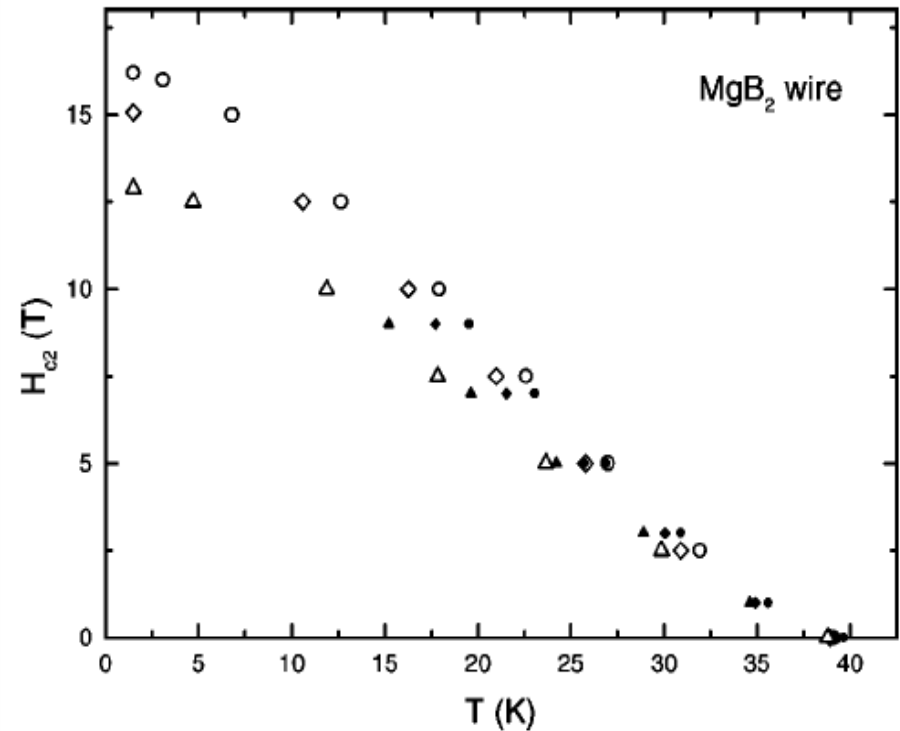
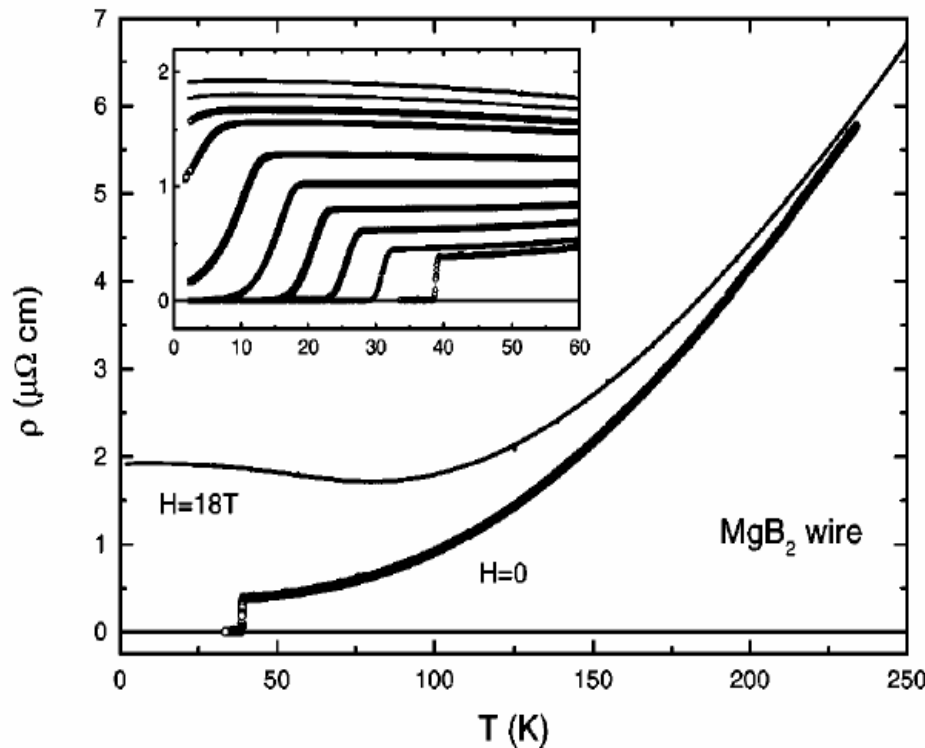
Canfield et al., PRL 86 (2001) 2423

- Over 90% dense => good grain to grain coupling.
- $T_c$  above 39K
- Full diamagnetic screening.
- Low normal state resistivity.



# $H_{c2}$ of $MgB_2$

Using resistivity data we can determine  $H_{c2}$  of the polycrystalline sample.

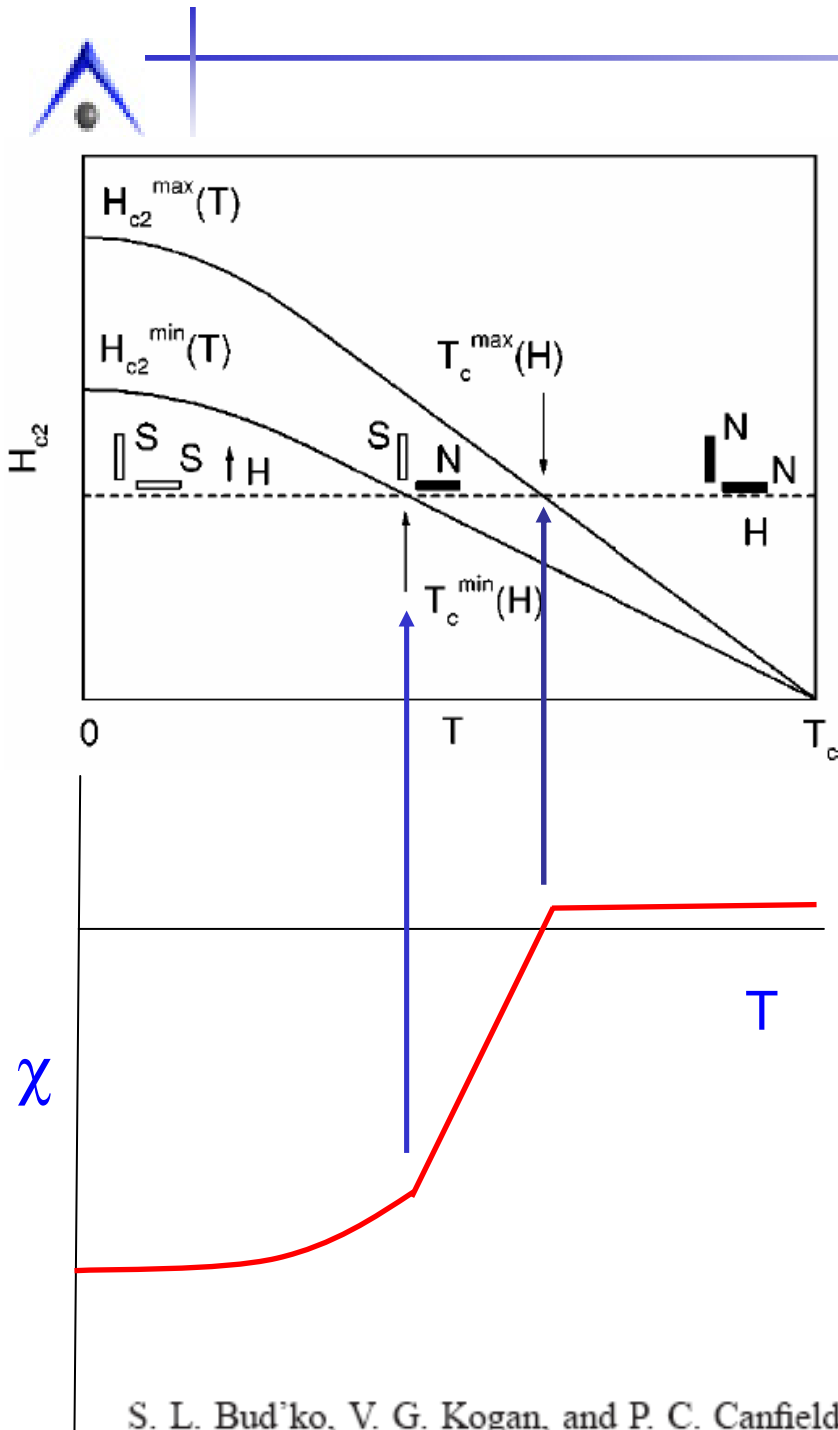


S.L. Bud'ko et al., Phys. Rev. B 63 (2001) 220503

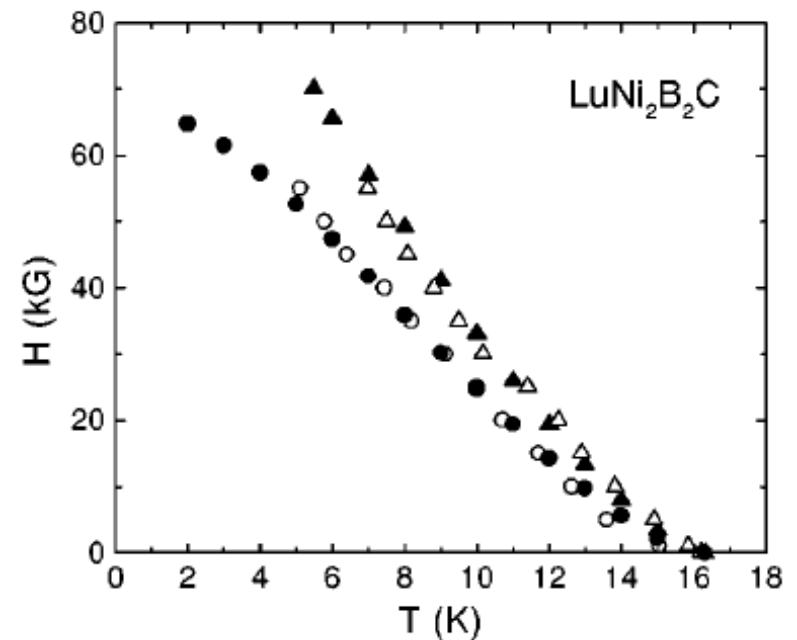
Two questions arise:

Are we missing hidden anisotropies?

Can we improve  $H_{c2}$  to be better than the  $\sim 30$  T of  $Nb_3Sn$ ?

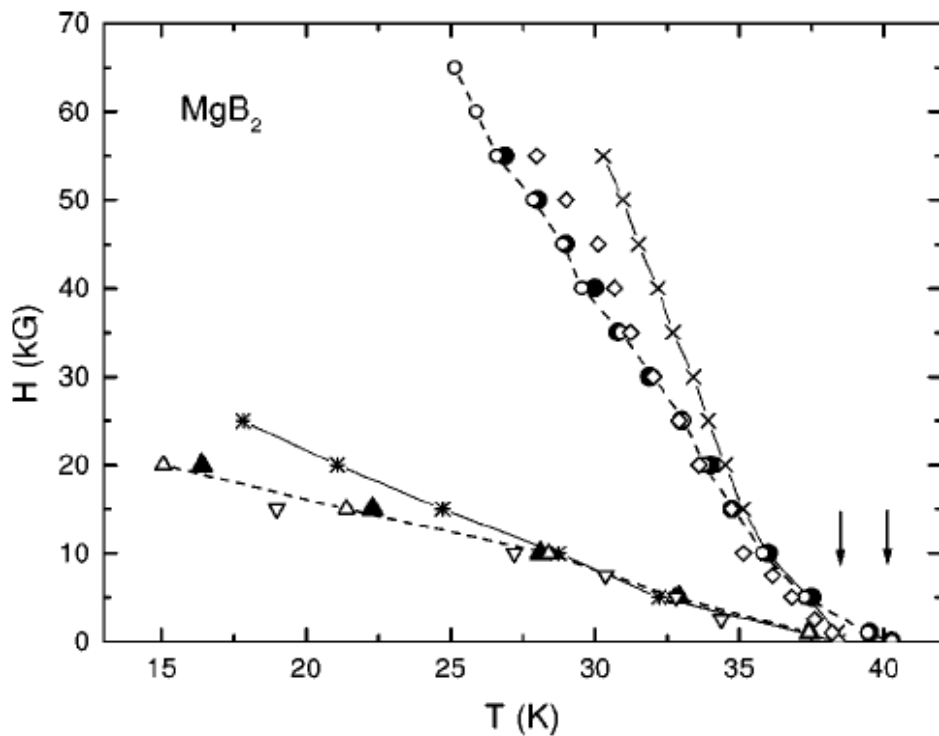
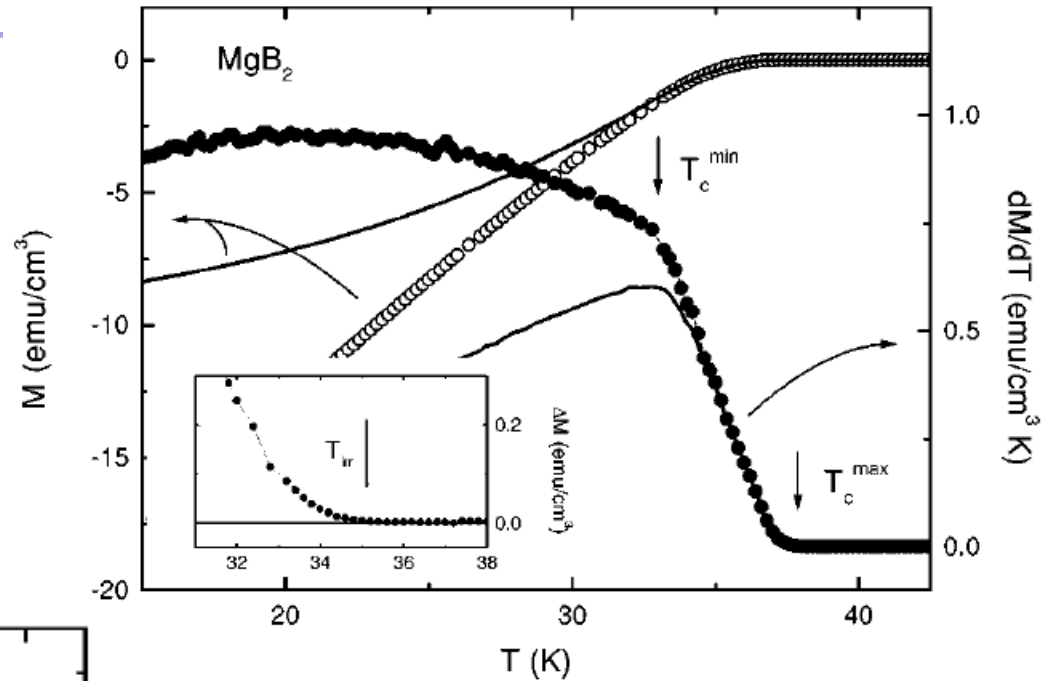


Given our experience with single crystals and anisotropy we were motivated to develop a new method for inferring  $H_{c2}$  anisotropy from polycrystalline data. We tested it on  $\text{YNi}_2\text{B}_2\text{C}$  and  $\text{LuNi}_2\text{B}_2\text{C}$  since we could measure directly on these samples as well.





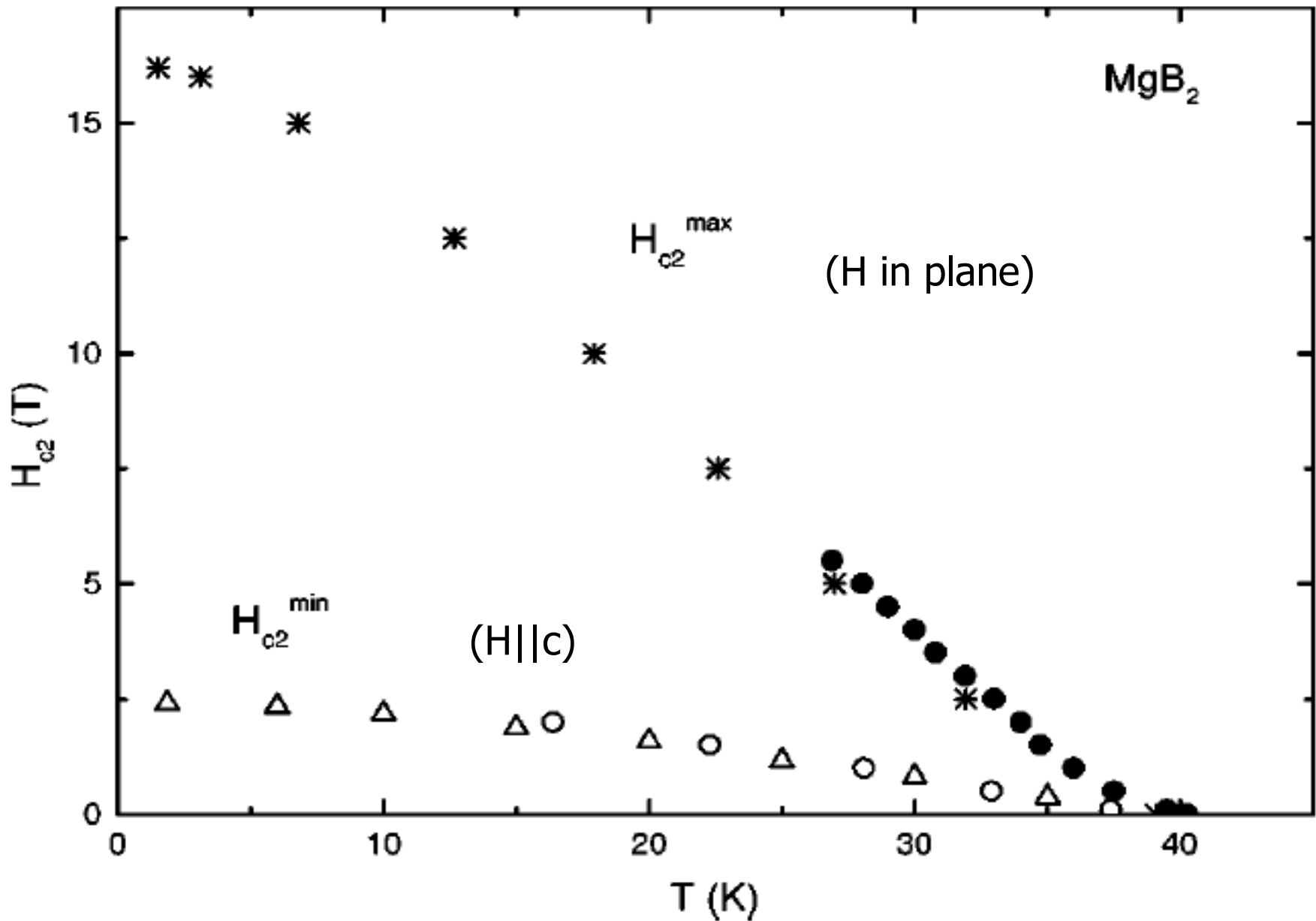
We used this method to deduce that  $\text{MgB}_2$  had an exceptionally large anisotropy in  $H_{c2}$ .



Anisotropic  $H_{c2}(T)$  data inferred from polycrystalline measurements!



The full  $H_{c2}$  anisotropy as well as direction could be determined by adding in resistivity data.



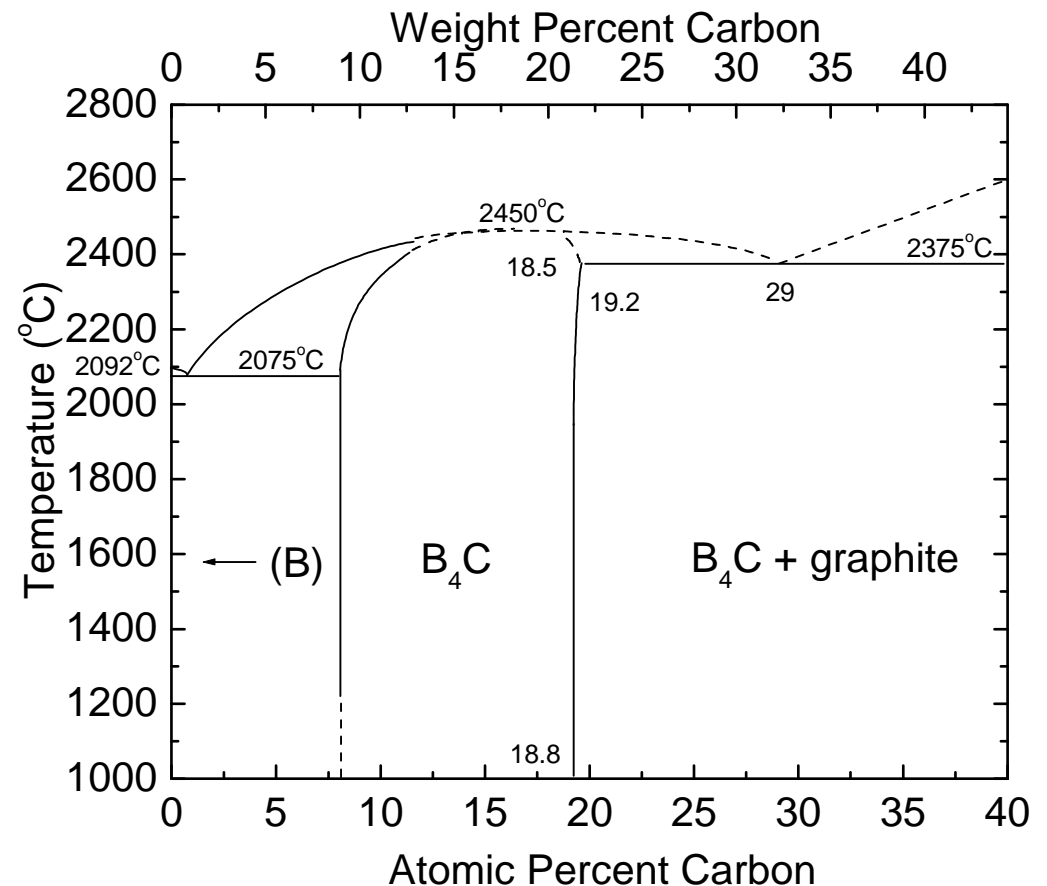


The other  $H_{c2}$  related question was, can we improve it? This can sometimes be done by adding non-magnetic impurities....Which one and how?

Recall we make  $MgB_2$  by diffusing the Mg vapor into B. Any viable dopant will need to be compatible with this synthesis technique.

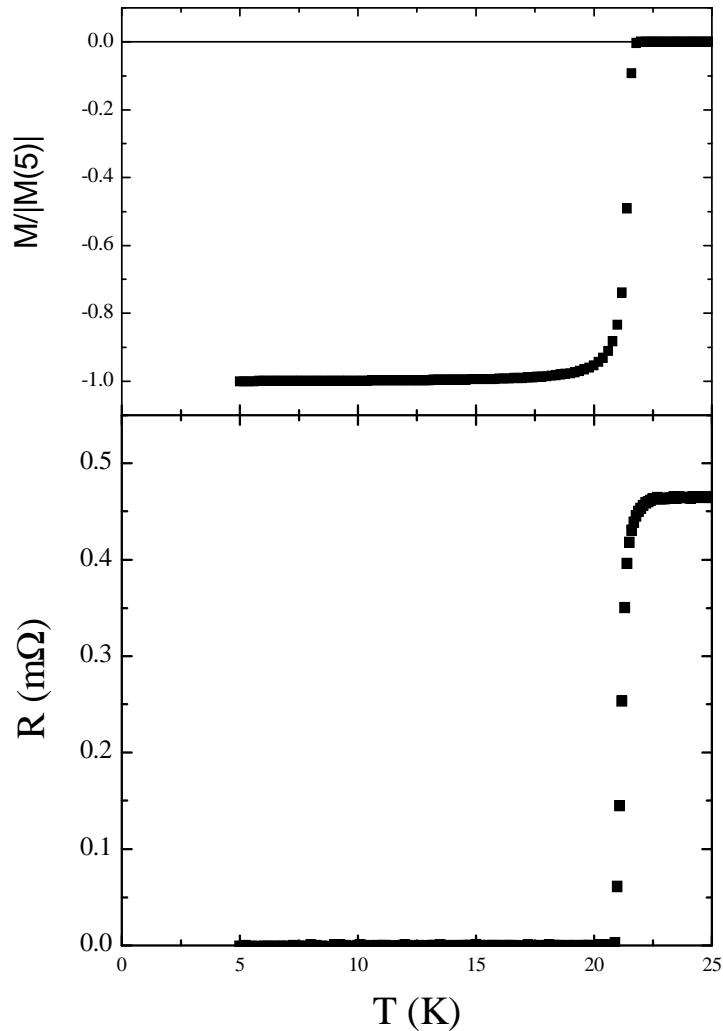
- $B_4C$  only stable Binary
- Width of formation allows for ~ 8-19% C doping

## B/C Phase Diagram





# Ribeiro et al.: $5\text{Mg} + 2\text{B}_4\text{C} \rightarrow \text{Mg}(\text{B}_{1-x}\text{C}_x)_2$

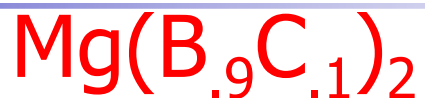


Ribeiro et al. Physica C 384 (2003) 227.

- Sharp superconducting phase transition with  $T_c \sim 22\text{K}$ .
- X-ray spectra indicate presence of  $\text{MgB}_2\text{C}_2$
- $x$  in  $\text{Mg}(\text{B}_{1-x}\text{C}_x)_2$  determined to be  $0.10 \pm 0.02$  by Rietveld Analysis of neutron diffraction pattern on  $\text{Mg}({}^{11}\text{B}_{1-x}\text{C}_x)_2$ .

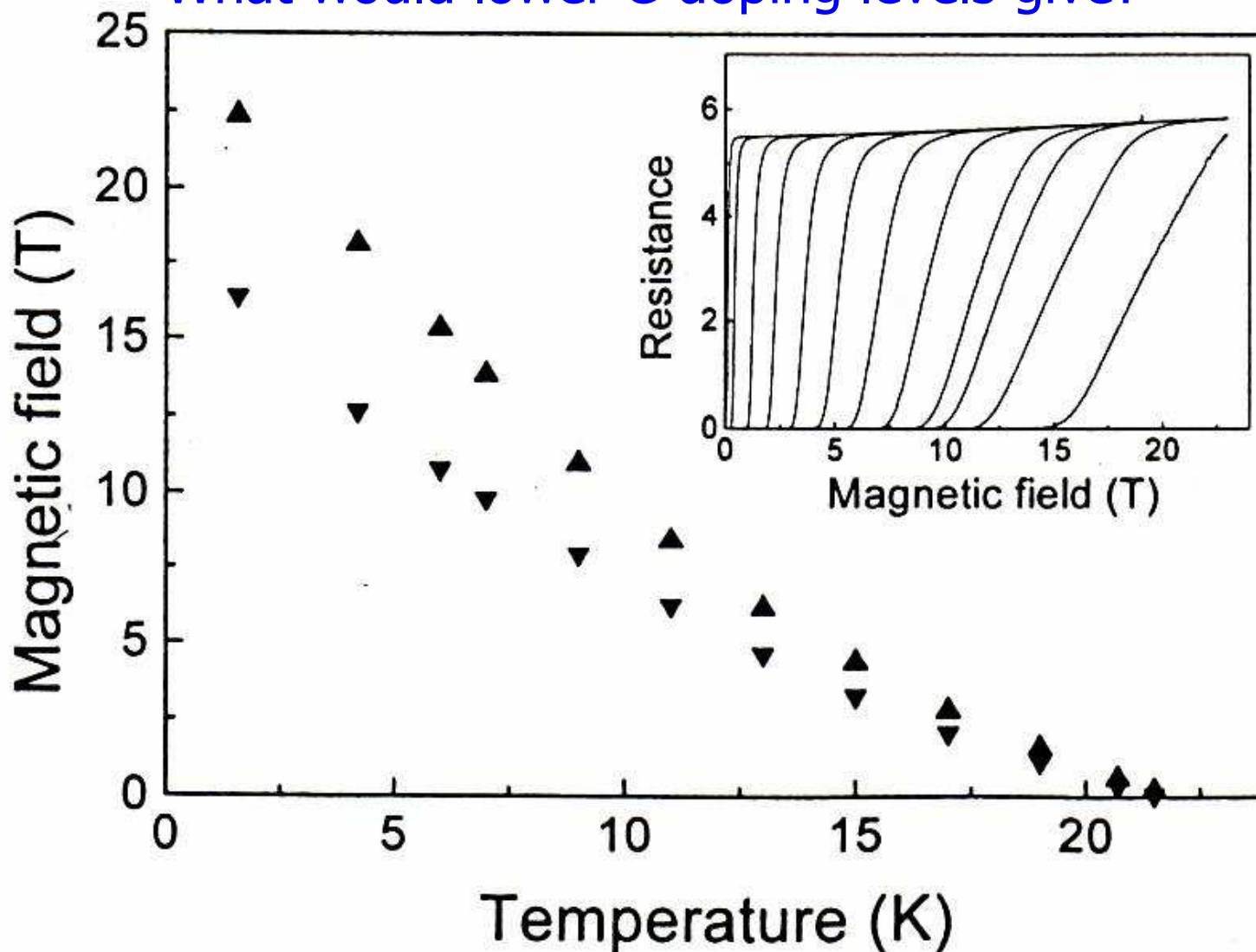
Phase	Weight fraction
$\text{Mg}(\text{B}_{1-x}\text{C}_x)_2$	73.4(1)
$\text{MgB}_2\text{C}_2$	20.4(1)
Mg	4.6(2)
MgO	1.6(1)

Avdeev et al. Physica C 387 (2003) 301



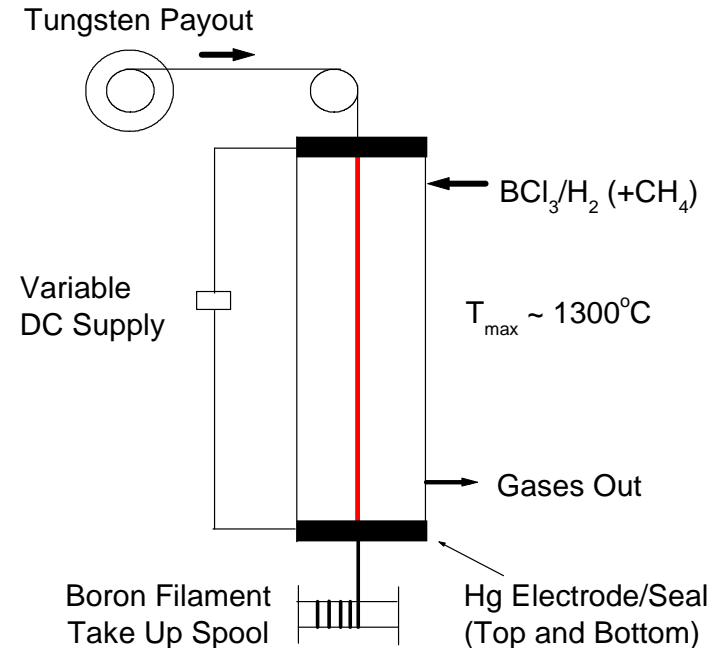
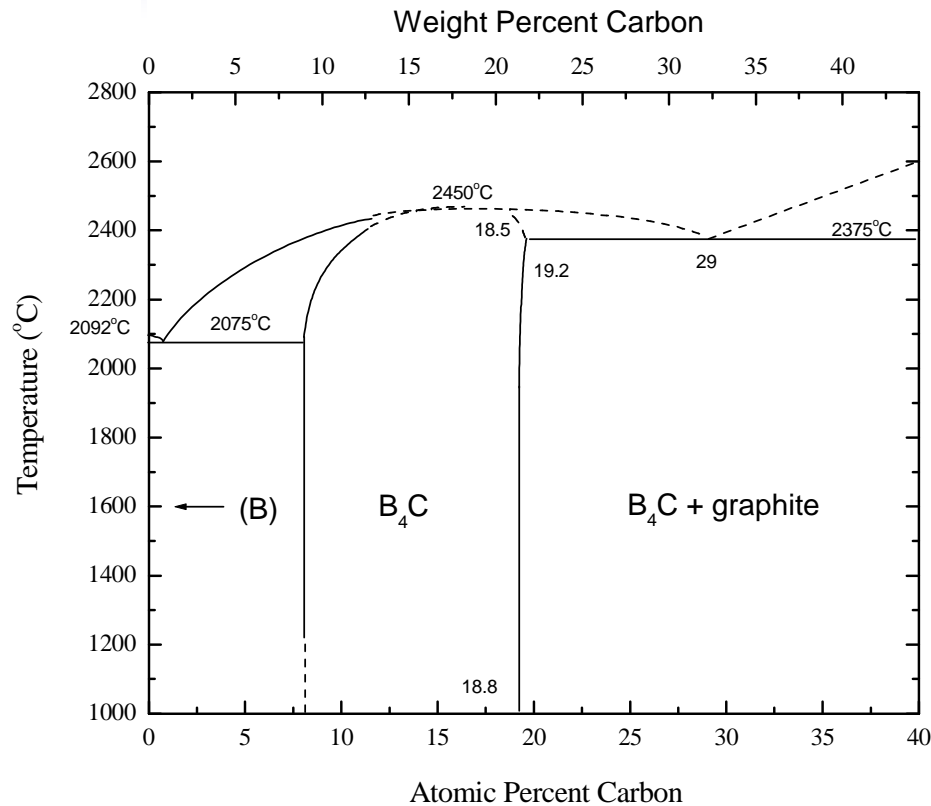
Even with  $T_c$  reduced to  $\sim 20$  K  $H_{c2}$  is clearly enhanced!!

What would lower C-doping levels give?





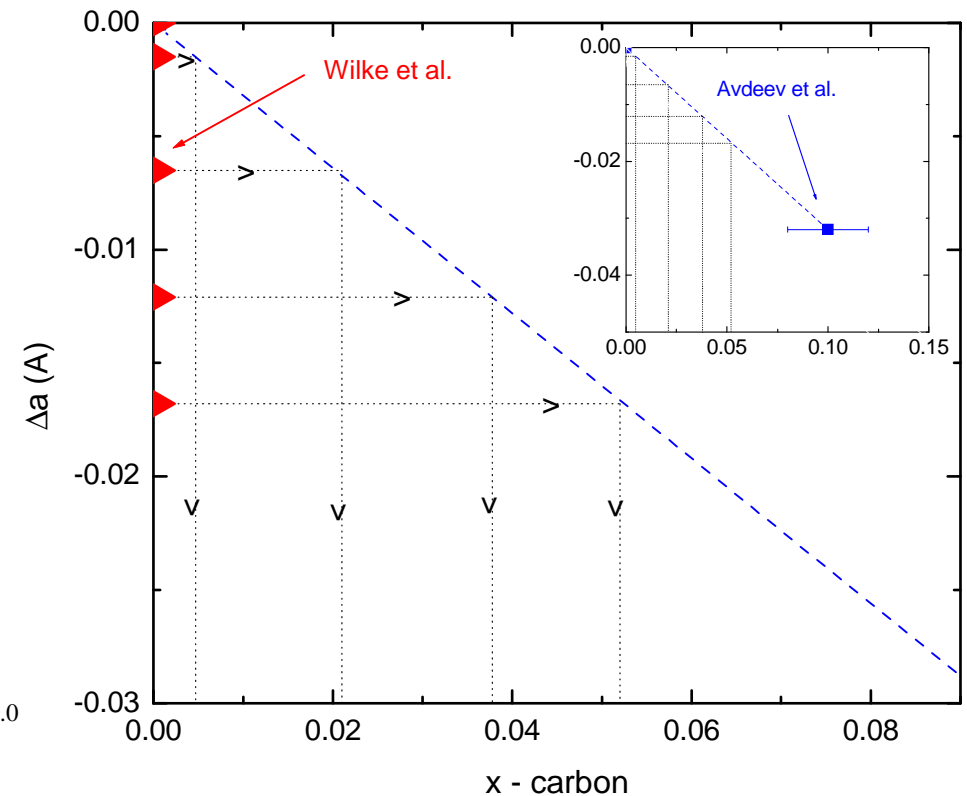
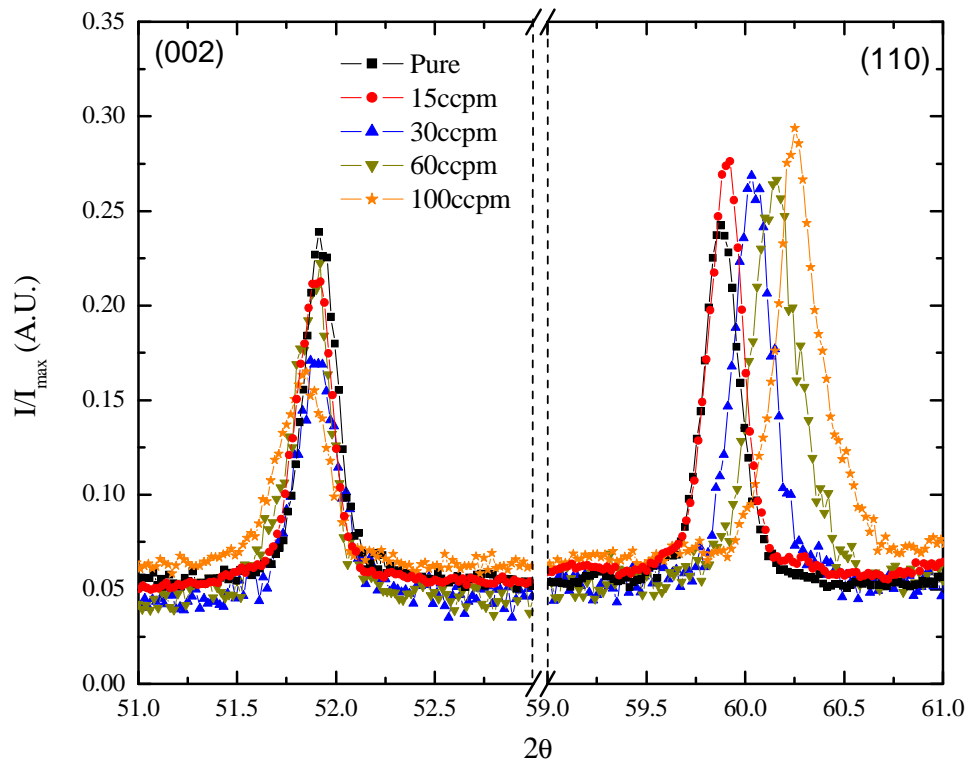
# Systematic Carbon Doping



- If  $Mg(B_{.9}C_{.1})_2$  has  $T_c \sim 22K$  and  $H_{c2}(T=0)$  near 25T,  $Mg(B_{1-x}C_x)_2$  may have maximum  $H_{c2}$  for  $x < 0.10$ .
- Width of formation in  $B_4C$  limits C doping to above 8%.
- Carbon incorporation: add  $CH_4$  to gas stream. Methane flow rates 15, 30, 60, and 100 ccpm (3000 ccpm  $BCl_3$ )



# Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> Filaments

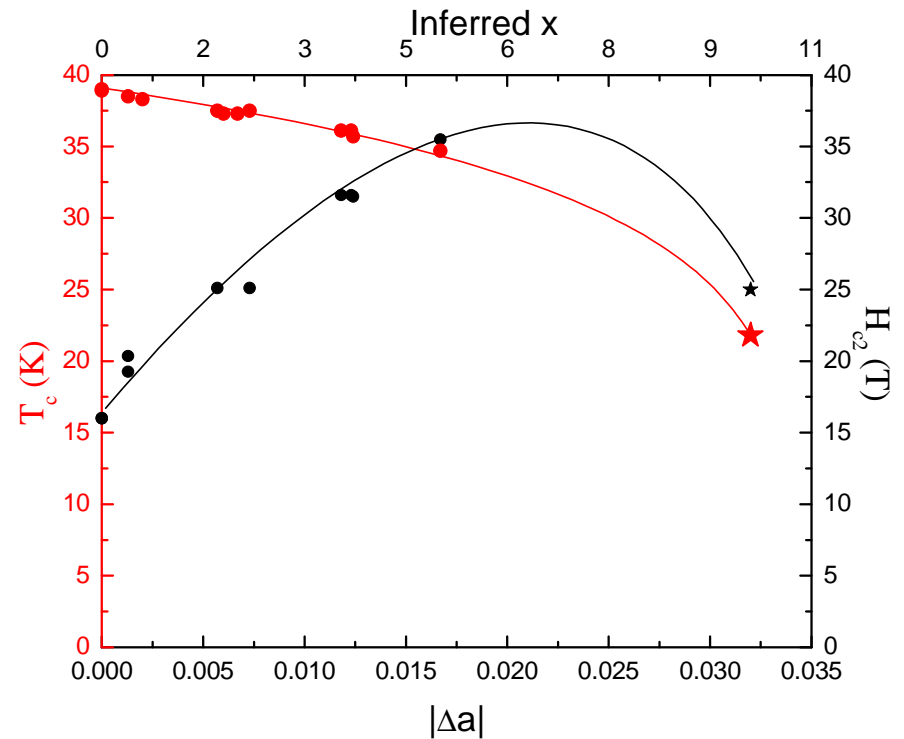
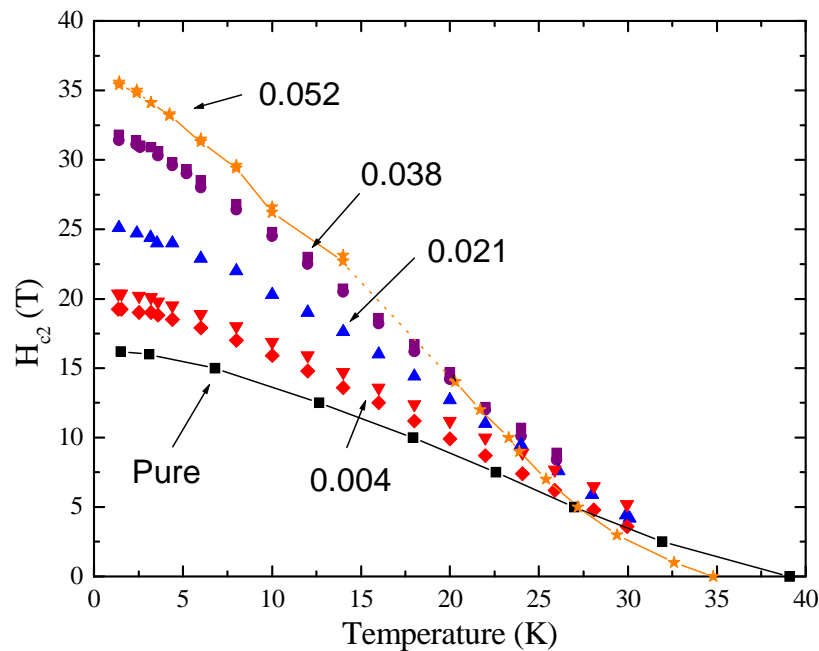
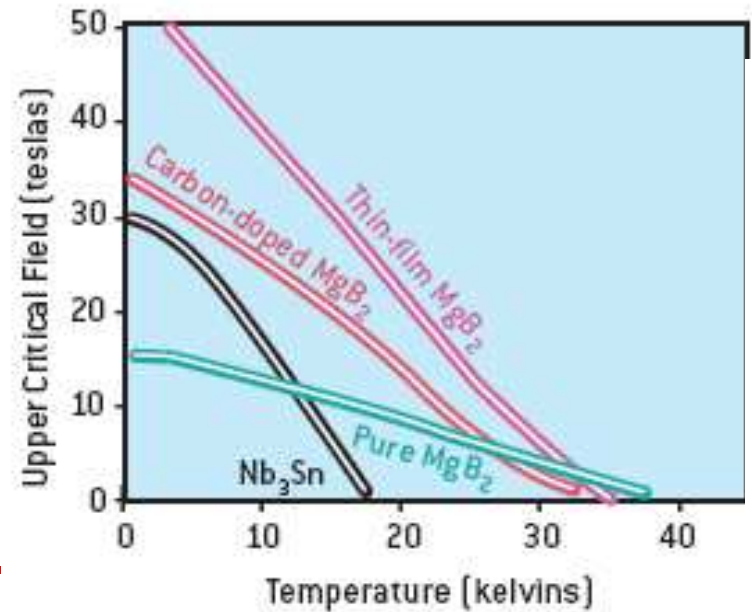


- Shift of (110) peak yields calculated carbon concentrations of  $x = 0.004, 0.021, 0.038,$  and  $0.052$ .



# Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> studies

Tuning of  $T_c$  and  $H_{c2}$  with light carbon doping makes MgB<sub>2</sub> superior to Nb<sub>3</sub>Sn (in H-T space).

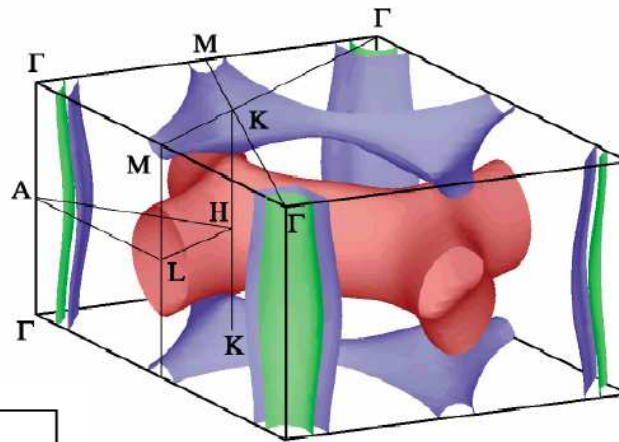
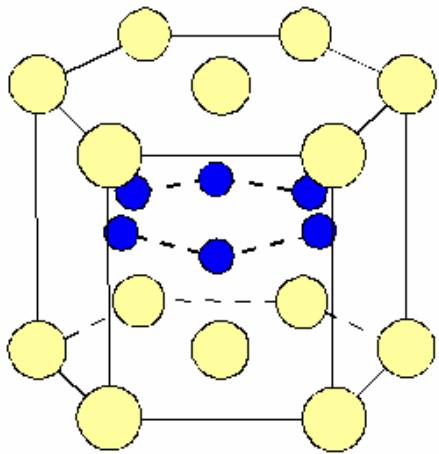


Phys. Rev. Lett. 92, 217003 (2004)

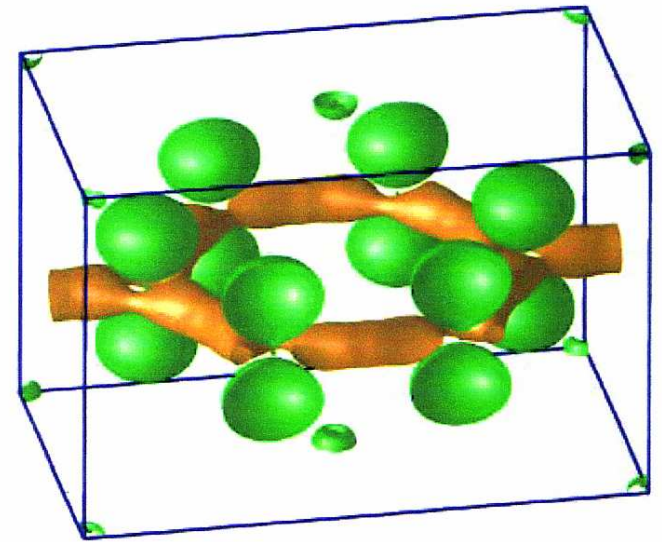
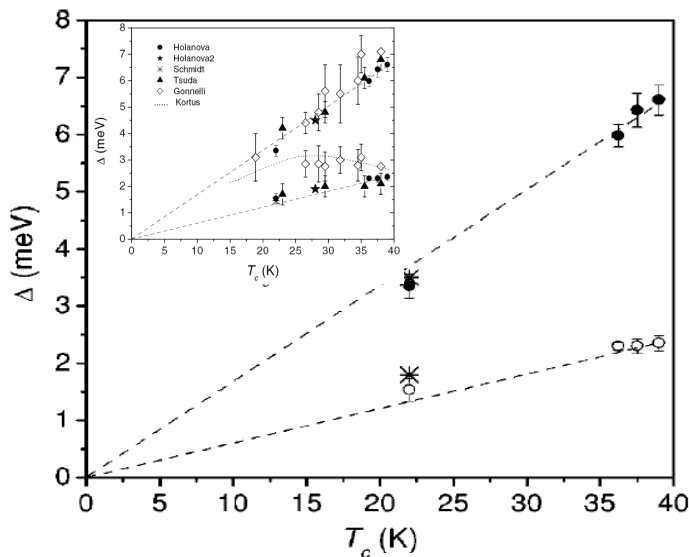


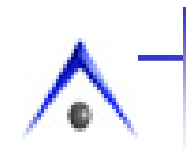
There is much, much more to  $\text{MgB}_2$  than just  $H_{c2}(T)$  and its tuning. There is the exceptionally clear two gap nature of the superconductivity as well as issues about  $J_c$  and making a practical, applied material.

## Origins of $\sigma$ and $\pi$ Bands

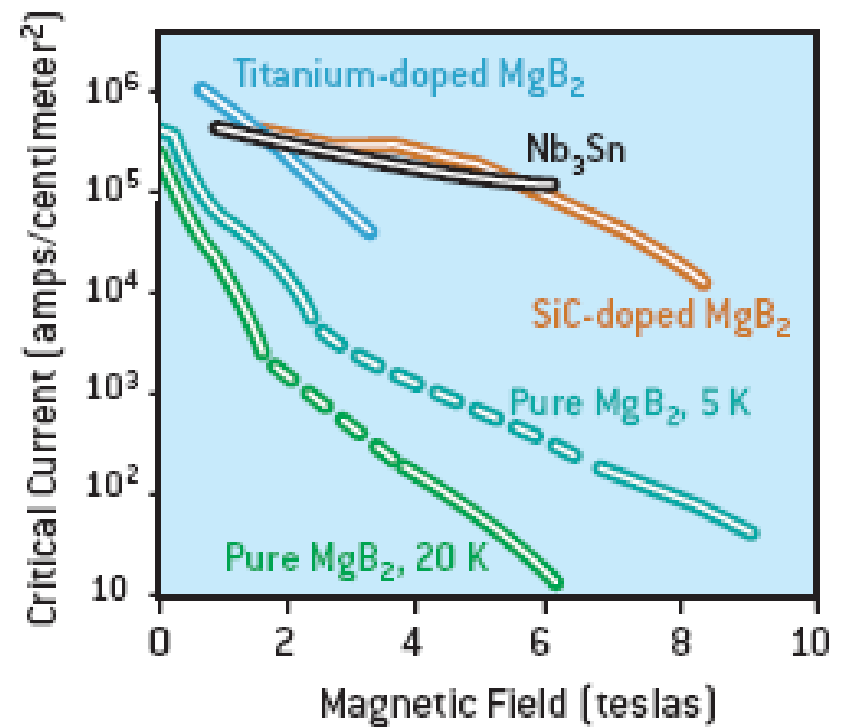
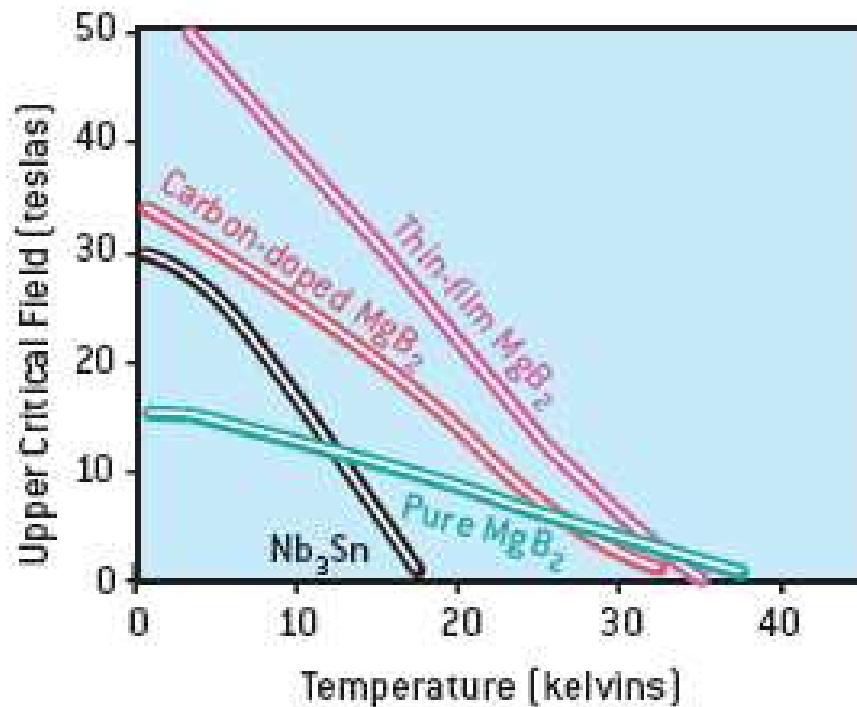


Korus et. al. Phys. Rev. Lett. 86 (2001) 4656





In addition, the current focus of applied research on  $\text{MgB}_2$  is how to improve and preserve a good critical current,  $J_c$ , especially for high  $H_{c2}$  samples. This too goes beyond an introductory lecture and into advanced topics in superconductivity.



# LOW-TEMPERATURE SUPERCONDUCTIVITY IS WARMING UP

Magnesium diboride defies the once conventional wisdom about what makes a good superconductor. It becomes superconducting near the relatively warm temperature of 40 kelvins—which promises a variety of applications

By **PAUL C. CANFIELD AND SERGEY L. BUD'KO**

*NOTE:*

*This article is available on our web page in nine different languages.*



To end this lecture, here is a final thought related to the search for new materials and ground states:

It is important to note that superconductivity in **both**  $\text{RNi}_2\text{B}_2\text{C}$  as well as  $\text{MgB}_2$  was discovered by accident (as part of a search for other compounds), illustrating perhaps one of the most important aspects of new materials research: the importance of keeping our eyes open for new phases / ground states.

This is not a new idea:



To end this lecture, here is a final thought related to the search for new materials and ground states:

It is important to note that superconductivity in **both**  $\text{RNi}_2\text{B}_2\text{C}$  as well as  $\text{MgB}_2$  was discovered by accident (as part of a search for other compounds), illustrating perhaps one of the most important aspects of new materials research: the importance of keeping our eyes open for new phases / ground states.

This is not a new idea:

*Search and you shall find -- what is unsought goes undetected.*                      Sophocles



*That's All Folks*