

Stochastic thermodynamics

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thanks to:

- F. Berger, J. Mehl, T. Schmiedl and Th. Speck (theory)
- V. Blickle and C. Bechinger (expt's on colloidal particles)
- C. Tietz, S. Schuler and J. Wrachtrup (expt's on single atoms)

Review: U.S., Eur. Phys. J. B, 64 : 423-431, 2008.

LECTURE II

- Stochastic entropy
- Non-equilibrium steady states (NESSs)

- Recap: Paradigm for mechanical driving:



- Langevin dynamics: $\dot{x} = \mu F(x, \lambda) + \zeta$,
- Total force: $F(x, \lambda) = -\partial_x V(x, \lambda) + f(\lambda)$ depends on $[\lambda(\tau)]$
- Gaussian noise: $\langle \zeta(\tau) \zeta(\tau') \rangle = 2D \delta(\tau - \tau')$ with $D = k_B T \mu$
- Fokker-Planck equation:

$$\partial_\tau p(x, \tau) = -\partial_x j(x, \tau) = -\partial_x (\mu F(x, \tau) - D \partial_x) p(x, \tau)$$

- Stochastic entropy [U.S., PRL 95, 040602, 2005]

- Fokker-Planck equation

$$\partial_\tau p(x, \tau) = -\partial_x j(x, \tau) = -\partial_x (\mu F(x, \lambda) - D \partial_x) p(x, \tau)$$

- Common non-eq **ensemble** entropy

$$S(\tau) \equiv - \int dx p(x, \tau) \ln p(x, \tau)$$

- Stochastic entropy for a **single trajectory** $x(\tau)$

$$s(\tau) \equiv - \ln p(x(\tau), \tau) \quad \text{with} \quad \langle s(\tau) \rangle = S(\tau)$$

- equation of motion

$$\dot{s}(\tau) = \underbrace{-\frac{\partial_\tau p(x, \tau)}{p(x, \tau)} + \frac{j(x, \tau)}{D p(x, \tau)}}_{\dot{s}_{\text{tot}}} \dot{x} - \underbrace{\frac{\mu F(x, \lambda)}{D}}_{\dot{s}_{\text{m}}} \dot{x}.$$

- $\Delta s = \Delta s_{\text{tot}} - \Delta s_{\text{m}} (= \Delta q/T)$

- Integral fluctuation theorem for entropy

$$1 = \langle \exp[\underbrace{-\beta q[x(\tau)]}_{-\Delta s_m} + \ln p_1(x_t)/p_0(x_0)] \rangle$$

- for any (normalized) $p_1(x_t)$
- with $p_1(x_t) = p(x, t) = \exp[-s(t)]$

- $\langle \exp[-\Delta s_{\text{tot}}] \rangle = 1 \Rightarrow \langle \Delta s_{\text{tot}} \rangle \geq 0$

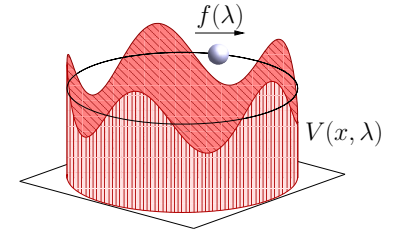
- integral fluctuation theorem for total entropy production
- arbitrary initial state, driving, length of trajectory

- Unification through a general integral fluctuation theorem

$$\langle \exp[-\Delta s_m] p_1(x_t)/p_0(x_0) \rangle = 1$$

for any (normalized) $p_1(x_t)$

$p_1(x)$	Integral F'theorem	
$p(x, t)$	$\langle \exp[-\Delta s_{\text{tot}}] \rangle = 1$	U.S. '05
$p_{\text{eq}}(x, \lambda_t)$	$\langle \exp[-W] \rangle = \exp[-\Delta F]$	Jarzynski '97
$p_{\text{eq}}(x, \lambda_0)$	$\langle \exp[-W_0] \rangle = 1$	Bochkov-Kuzovlev '81
$f(x)p_{\text{eq}}(x, \lambda_t)/\mathcal{N}$...	$\langle f(x_t) \exp[-W_d] \rangle = \langle f(x) \rangle_{\text{eq}(\lambda_t)}$	Crooks '00 ...



- Non-equilibrium steady state (NESS)

- $\dot{x} = \mu[-\partial_x V(x) + f] + \zeta \equiv \mu F(x) + \zeta,$

- stat distribution $p^s(x) \equiv \exp[-\phi(x)]$ with $\phi(x) \neq V(x)/T$

- broken detailed balance $p(x_2(t)|x_1(0))p^s(x_1) \neq p(x_1(t)|x_2(0))p^s(x_2)$

- mean global velocity $v_s = \langle \nu_s(x) \rangle = j_s$

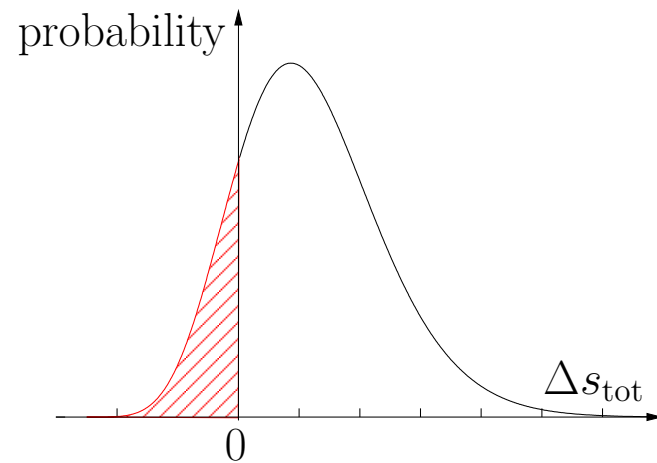
- mean local velocity $\nu_s(x) \equiv \langle \dot{x}|x \rangle = j_s/p^s(x) = \mu F(x) + D\phi'(x)$

- NESS I: Fluct'n theorem

$$p(-\Delta s_{\text{tot}})/p(\Delta s_{\text{tot}}) = \exp(-\Delta s_{\text{tot}})$$

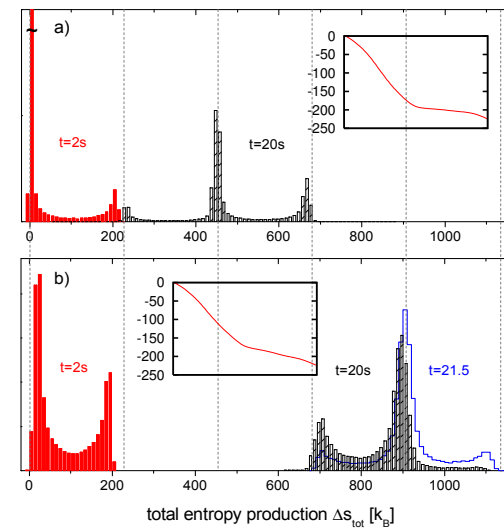
long-time limit: Evans et al (1993), Gallavotti & Cohen (1995), Kurchan (1998), Lebowitz & Spohn (1999) ... **finite times:** U.S., PRL'05

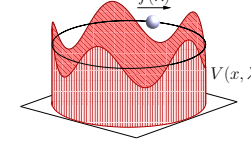
– violation of the 2nd law ??



– experimental data

[Speck, Blickle, Bechinger, U.S., EPL **79** 30002 (2007)]

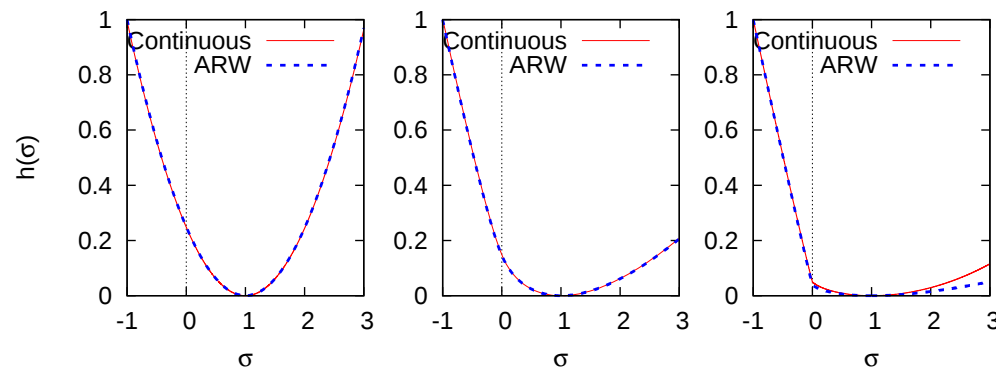




- Long time: Large deviation form

[J. Mehl, T. Speck, U.S., Phys. Rev. E, 78 : 011123, 2008.]

- Scaled entropy production rate $\sigma \equiv \Delta s_m / \langle \dot{s}_m \rangle t$
- “Large deviation” for $t \rightarrow \infty$ $p(\Delta s_m, t) \sim \exp[-t h(\sigma)]$
- Rate function $h(\sigma) = h(-\sigma) - \sigma$



- mapping to asymm random walk

- Proof of finite time DFT [U.S., PRL 95, 040602, 2005]

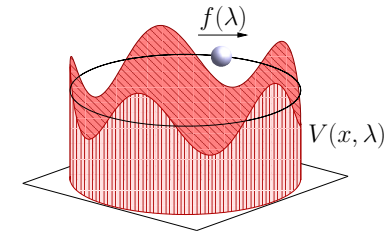
- steady state distribution $p^s(x)$
- total entropy production $\Delta s_{\text{tot}} = \ln \frac{p[x(\tau)|x_0] p^s(x_0)}{p[\tilde{x}(\tau)|\tilde{x}_0] p^s(\tilde{x}_0)}$
- for any function $g(\Delta s_{\text{tot}})$

$$\begin{aligned}
 \langle g(\Delta s_{\text{tot}}) \rangle &= \sum_{x(\tau), n_0} p[x(\tau)|x_0] p^s(x_0) g(\Delta s_{\text{tot}}) \\
 &= \sum_{x(\tau), x_0} p[\tilde{x}(\tau)|\tilde{x}_0] p^s(\tilde{x}_0) \exp \Delta s_{\text{tot}} g(\Delta s_{\text{tot}}) \\
 &= \sum_{\tilde{x}(\tau), \tilde{x}_0} p[\tilde{x}(\tau)|\tilde{x}_0] p^s(\tilde{x}_0) \exp \Delta s_{\text{tot}} g(\Delta s_{\text{tot}}) \\
 &= \sum_{x(\tau), n_0} p[x(\tau)|x_0] p^s(x_0) \exp(-\Delta s_{\text{tot}}) g(-\Delta s_{\text{tot}}) \\
 &= \langle \exp(-\Delta s_{\text{tot}}) g(-\Delta s_{\text{tot}}) \rangle
 \end{aligned}$$

$$\text{– } \boxed{p(-\Delta s_{\text{tot}})/p(\Delta s_{\text{tot}}) = \exp(-\Delta s_{\text{tot}})}$$

- NESS II: FDT violation and restoration

[T. Speck and U.S., Europhys. Lett. **74**, 391, 2006]



– FDT in eq:

$$T \frac{\delta \langle \dot{x}(t) \rangle}{\delta f(\tau)} \Big|_{f=0} = \langle \dot{x}(t) \dot{x}(\tau) \rangle_{\text{eq}}$$

– violation of FDT in non-eq:

$$T \frac{\delta \langle \dot{x}(t) \rangle}{\delta f(\tau)} \Big|_{f \neq 0} = \langle \dot{x}(t) \dot{x}(\tau) \rangle_{\text{neq}} - \langle \dot{x}(t) \nu_s(x(\tau)) \rangle_{\text{neq}}$$

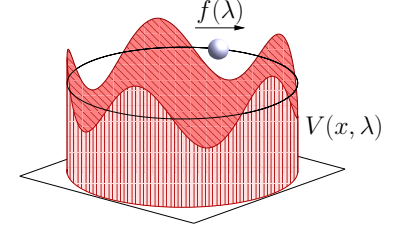
– restoration of FDT in non-eq for renormalized velocity:

$$v(\tau) \equiv \dot{x}(\tau) - \nu_s(x(\tau))$$

$$T \frac{\delta \langle v(t) \rangle}{\delta f(\tau)} \Big|_{f \neq 0} = \langle v(t) v(\tau) \rangle_{\text{neq}}$$

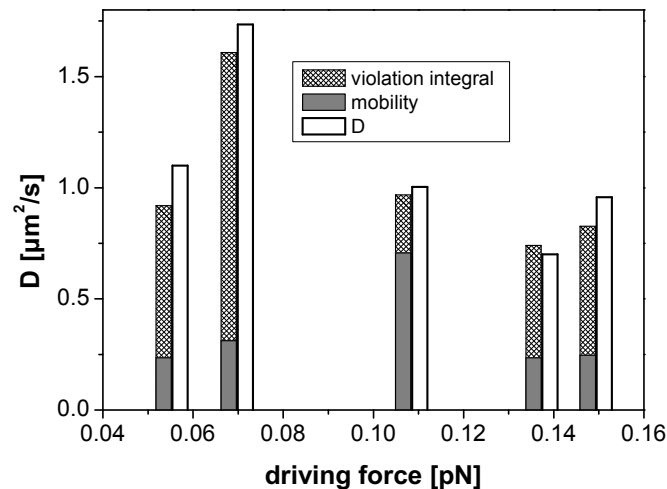
- NESS III: Generalized Einstein relation

[Blickle, Speck, Lutz, U.S., Bechinger, PRL **98**, 210601 (2007)]



$$\frac{\delta \langle \dot{x}(t) \rangle}{\delta f(\tau)} \Big|_{f \neq 0} = \langle \dot{x}(t) \dot{x}(\tau) \rangle_{\text{neq}} - \langle \dot{x}(t) \nu_s(x(\tau)) \rangle_{\text{neq}}$$

$$\mu_{\text{eff}}(f) = D_{\text{eff}}(f) - \int_0^\infty dt I(t) \quad \text{with} \quad I(t) \equiv \langle \dot{x}(t) \nu_s(x(0)) \rangle - \langle \nu_s(x) \rangle^2$$



cf giant diffusion [Reimann et al 2001]

$$\neq D_{\text{eff}}(f) = T_{\text{eff}}(f) \mu_{\text{eff}}(f)$$

- NESS IV: : Entropy production for dumbbell in shear flow

[T. Speck, J. Mehl and U.S., PRL **100** 178302, 2008]

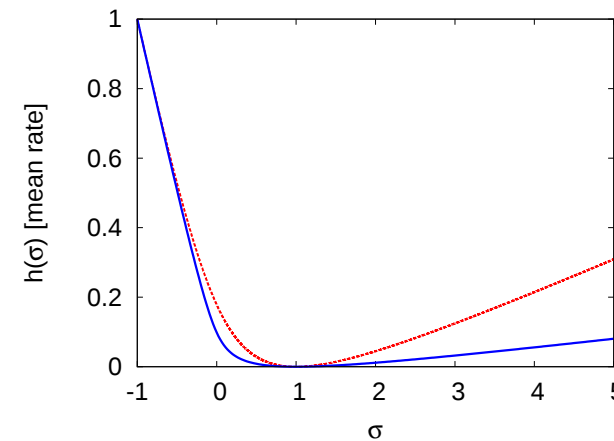
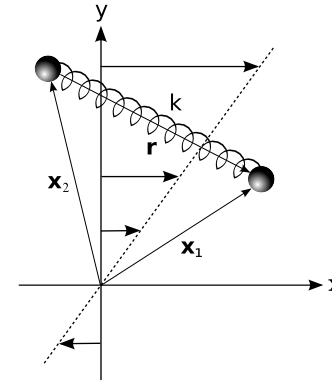
- Scaled entropy production rate

$$\sigma \equiv \Delta s_m / \langle \dot{s}_m \rangle t$$

- “Large deviation” for $t \rightarrow \infty$

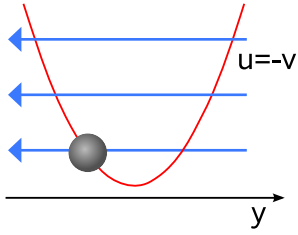
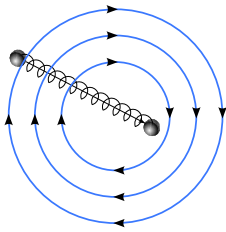
$$p(\Delta s_m, t) \sim \exp[-t h(\sigma)]$$

- Rate function $h(\sigma)$



[cf. Turitsyn et al, PRL 2007]

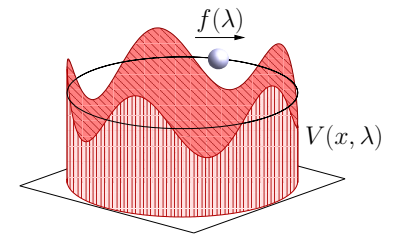
- Equilibrium vs. non-equilibrium for flow driven systems

system	det bal	prob currents	
 particle trapped in flow	fulfilled	zero	non-eq
 dumbbell in pure rot flow	broken	non-zero	eq

- better criterion for eq:

$$dw = 0 = D_t V + \mathbf{f}[d\mathbf{r} - \mathbf{u}dt]$$

- Transitions between NESSs (SST) [Oono and Paniconi, Hatano and Sasa]
 - $V(x)$ time-independent, $f = f(\lambda(\tau))$ switches from f_1 to f_2
 - $\phi(x, \lambda) \equiv -\ln p^s(x, \lambda)$ ($\neq s(\tau)$)

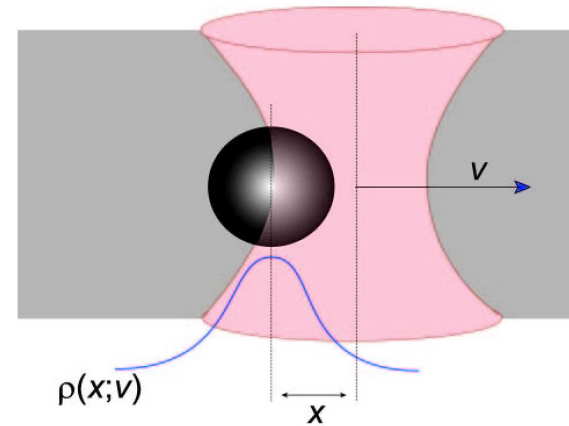
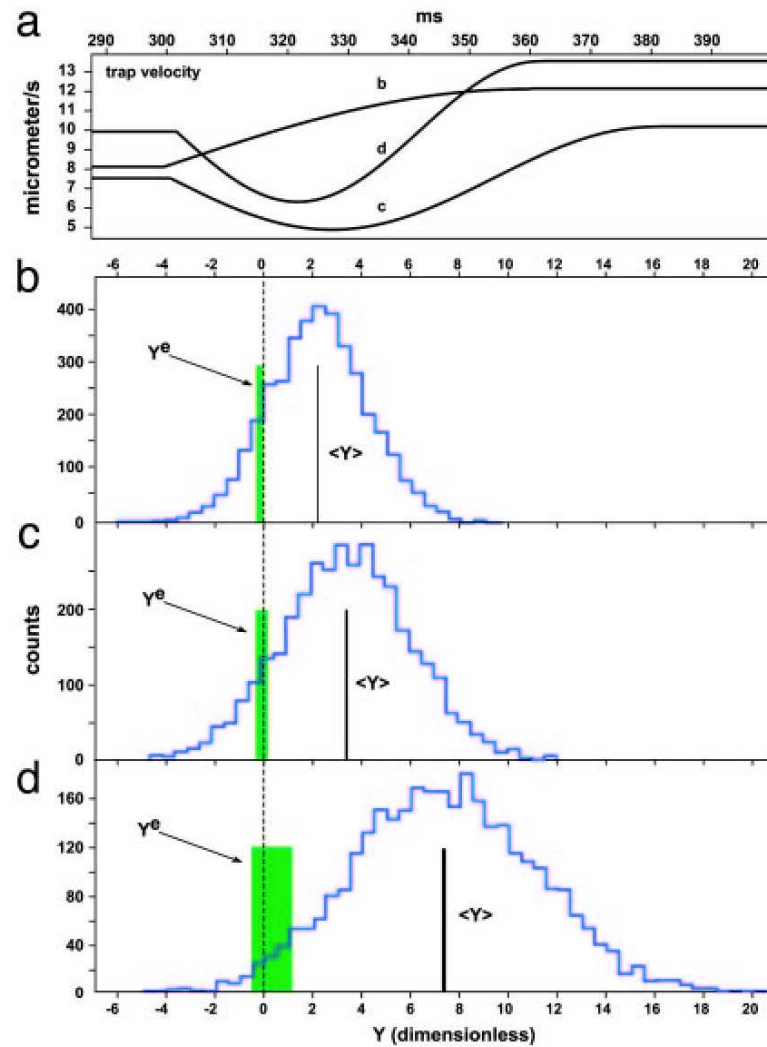


- total heat

$$\begin{aligned}
 q &= \int_0^t d\tau F(x, \lambda) \dot{x} = \Delta s_m \\
 &= \underbrace{\int_0^t \nu_s(x(\tau), \lambda) \dot{x}}_{q_{hk}} \underbrace{- D[(\phi(x_t, \lambda_t) - \phi(x_0, \lambda_0))] + D \int \dot{\lambda} (\partial \phi / \partial \lambda)}_{q_{ex}} \\
 &\quad \quad \quad \underbrace{\hspace{10em}}_{\Delta S} \quad \quad \quad \underbrace{\hspace{10em}}_{\equiv Y}
 \end{aligned}$$

- $\langle \exp[-Y] \rangle = 1 \Rightarrow \boxed{\Delta S \geq -\langle q_{ex} \rangle}$ (“2nd law for NESSs”)

- Experimental test of the Hatano-Sasa relation $\langle \exp -Y \rangle = 1$
[E.H. Trepagnier et al, PNAS 2004]



- Further FTs: (T. Speck, U.S, J Phys A 38, L581, 2005)

- $\langle \exp(-q_{hk}) \rangle = 1$

- $\langle \exp[-\underbrace{\Delta s_m + \Delta \phi}_R] \rangle = 1$ (generalized JR)

	General	DBS	NESS
		$Q_{hk} = 0, R = Y = \beta W_{\text{dis}}$	$Y = 0, R = \beta Q_{hk}$
(1)	$\langle \exp[-Y] \rangle = 1$	$\langle \exp[-\beta W_{\text{dis}}] \rangle = 1$	$1 = 1$
(2)	$\langle \exp[-\beta Q_{hk}] \rangle = 1$	$1 = 1$	$\langle \exp[-\beta Q_{hk}] \rangle = 1$
(3)	$\langle \exp[-R] \rangle = 1$	$\langle \exp[-\beta W_{\text{dis}}] \rangle = 1$	$\langle \exp[-\beta Q_{hk}] \rangle = 1$