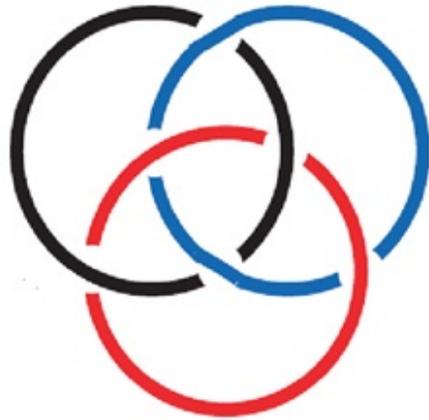
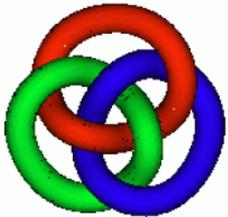


- **Expect tremendous experimental differences**
- **Mathematically extremely difficult to consider**
- **Option for computer simulations**





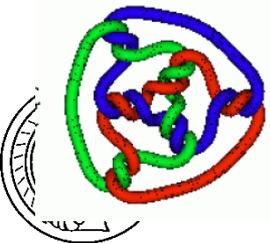
- Expect tremendous experimental differences
- Mathematically extremely difficult to consider
- Option for computer simulations

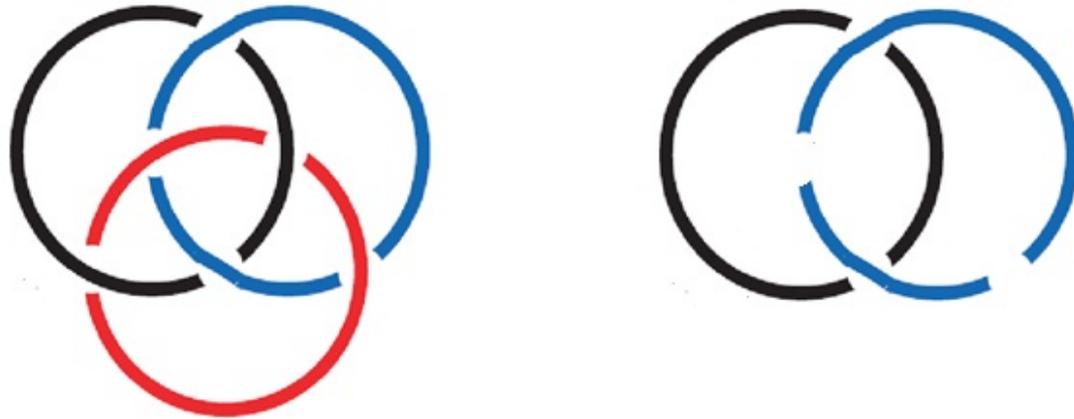


Borromean rings

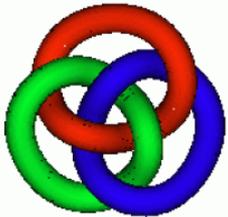
DNA rings, Nad Seeman, NYU

Door of an Italian church





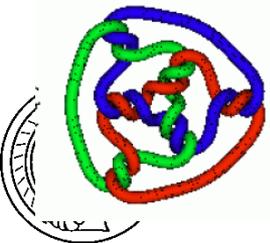
- Expect tremendous experimental differences
- Mathematically extremely difficult to consider
- Option for computer simulations



Borromean rings

DNA rings, Nad Seeman, NYU

Door of an Italian church



Topological Constraints Matter



- **Tube concept**

- R. Everaers, G. S. Grest, M. Pütz, ...

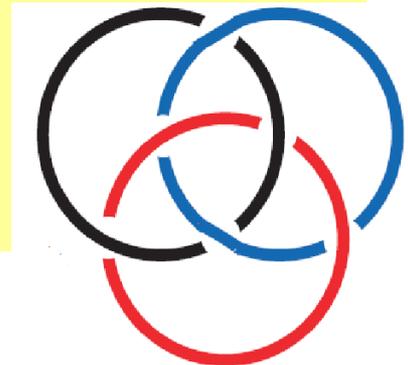
- Gels etc.

- R. Everaers, G. S. Grest, C. Holm, ...

- Melt of non concatenated rings, melt of globules

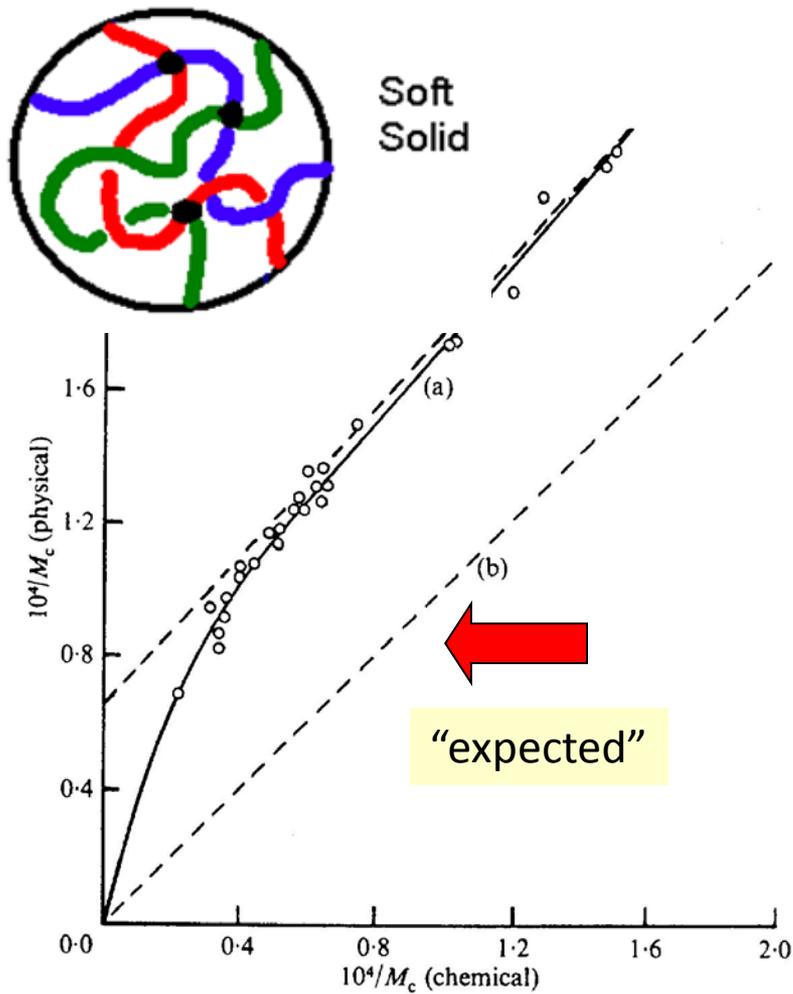
- T. Vettorell, J. Halverson, A. Yu Grosberg

- Questions

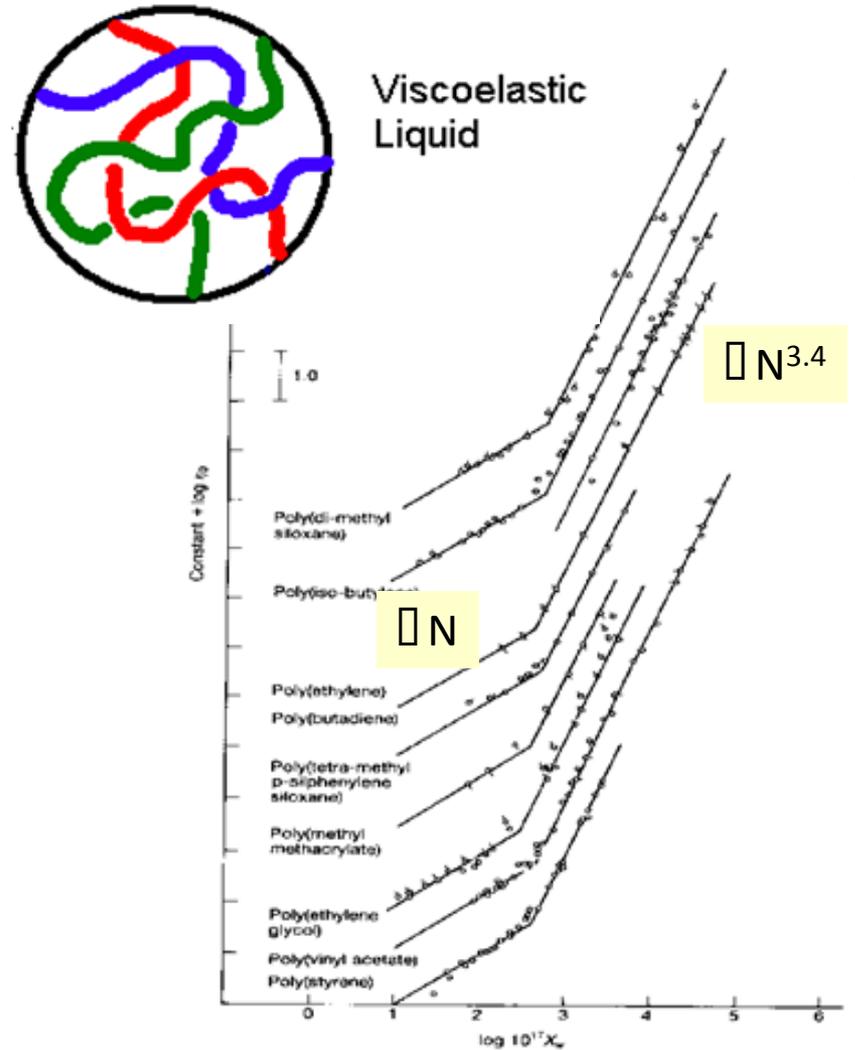




Elastic modulus of network vs $1/N$, N chain length



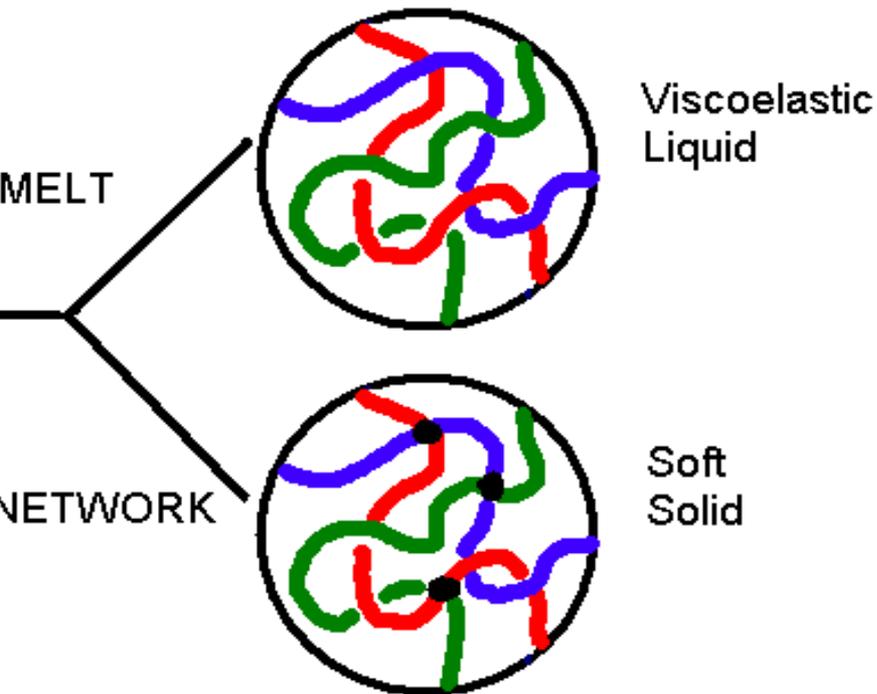
Melt viscosity vs N



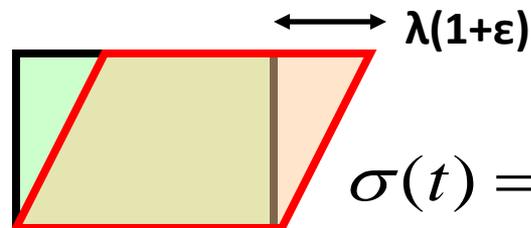
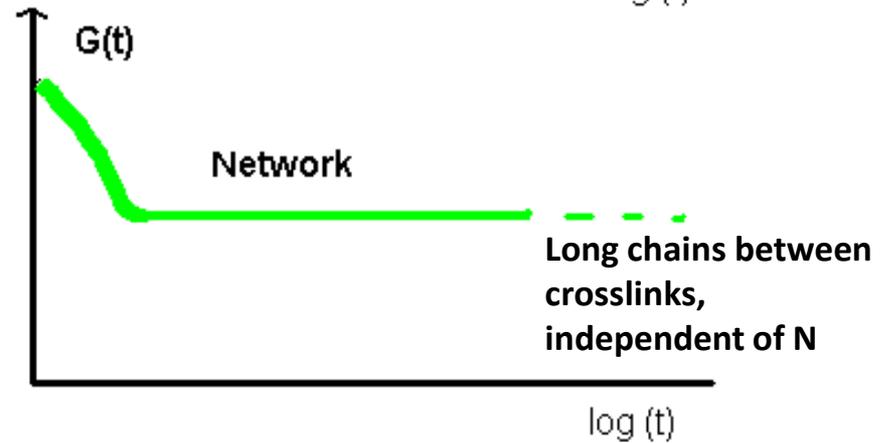
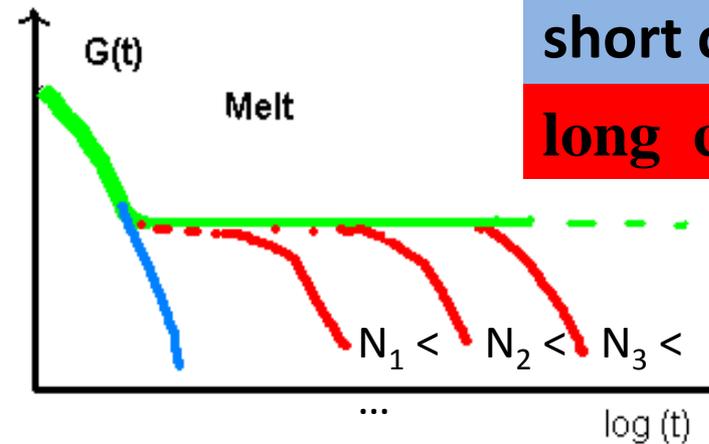
(Moore & Watson 1956
Mullins 1959)

From Doi and Edwards

Classical Problem: MELT \leftrightarrow NETWORK



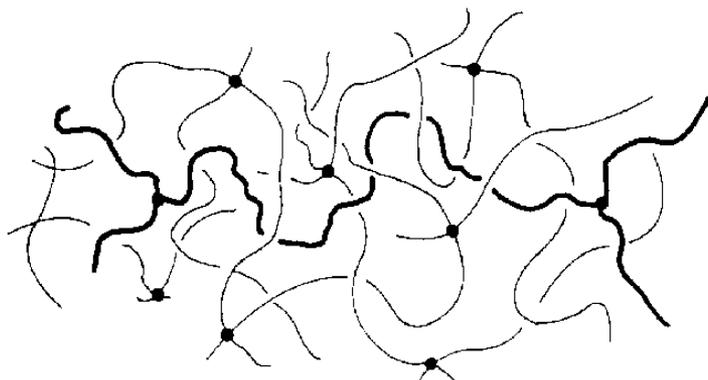
□ quenched



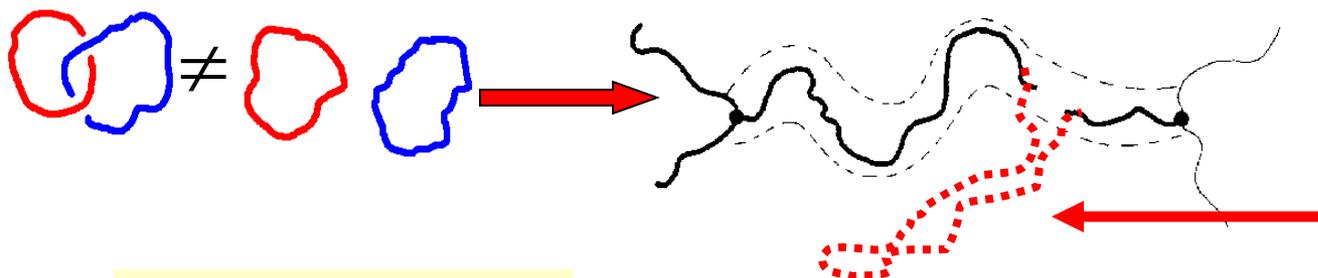
$$\sigma(t) = G(t)(\lambda - 1)$$

G(t): Time dependent (shear) modulus:

Tube Model



Edwards '67

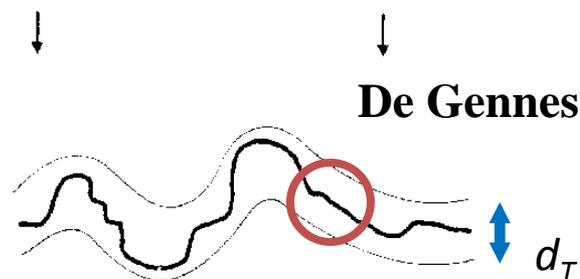


“Tube leakage”:
Entropy penalty ($O(L/2)$)
for defects too large

$$d_T \propto N_e^{1/2}$$

$$G^0 \propto N_e^{-1}$$

$$\tau_e \propto N_e^2$$



De Gennes '71

Unique length d_T



Tube Model

Concept, based entirely on conformational statistics of networks (soft solid) and melts (liquid!).

→ $d_T \propto N^{1/2}$
 N_e should be determined from conformational statistics of the polymers(?)

$$\tau_e \propto N_e$$

Unique length d_T

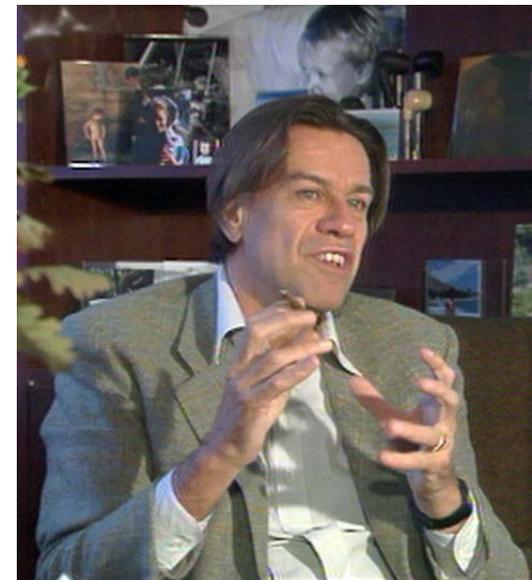
Edwards '67

De Gennes '71

d_T



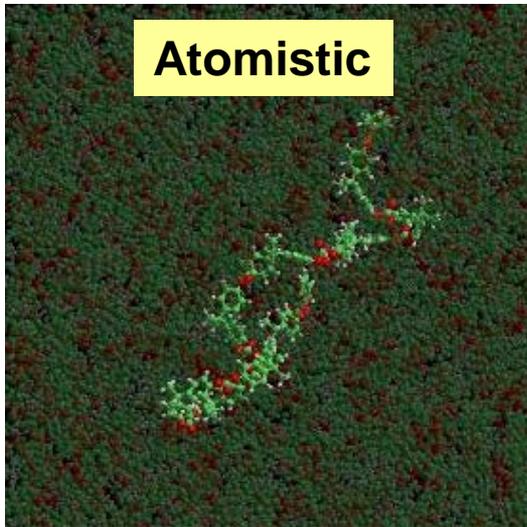
“Tube leakage”:
Entropy penalty ($O(L/2)$)
for defects too large



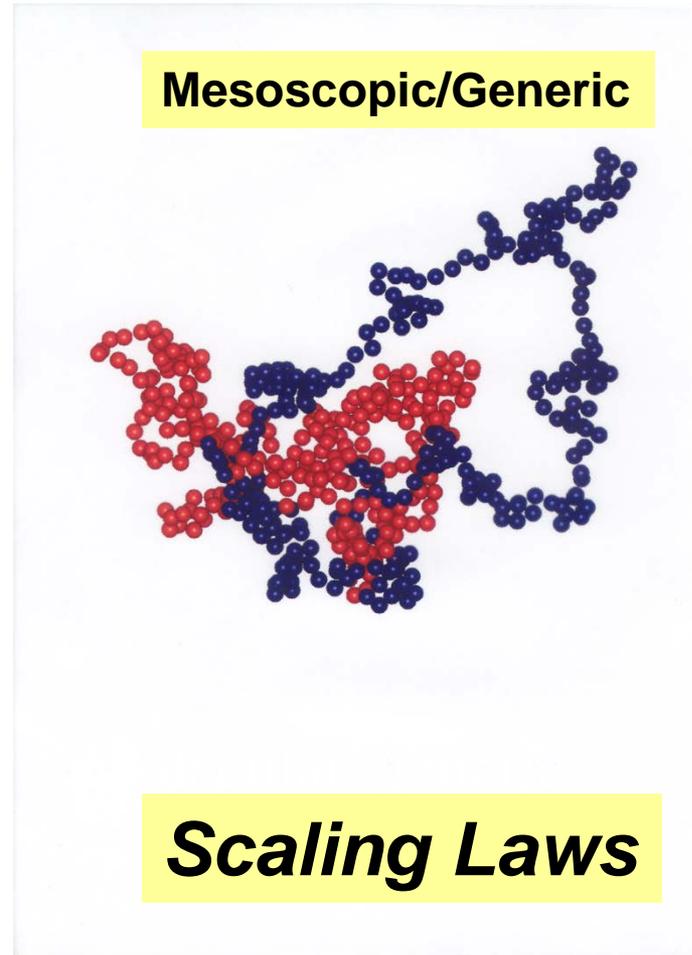


Polymers: Simulation

Structure Property Relations

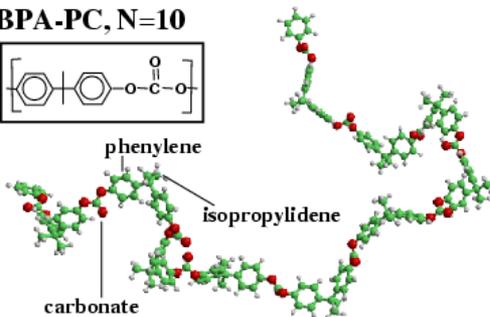
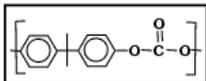


versus



...

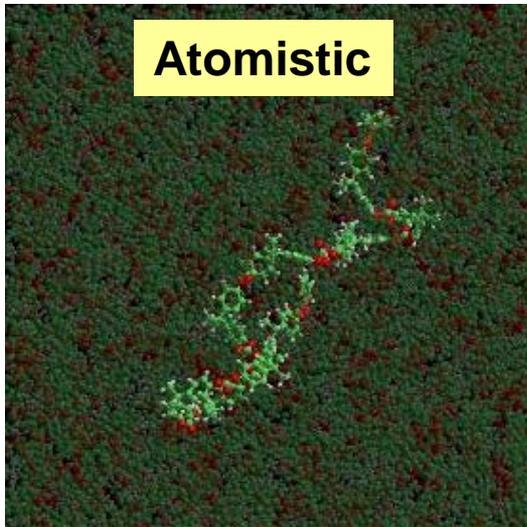
BPA-PC, N=10



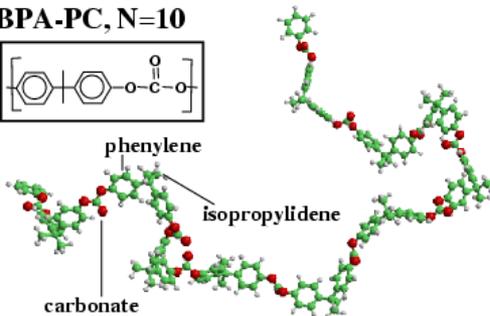
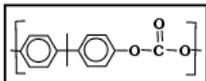


Polymers: Simulation

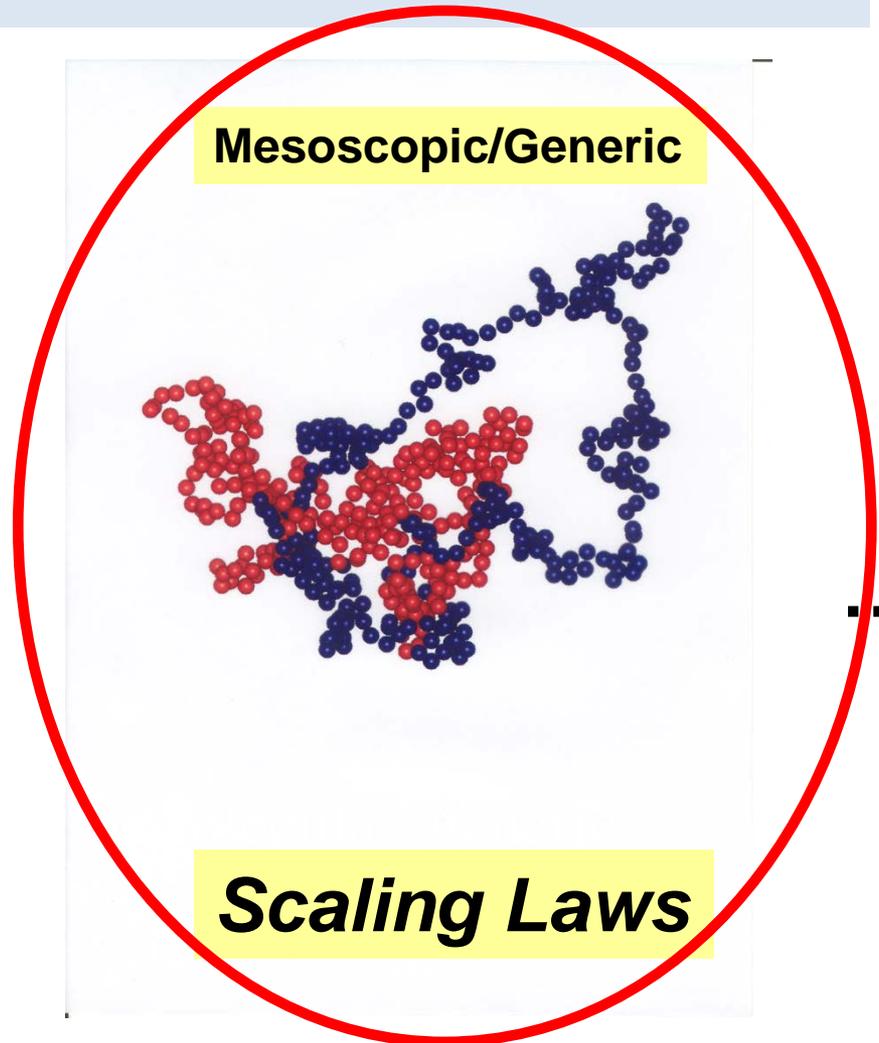
Structure Property Relations



BPA-PC, N=10

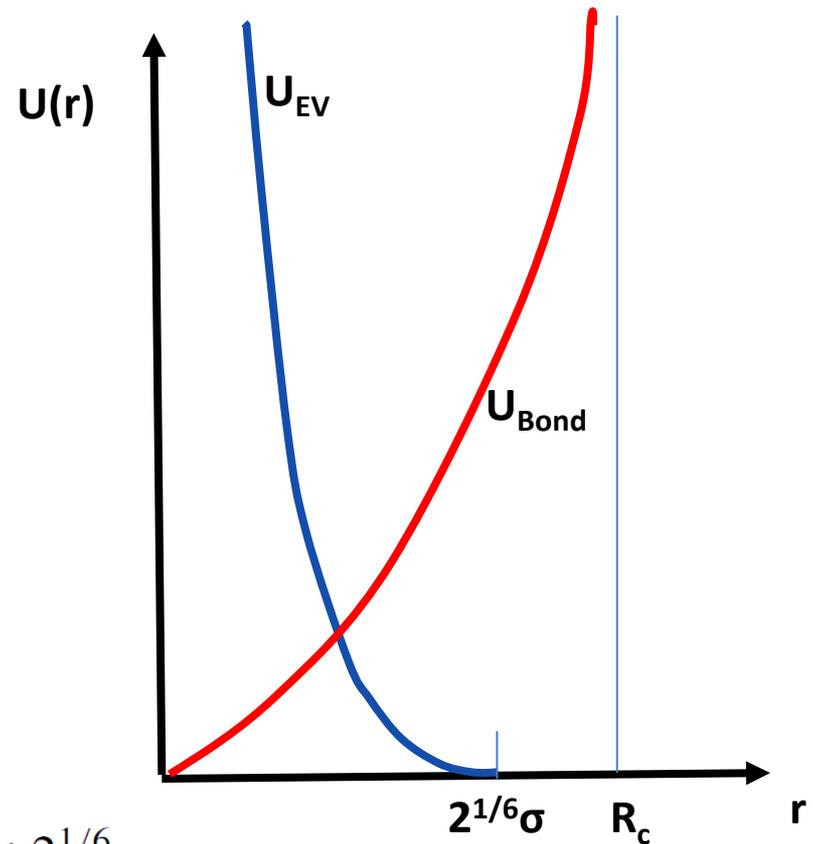
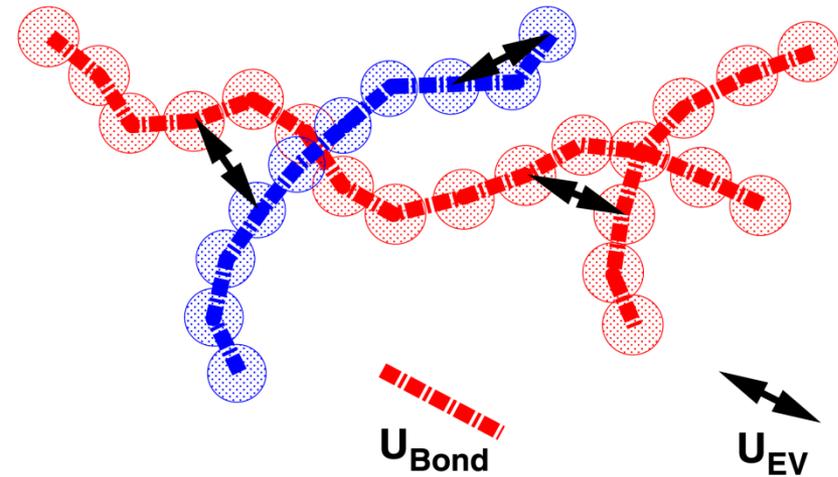


versus



“Simplest Model” Bead-spring model

K.K & G.S. Grest



$$U_{Bond} = -30 k_B T \ln \left\{ 1 - \left(\frac{r}{R_c} \right)^2 \right\}$$

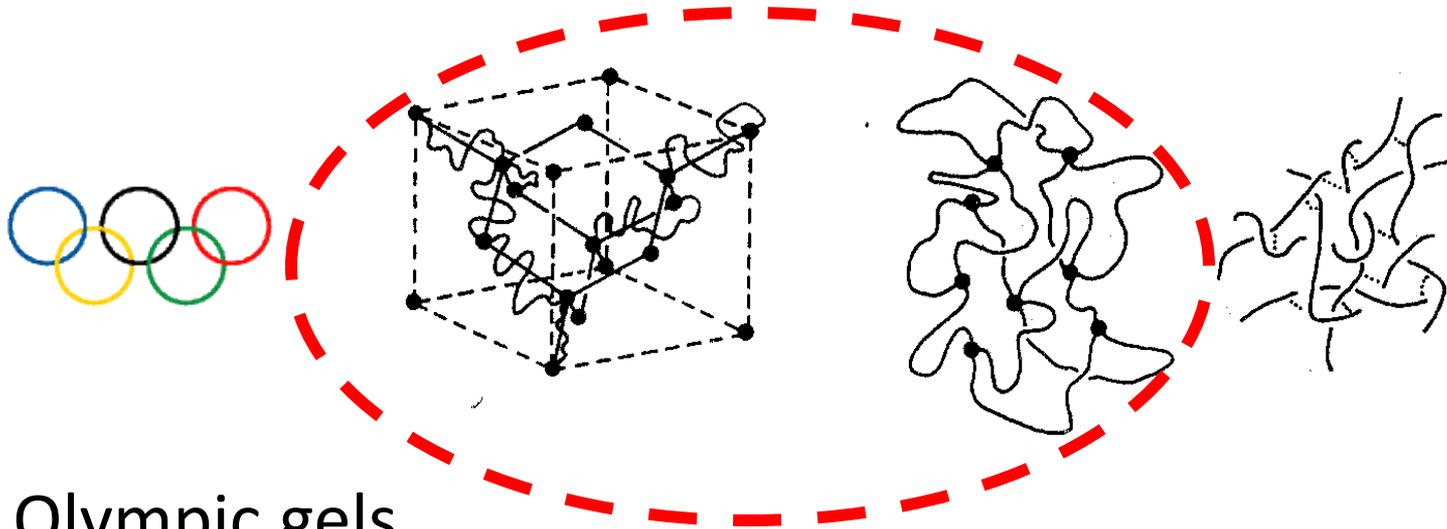
$$U_{pair}(r_{ij}) = \begin{cases} 4\epsilon [(\sigma/r_{ij})^{12} - (\sigma/r_{ij})^6] + \epsilon, & r_{ij} \leq 2^{1/6}\sigma \\ 0, & r_{ij} > 2^{1/6}\sigma \end{cases}$$

- **Topology conservation**

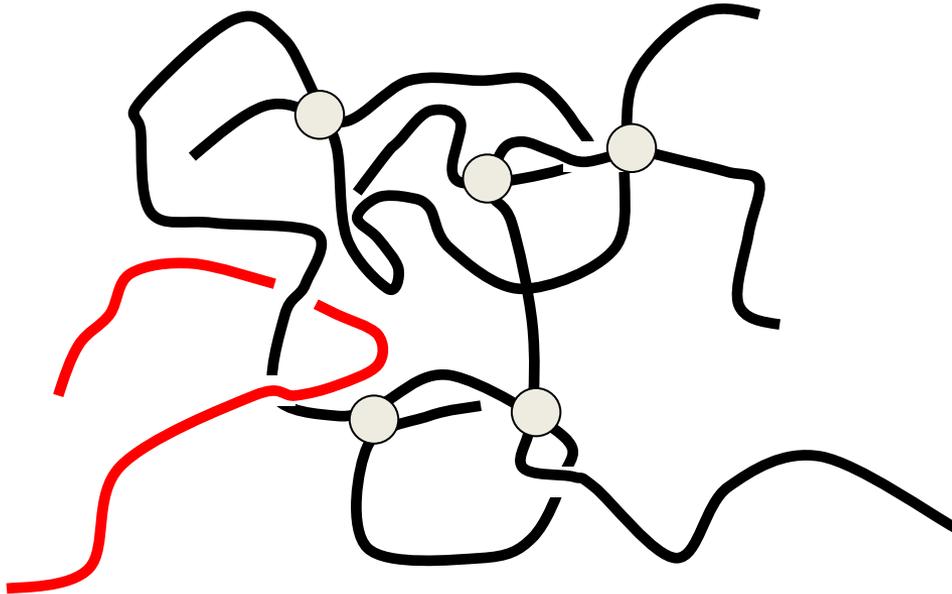
- $U_{Cross} \approx 70k_B T$



Model Networks

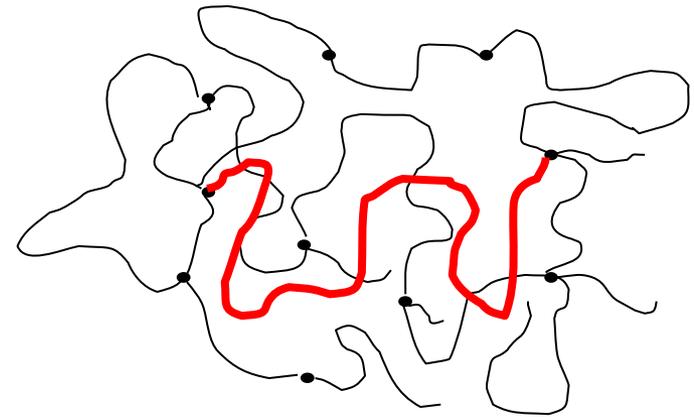
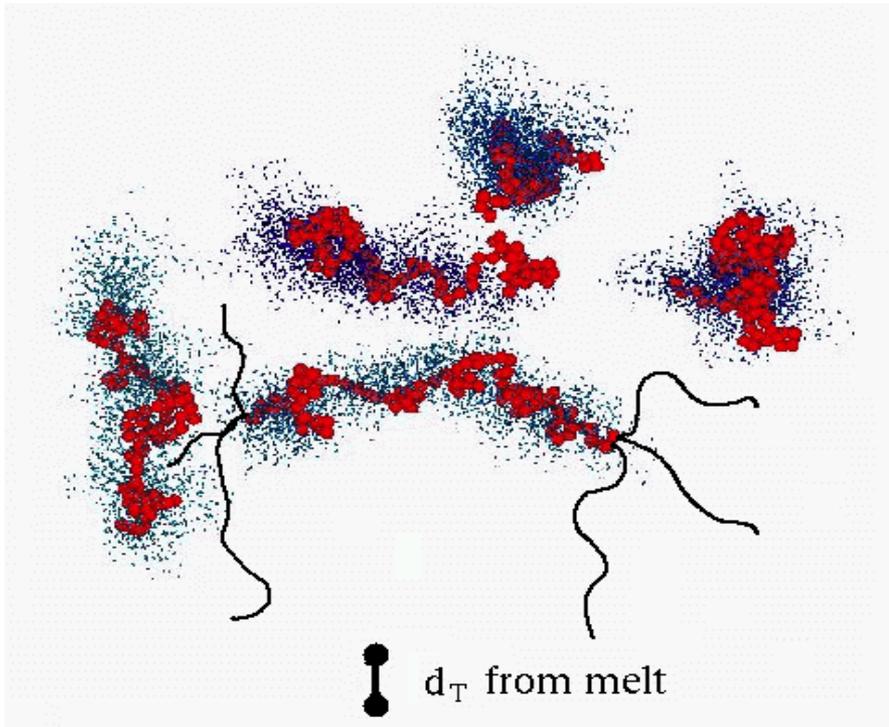


- Olympic gels
- Lattice Networks
- Endlinked melts
- Crosslinked melts



- **Polydispersity**
- **Dangling ends**
- **Self loops**
- **Entanglements**

Confinement of network chains to a “tube”



$$N = 100 \quad N_e \approx 70$$

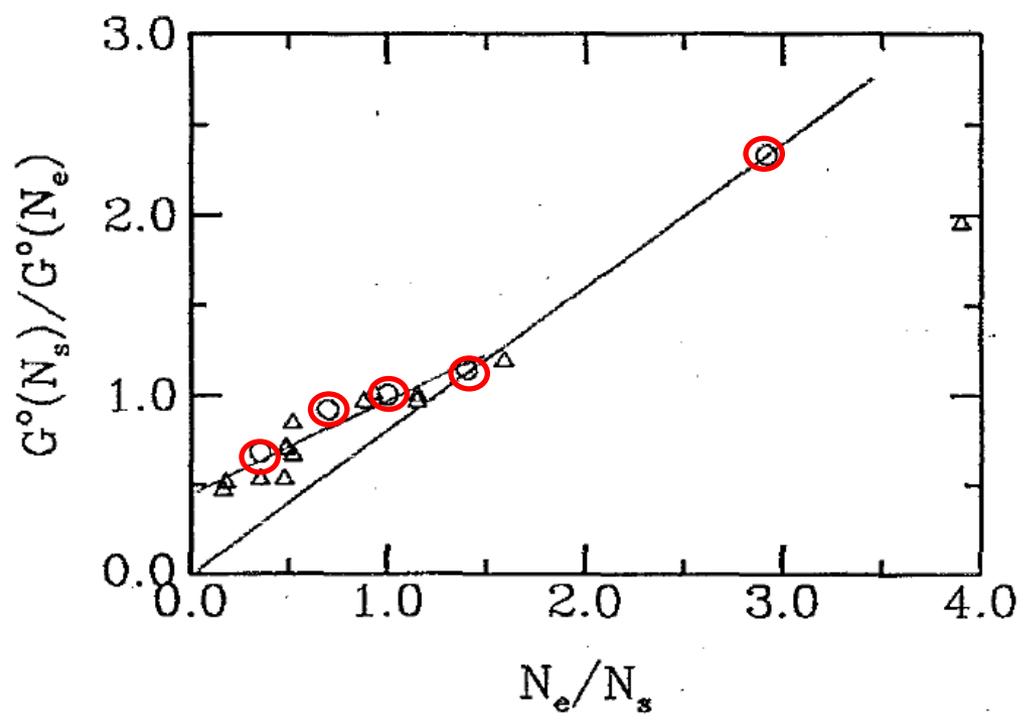
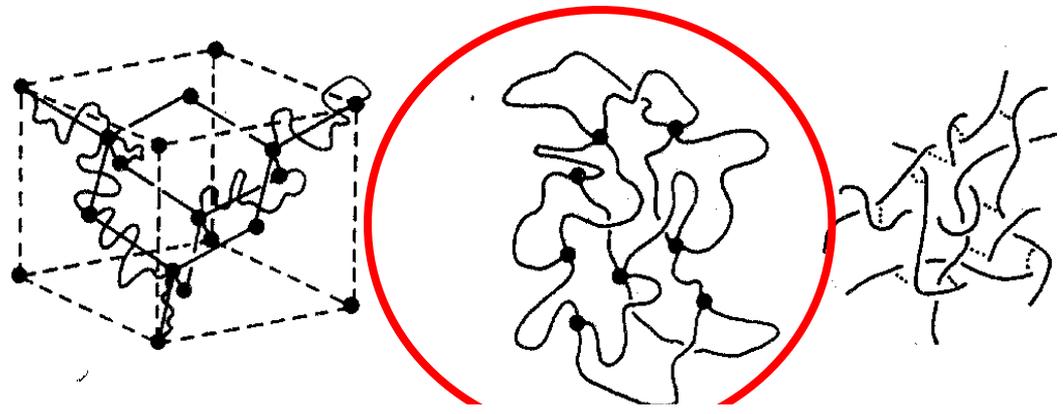
●●●●● initial conformations

■ ■ ■ ■ ■ all subsequent positions



Edgardo R. Duering, KK, Gary S. Grest
JCP, 1994

Elastic Modulus

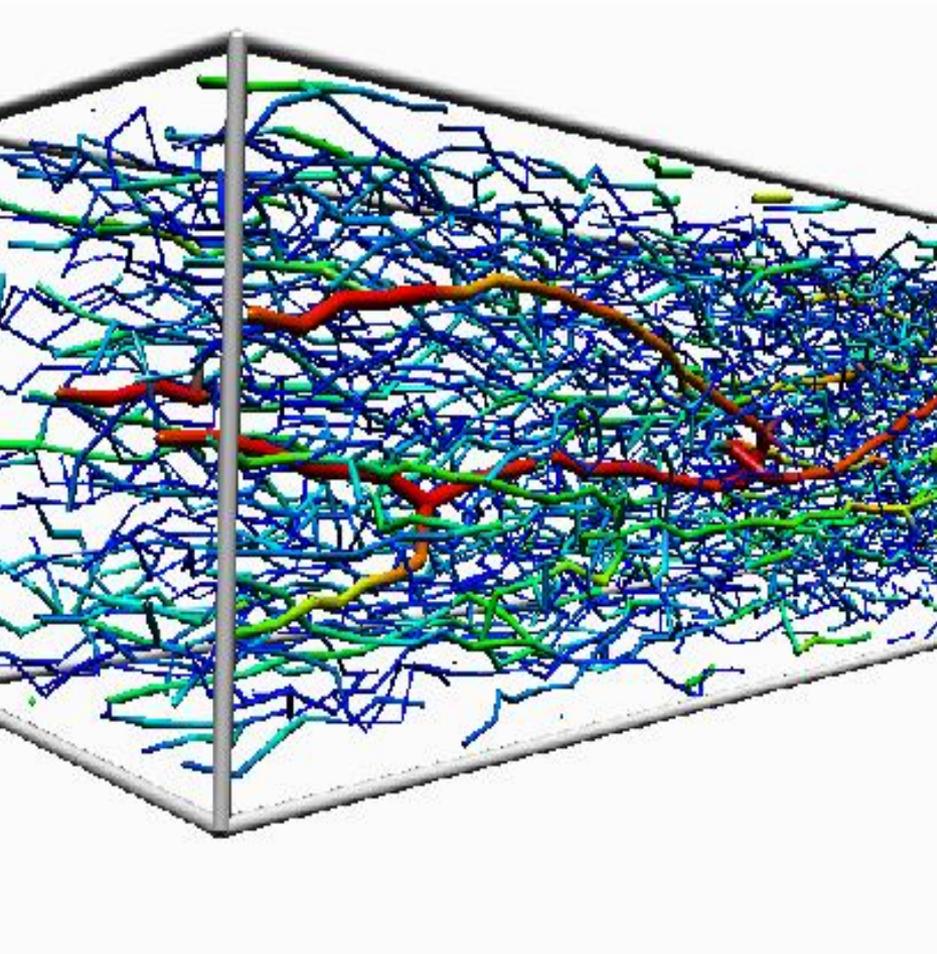
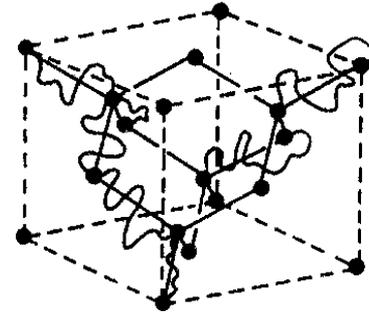


Comparison to PDMS
(Colby & Cohen et al.)





Links & Stress randomly interpenetrating networks

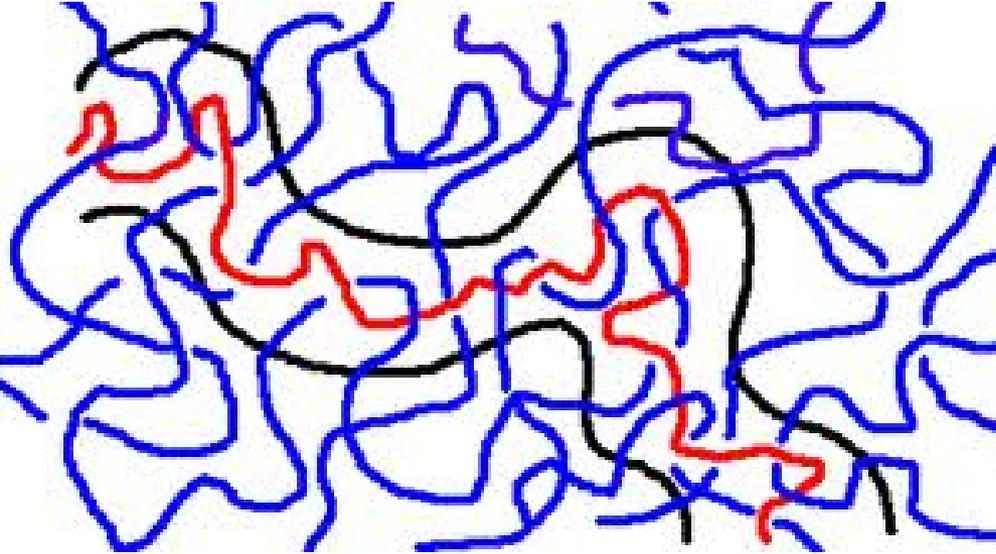


Local stress distribution

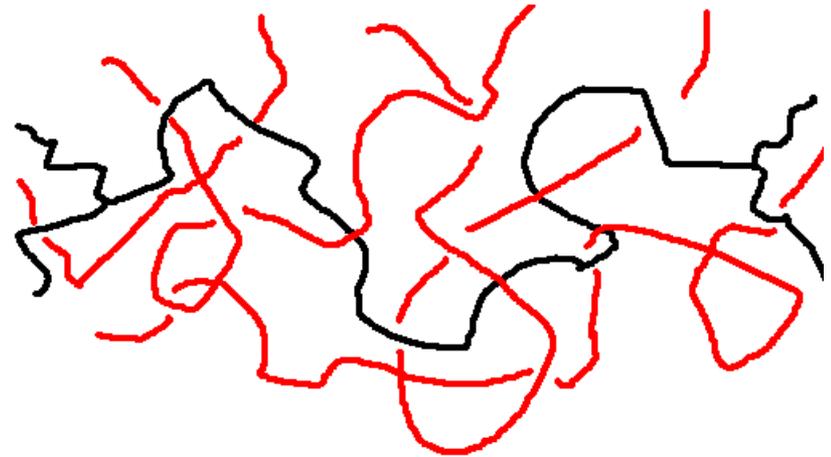
- SMALL STRESS
- LARGE STRESS



Entangled Melts



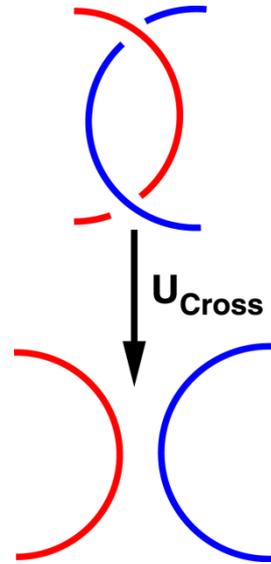
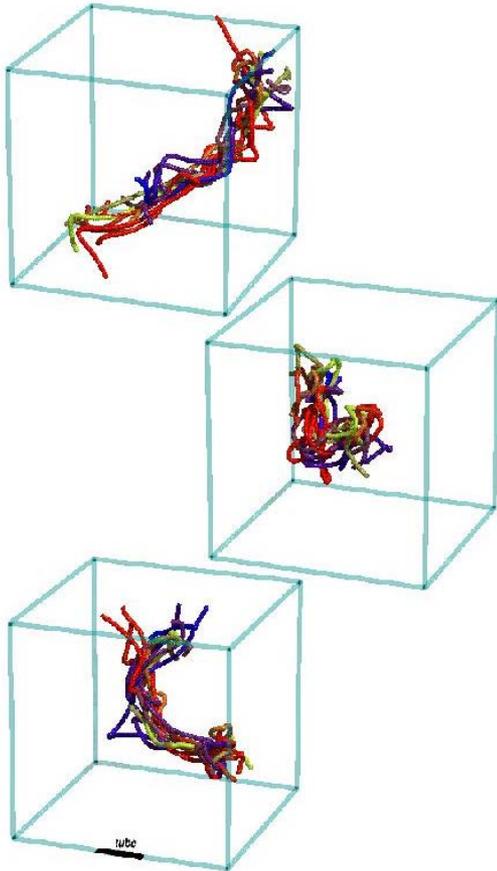
and Solutions



Time Evolution of Melt Chain Conformation (primitive chain in melt)



$$U_{Cross} \approx 70k_B T$$

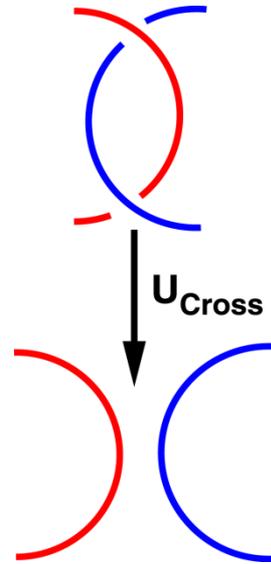
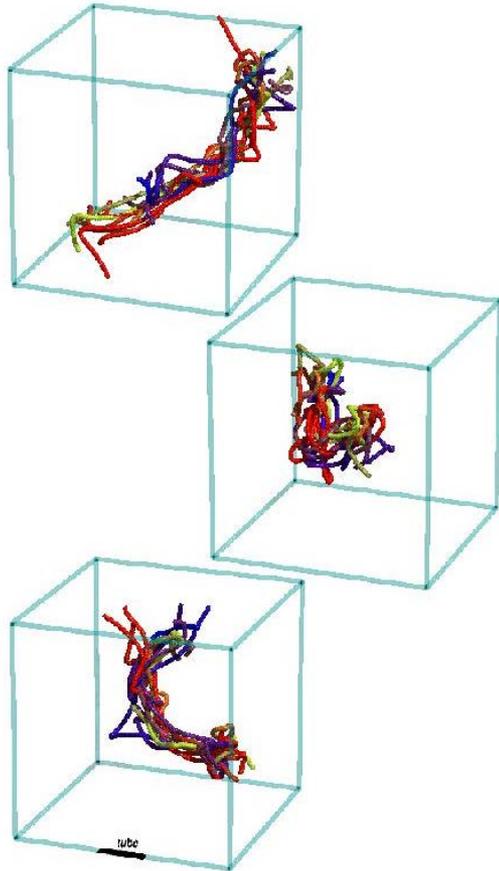


$$t_{max} \approx \tau_{Rouse}$$

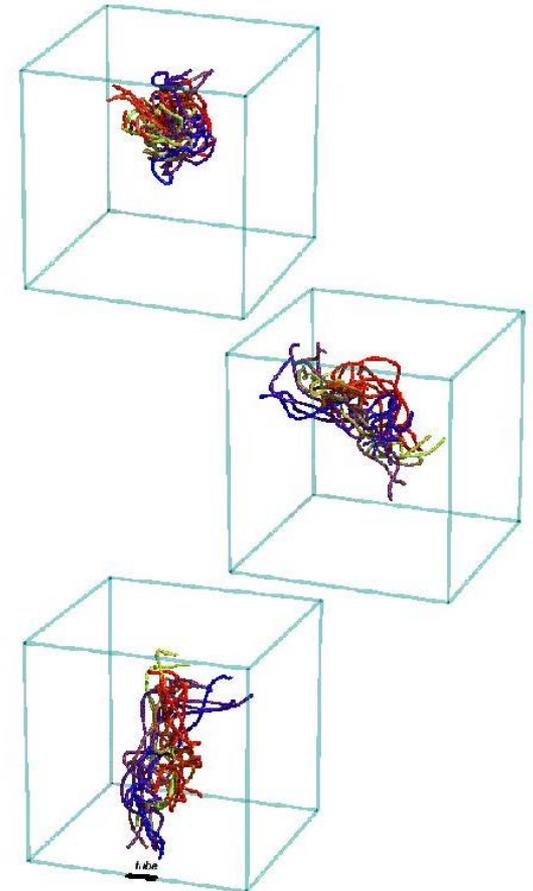
Time Evolution of Melt Chain Conformation (primitive chain in melt)



$$U_{Cross} \approx 70k_B T$$



$$U_{Cross} \approx 4k_B T$$



Crossing barrier CAN
be switched OFF
Without changing $\langle R^2 \rangle$

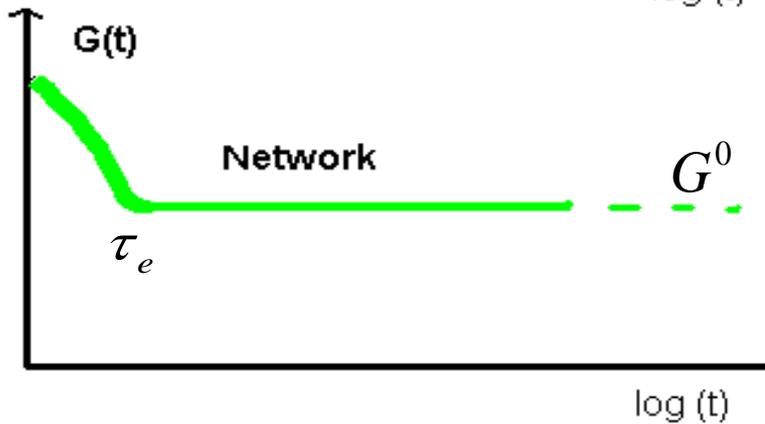
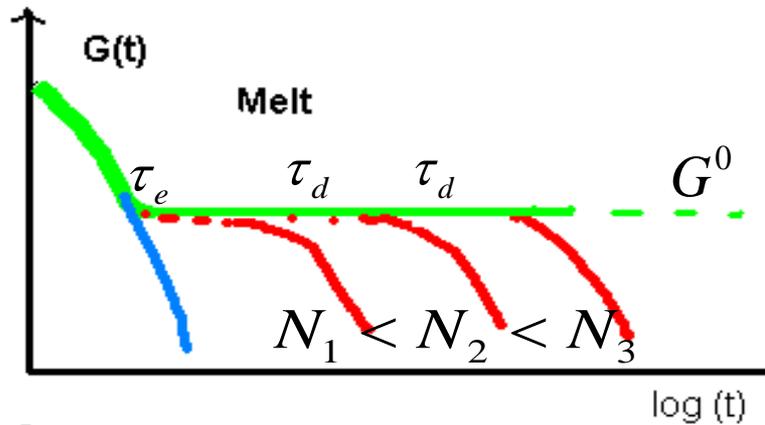
$$t_{\max} \approx \tau_{Rouse}$$



Dynamic Response: Restoring modulus G' after step strain

short chains

long chains



$$\tau_d \sim N^{3\dots}$$

$$G^0 \propto \frac{1}{N_e}$$

$$G^0 \propto \frac{1}{N_s}, \quad N_s \leq N_e$$

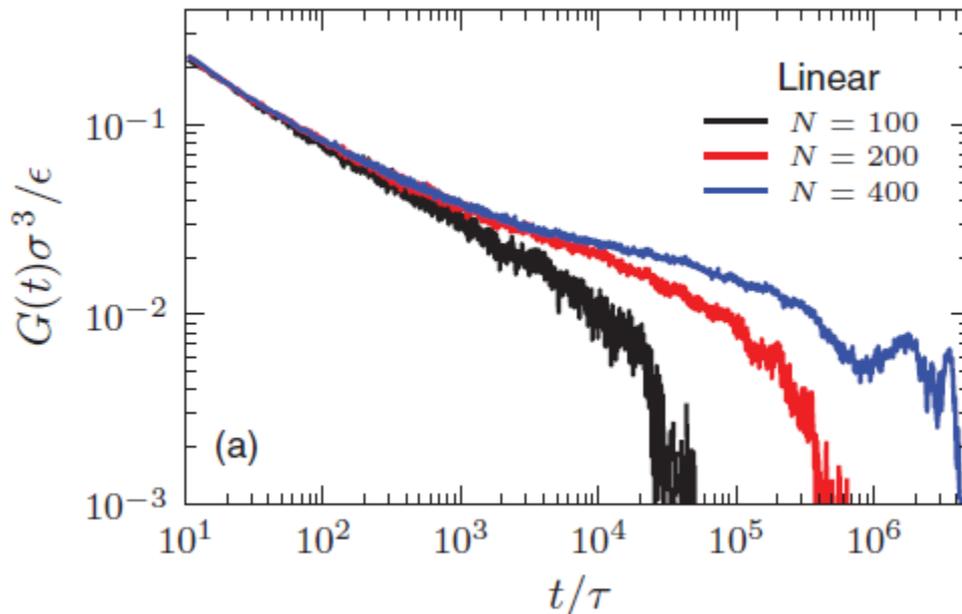
$$\propto \frac{1}{N_e}, \quad N_s \gg N_e$$

N, N_s : chain length
 G^0 independent of N , if $N \gg N_e$



Dynamic Response: Restoring modulus $G'(t)$

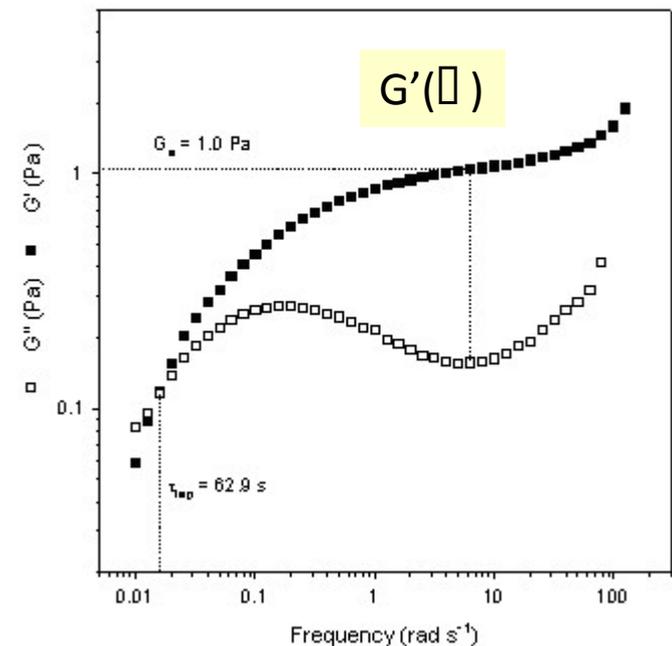
Simulation, $N \lesssim 15N_e$



Halverson et al, 2011

Dense DNA solution

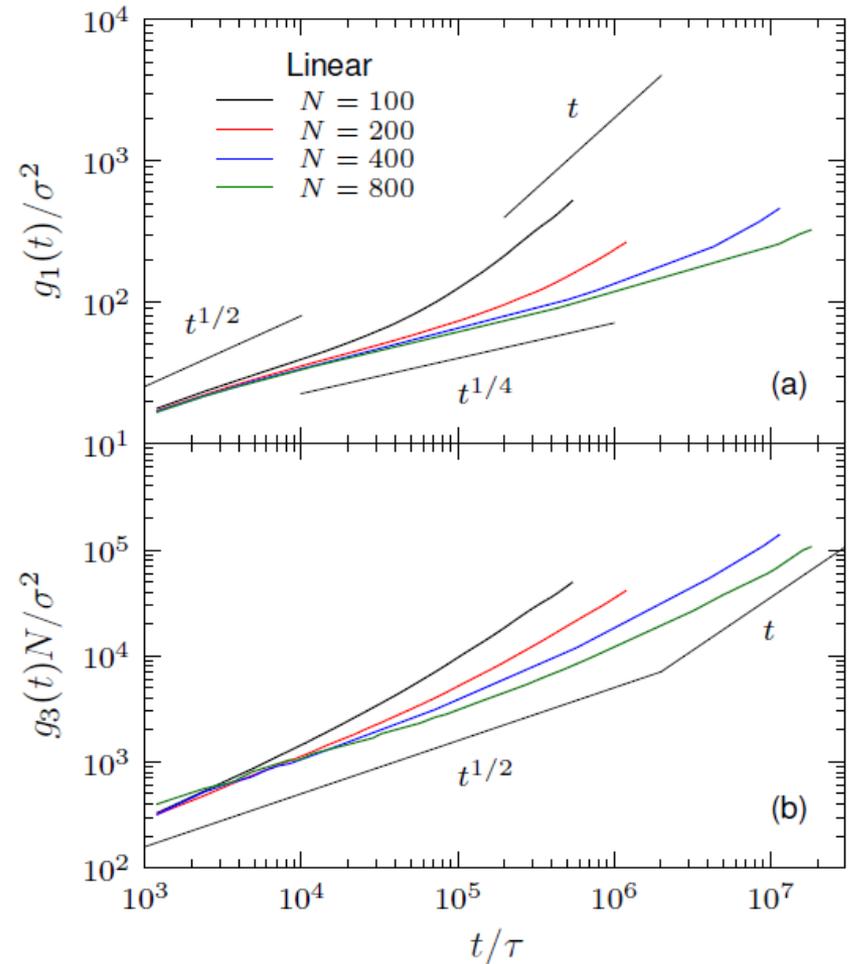
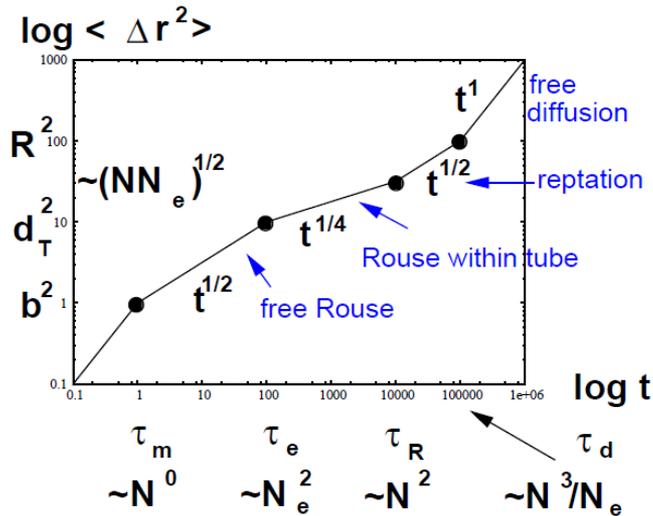
Dynamic Frequency Sweep
1.05 mg/mL DNA Solution





Dynamic Response: mean square displacement

Simulation, $N_e = 28$





How to determine N_e ?

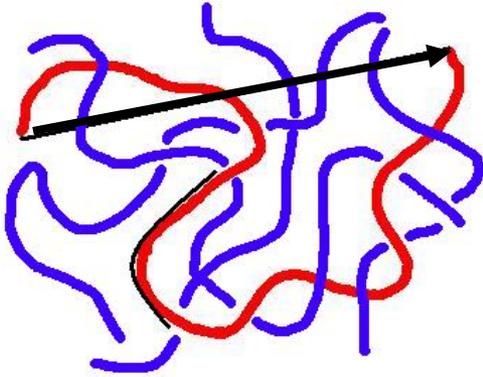
So far: Experiment and simulation measure dynamical quantity and fit N_e .

Original Edward's argument is only based on conformational statistics of the network strands.

Specific Polymers?



Melt: Polymer Conformations, d=3



N: number of bonds per chain

Polymer melt:
Excluded volume screened,
chains are random walks

$$\begin{aligned} \langle R^2(N) \rangle &\propto N \\ P(\vec{r}) &\propto \exp(-3r^2 / 2\langle R^2(N) \rangle) \\ \rho_{self} &\propto \frac{N}{\langle R^2(N) \rangle^{3/2}} \propto N^{-1/2} \end{aligned}$$

Melt density: $\rho_{melt} = \text{const}$

⇒ Volume of each chain is shared by $N^{1/2}$ other chains

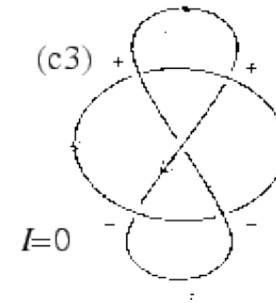
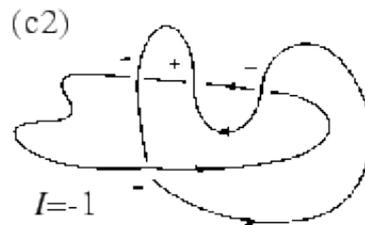
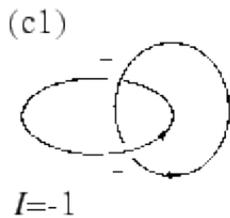
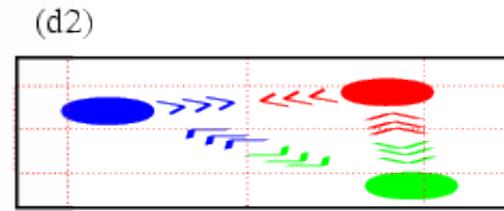
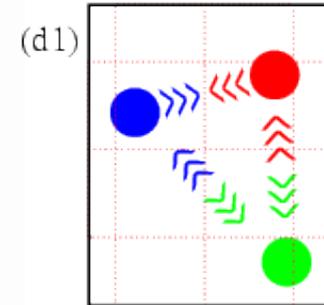
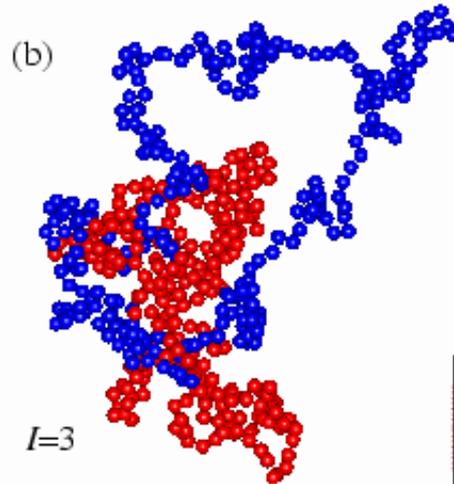
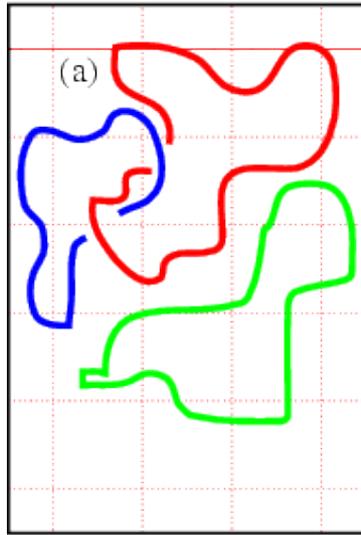
⇒ These $N^{1/2}$ other chains have to impose $O(N)$ constraints to form a tube!

Same holds for networks and gels



Entanglements: Melts

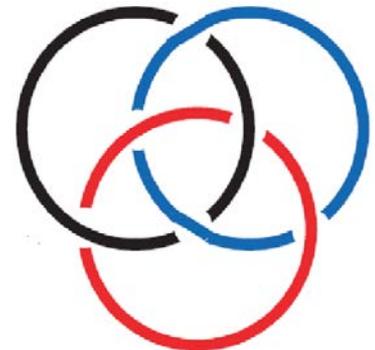
Back to topology arguments



RE&KK (1996)

Entanglements/Primitive Path (Large N)

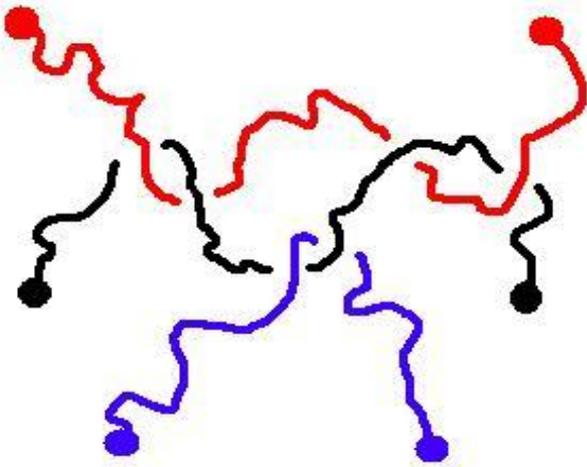
- $O(N)$ constraints per chain needed
⇒ $O(N^{1/2})$ chains sharing same volume!
- Pairwise (Gaussian) linking number in ring melts:
 $O(N^{1/2+1/4})$ NOT SUFFICIENT!
- **Multi chain effects needed, to give sufficient $O(N)$ local constraints!**





Constraints: Multi-Chain Effects

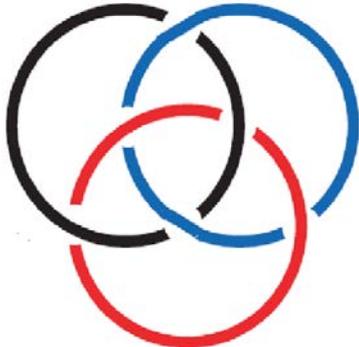
Constraints red-black



2 Constraints (links)



0 Constraints

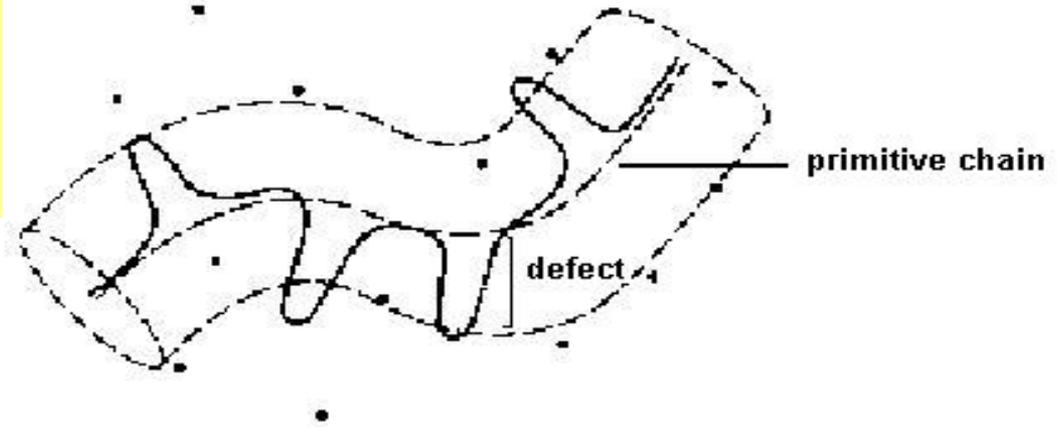


Origin of the Tube: Primitive Path Analysis for equilibrated long chain polymer melts



Need $O(N)$ constraints per chain (distributed along the backbone of the polymer)!

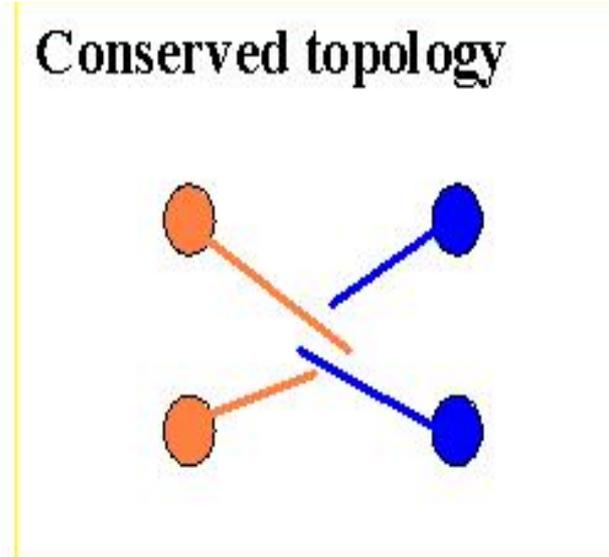
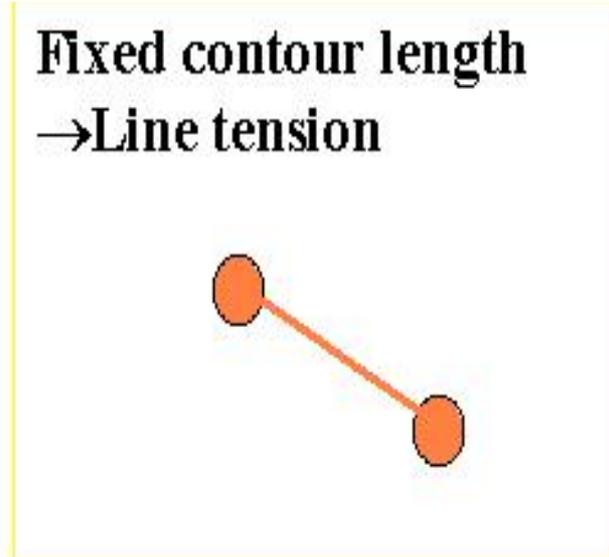
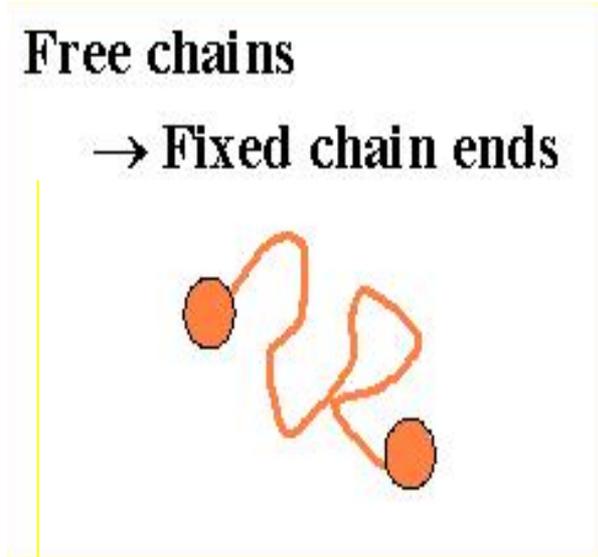
- Length of primitive path
⇒ tube diameter



“The primitive path is the shortest connection between the tube ends, which does not violate any constraint.”

(Edwards '69, Doi & Edwards '78ff)

Primitive Path Analysis for equilibrated long chain polymer melts



Problem: Obstacles are not fixed but other chains
→ primitive path determination needs to be
simultaneous for all chains!

Similar Idea: M. Rubinstein and E. Helfand, JCP 82, 2477 (1985)

Primitive Path Analysis: **Two** chains



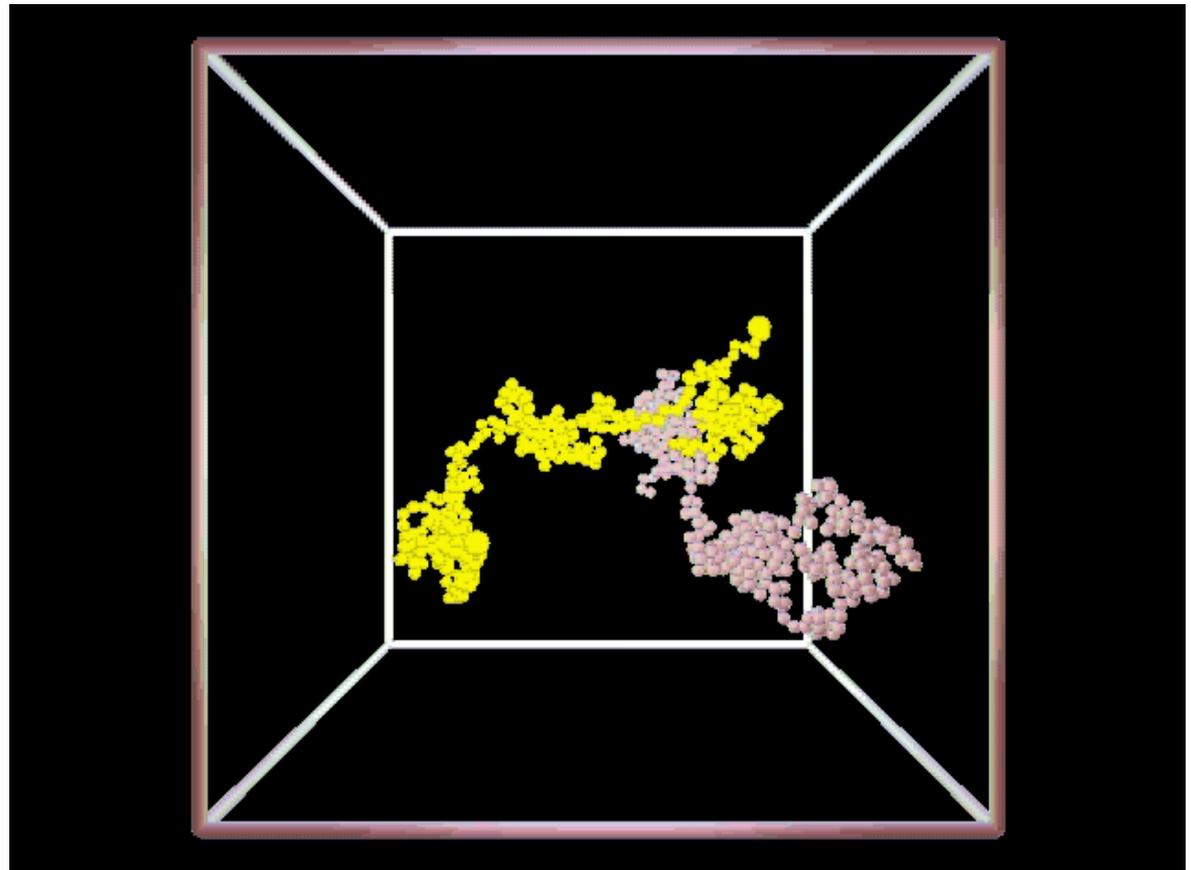
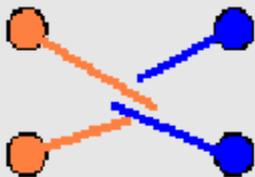
Fixed chain ends



Line tension



Conserved topology



Related work: Milner, Theodorou, Kröger, Larson, Briels,...

Primitive Path Analysis: **All** chains

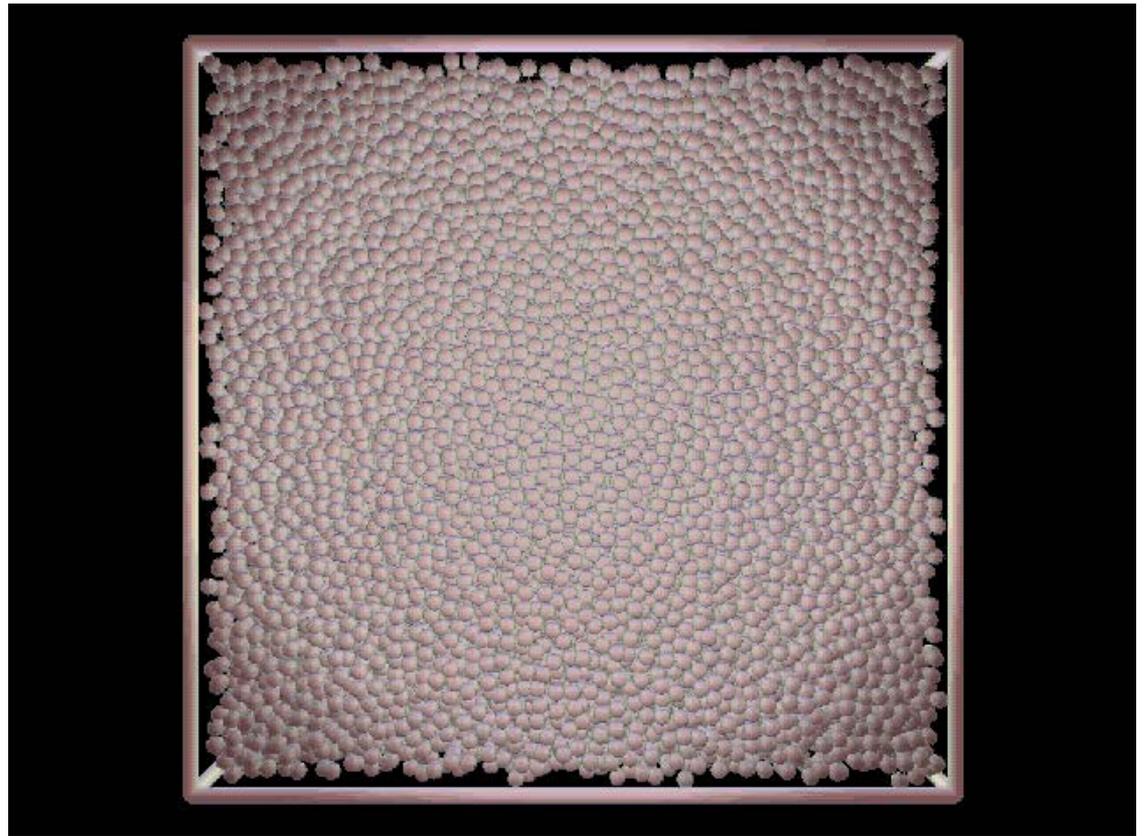
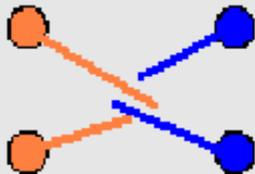
Fixed chain ends



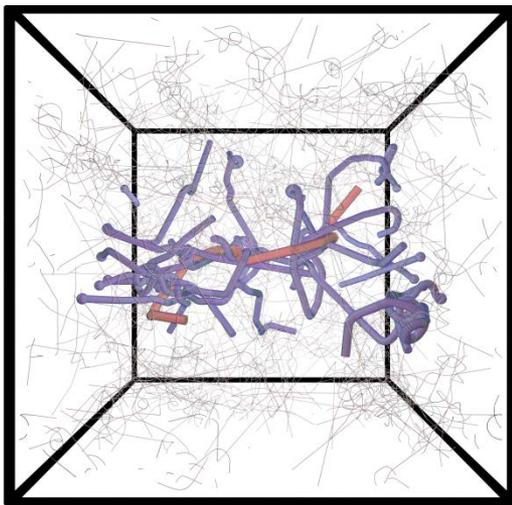
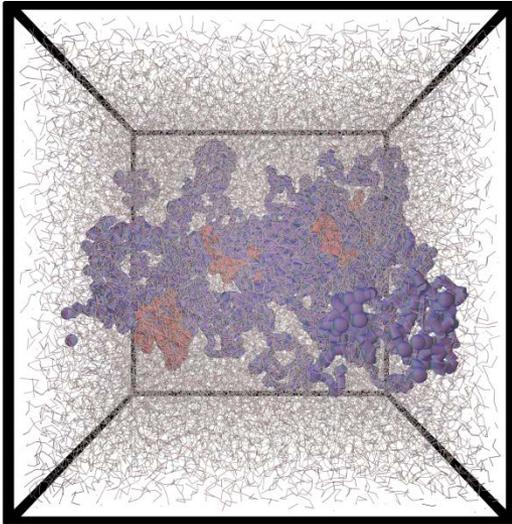
Line tension



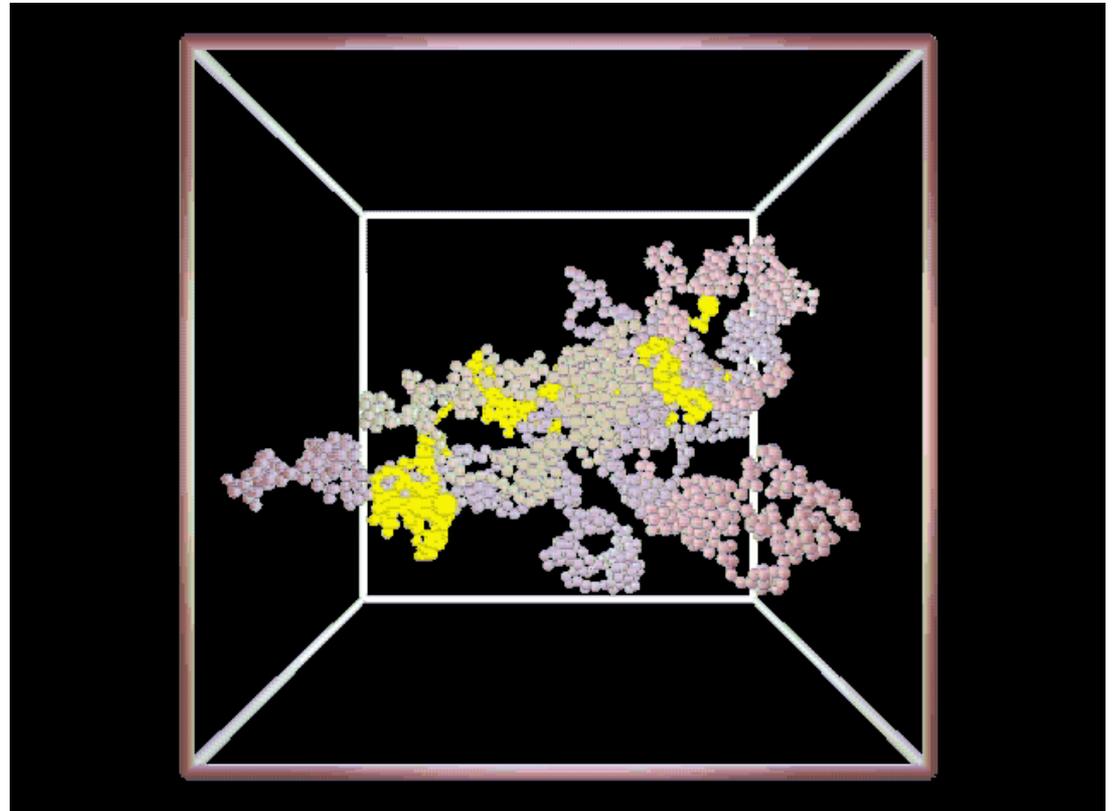
Conserved topology



Entanglements: Primitive Path Analysis



- Evolution of entangled chain cluster (Everaers et al, Science 2004)



Related work: Milner, Theodorou, Kröger, Larson, Briels,...



Characteristic Lengths: l_K and p

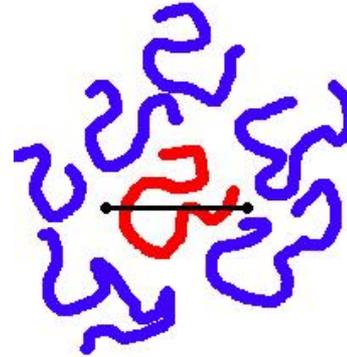
$$p = \frac{N}{\rho} \frac{1}{\langle R^2(N) \rangle}$$



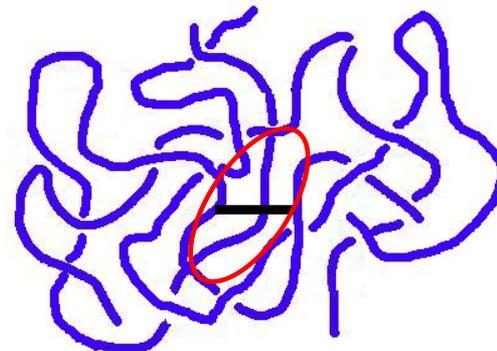
$\rho/N = \text{chain density}$

**Packing
Length p**

Solution at ρ^* : $p \ll R$



Melt: $p \approx \text{strand-strand distance}$



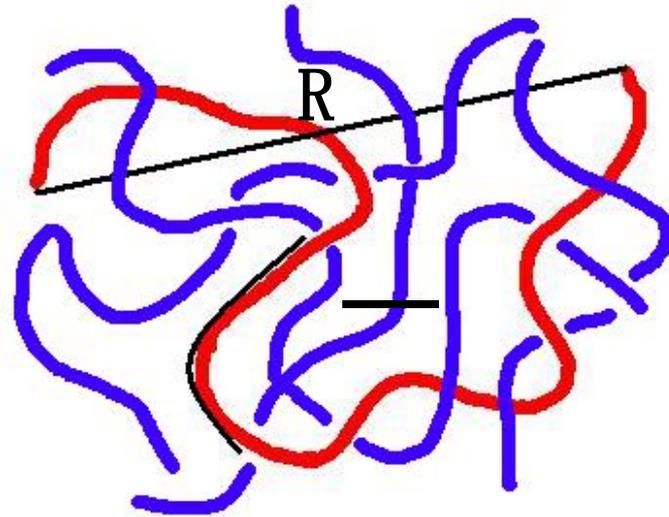
Packing Length ρ and Modulus G

Fetters et al



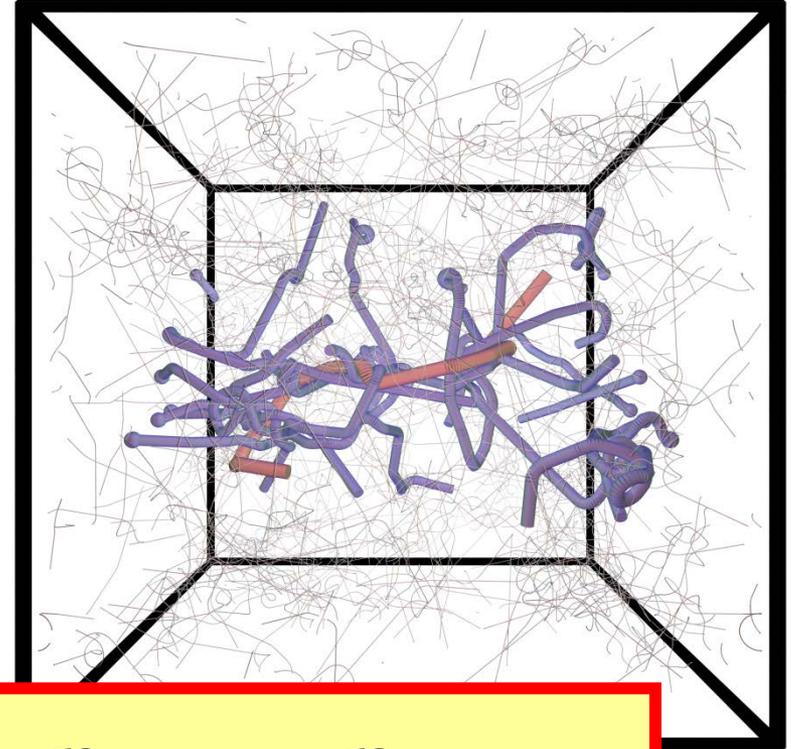
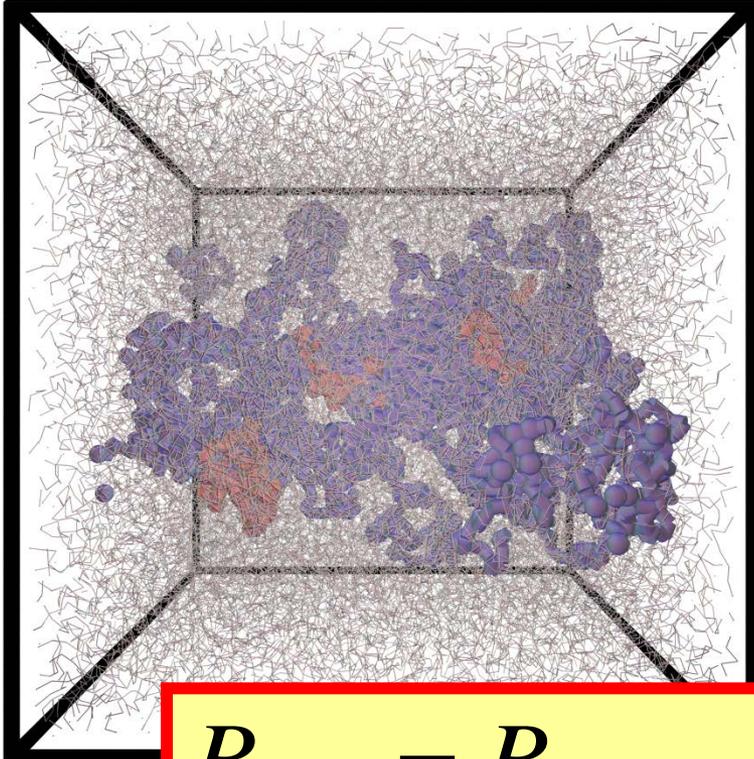
$$G_N \propto \rho^{-3}$$

$$\rho = \frac{N}{\rho} \frac{1}{\langle R^2(N) \rangle}$$



Experimental observation for many polymer melts

Packing length $\rho = \frac{N}{\rho} \frac{1}{\langle R^2(N) \rangle}$



$$R_{melt} = R_{primpath} \Rightarrow \rho_{melt} = \rho_{primpath}$$

$$l_k L = a_{pp} L_{pp}$$

Plateau Modulus:



Modulus (Doi-Edwards):

$$G^0 = 0.8 \rho kT / N_e$$

$$G = \frac{4}{5} \frac{kTN}{pR^2 N_e}$$

$$\dots \boxed{N_e / N = a_{pp} / L_{pp}} \dots$$

$$G = \frac{4}{5} \frac{kT}{p a_{pp}^2} \propto p^{-3}$$

=> Experiments Fetters et al



Plateau Modulus:
rescaled plot – dimensionless units: $\text{\AA} \Leftrightarrow \sigma$

Modulus (Simulation): $G^0 = 0.8 \rho kT / N_e$

Plot of Experimental data by Fetters et al. '99ff:

$$G_N^0 \quad \text{vs} \quad p^{-1}$$

Simulation + Experiment \Rightarrow Length, energy scaling needed:

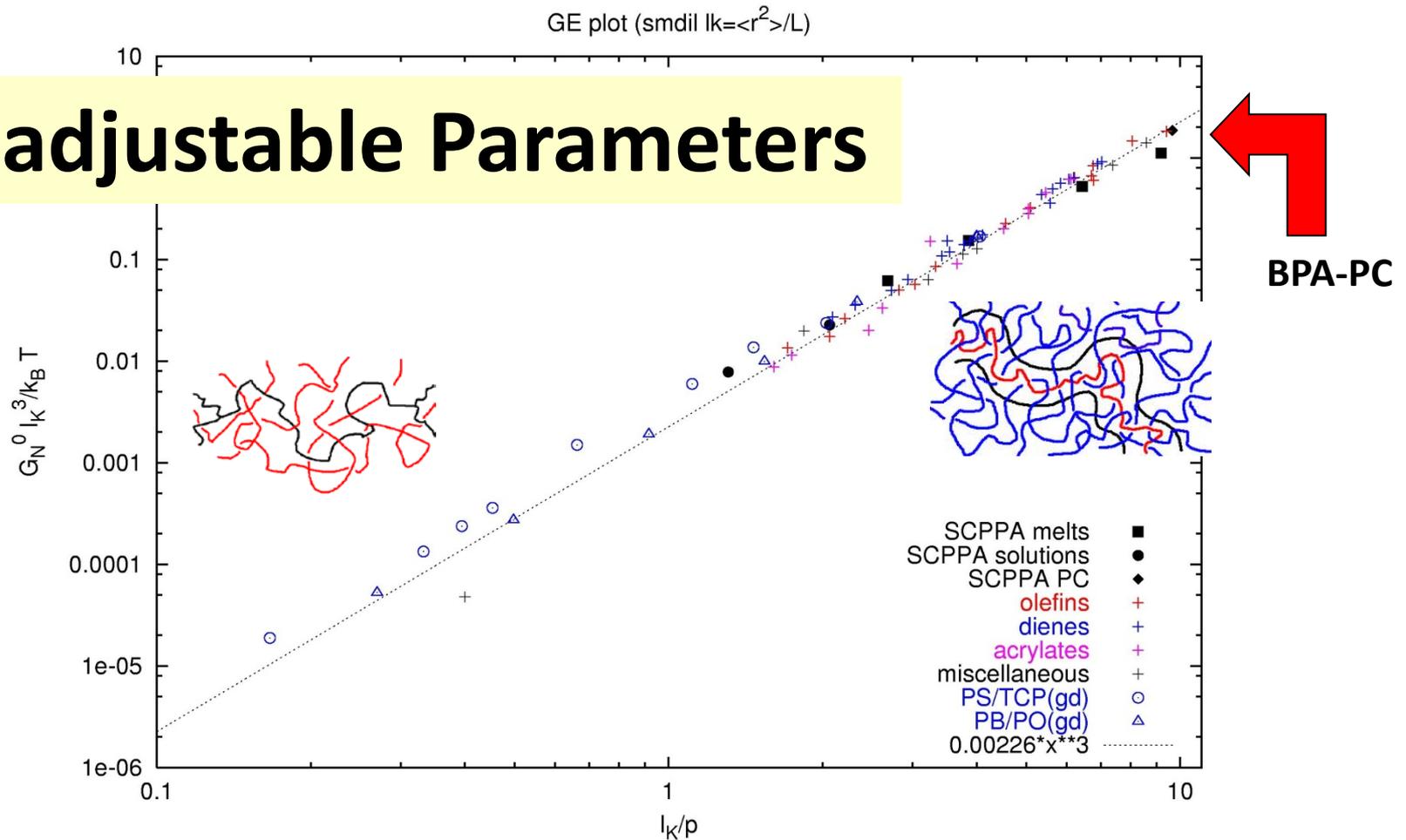
$$G_N^0 l_k^3 / kT \quad \text{vs} \quad l_k / p$$



Plateau Modulus Different Polymers

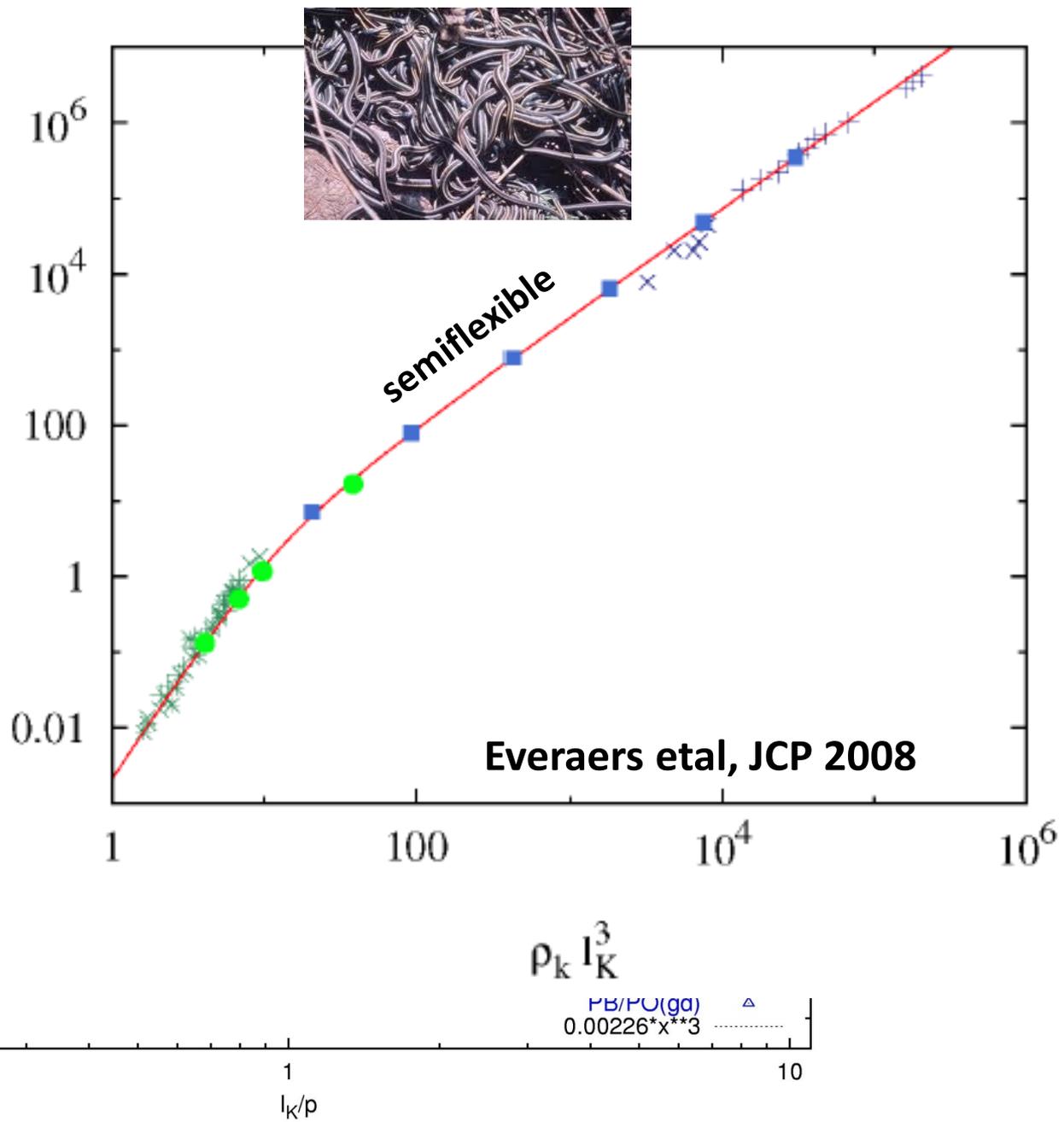
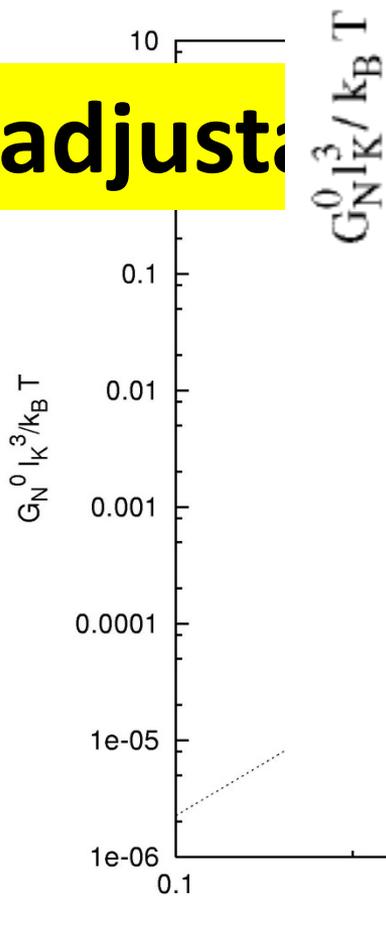
Melts and Solutions: Topological analysis and experiment

NO adjustable Parameters



Pla
Melts and S

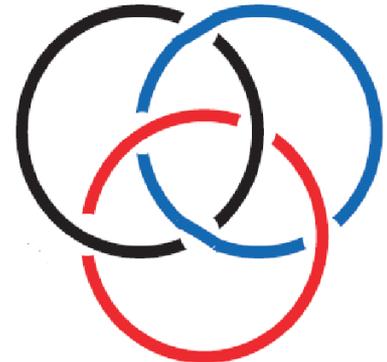
NO adjust



Topological Constraints Matter



- **so far:** topological constraints relevant for dynamics,relaxation
- for melts, solutions no general consequences for conformational statistics (RW character of chains)
- **Now:** constraints on conformations: gels, rings

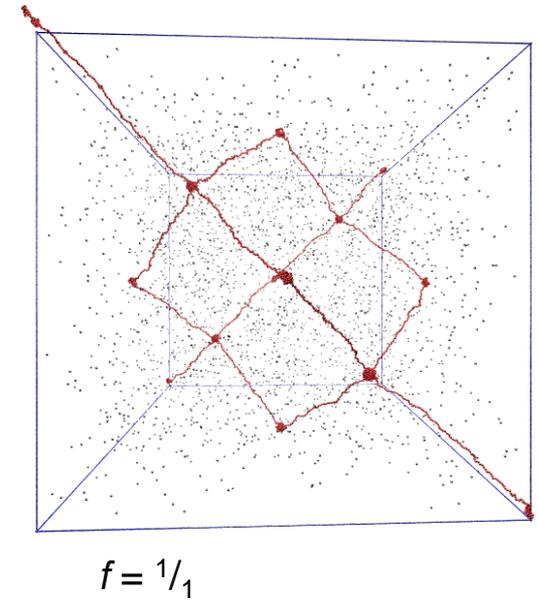
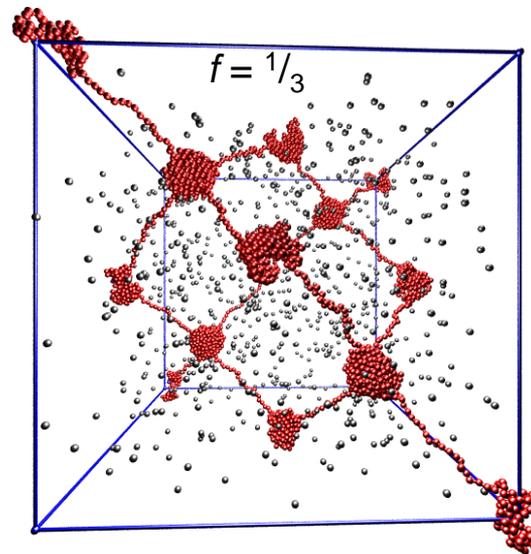
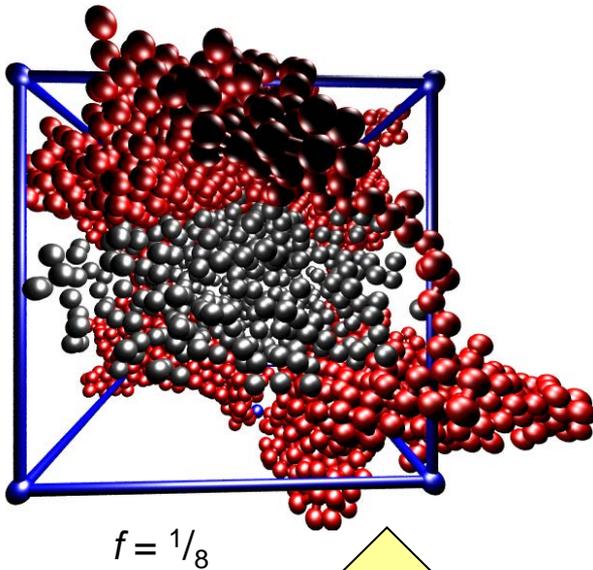


Shrinking Gels: Poor Solvent Hydrogels

high salt
Low charge fraction

...
...

low salt
high charge fraction

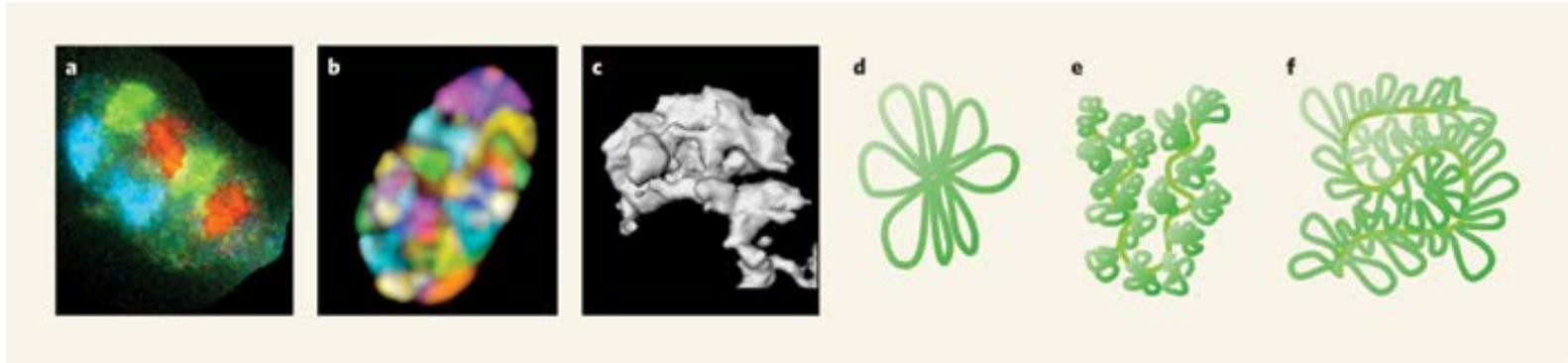


Segregated strands?

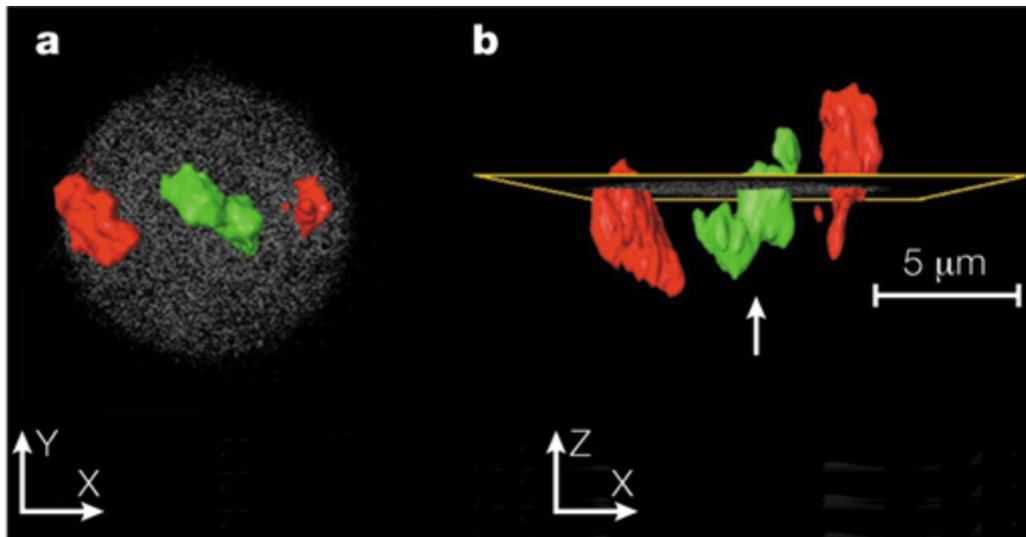
- osmotic pressure of “counterions” enforces swelling (\leftrightarrow good solvent picture)
- nodes act as condensation nuclei
- “pearls” form due to balance with surface tension (\leftrightarrow Rayleigh instability)

Chromosome Territories

K.J. Meaburn, T. Mistelli, Nature 445, 379-381(2007)



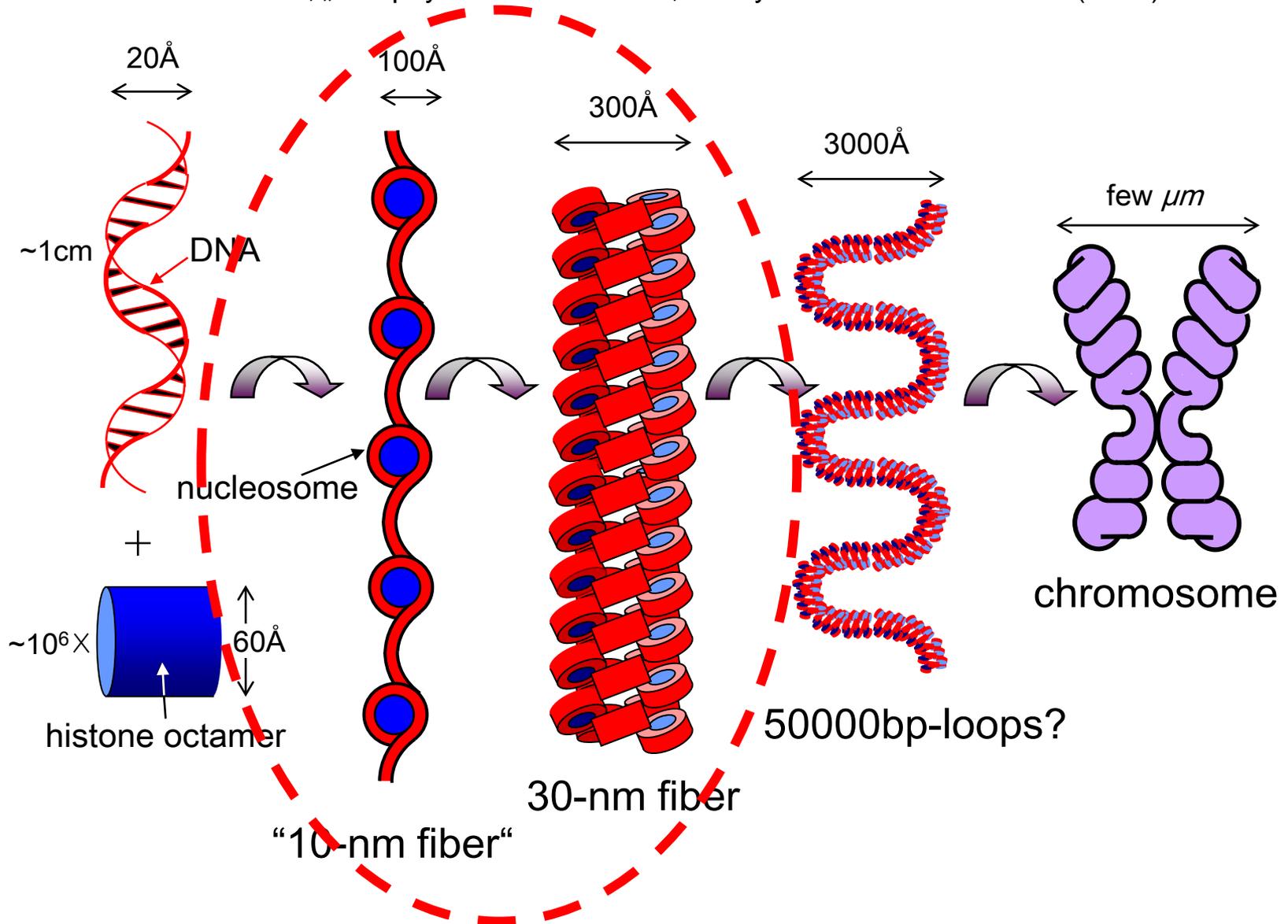
T. Cremer, C. Cremer, Nature Reviews Genetics vol. 2, pp. 292-301 (2001)





structure of chromatin

from review H. Schiessel, „The physics of chromatin“, J. Phys.: Condens. Matter 15 (2003) R699

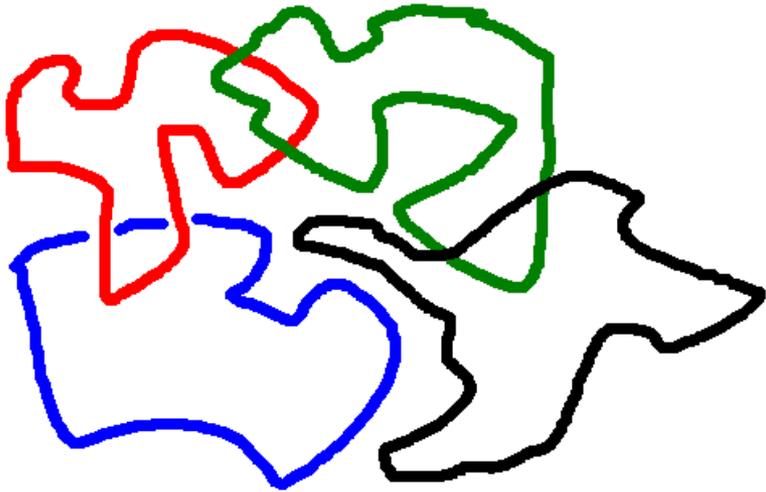


Questions



- Why do chains in a collapsed model (polyelectrolyte) gel not interpenetrate?
- Human DNA about 2m contour length, few cm of 10nm fiber in a nucleus of a few μm , how can they pack?
- Why seem DNA (10nm fibers) not to entangle, but segregate?
- **Look for a most simple model system**

Melt of non concatenated ring polymers



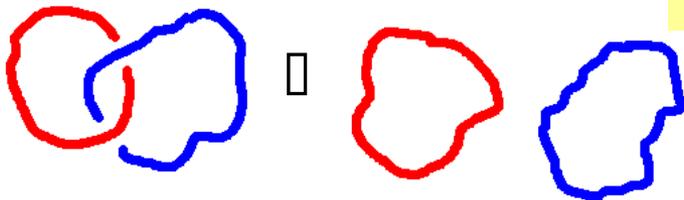
Conformation of the rings?

Random walk: $R^2(N) \propto N^{2\nu}$, $\nu=1/2$

\Rightarrow Rings interpenetrate each other

Collapsed ring: $R^2(N) \propto N^{2\nu}$, $\nu=1/3$

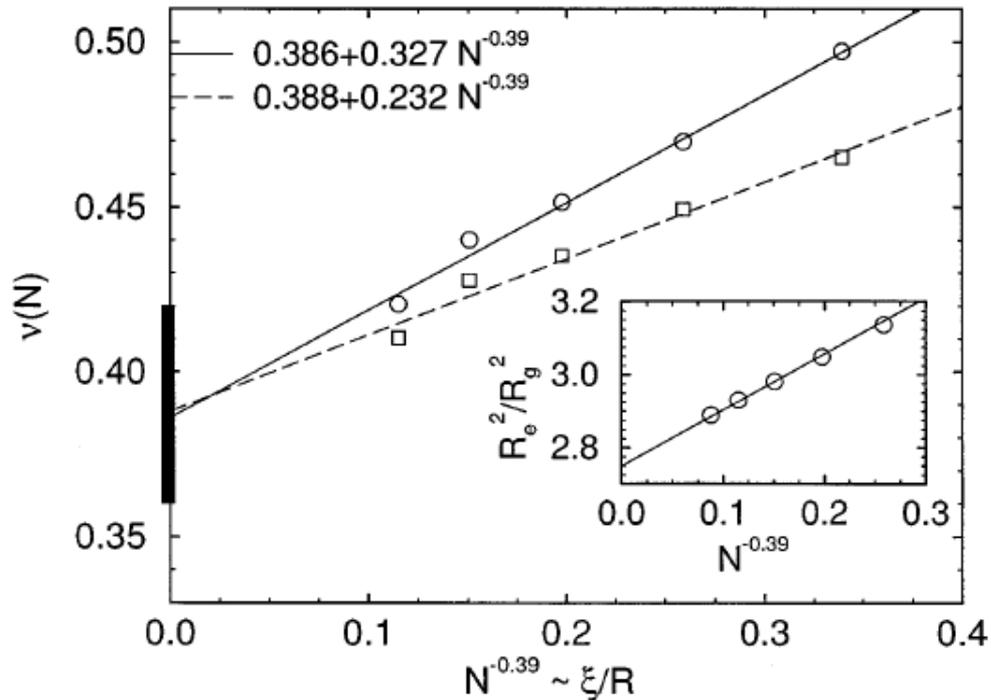
\Rightarrow Rings segregate,
finite number of neighbors





Topological effects in ring polymers: A computer simulation study

M. Müller,* J. P. Wittmer,† and M. E. Cates



**Contributions to Flory free energy
(Deutsch & Cates)**

'osmotic pressure' of other chains:

$$F \propto kTR^d/N$$

Surface term of EV interaction:

$$F \propto kTN/R^2$$

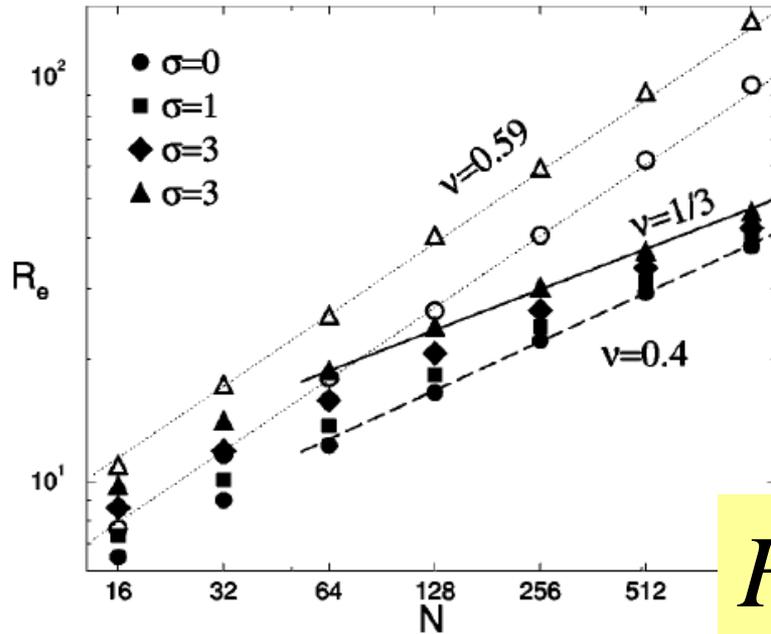
$$\nu = 2/(d+2)$$

$$\nu = 2/5$$



Topological effects in ring polymers. II. Influence of persistence length

M. Müller,^{1,2} J. P. Wittmer,^{2,3,*} and M. E. Cates²



Speculation of higher order contributions

$$F \propto kT (R^d / N)^\alpha + kTN / R^2$$

\Rightarrow

$$\nu = (\alpha + 1) / (3\alpha + 2)$$

$$\alpha = 1 \rightarrow \nu = 0.4$$

$$\alpha = \infty \rightarrow \nu = 1/3$$

(preliminary) Physical picture

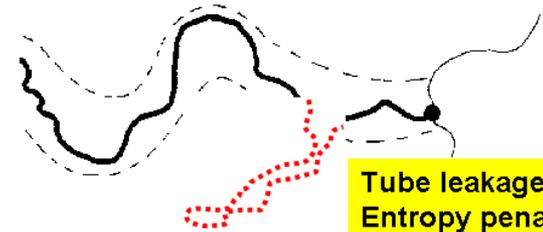
Grosberg, Vettorel, KK

- If rings collapse, (naively) expect core plus corona



$$d_T \ll N_e$$

(first simple picture)



Tube leakage:
Entropy penalty for
defects too large

$$N \leq N_e$$

$$N_e \leq N \leq O(\approx 10N_e)$$

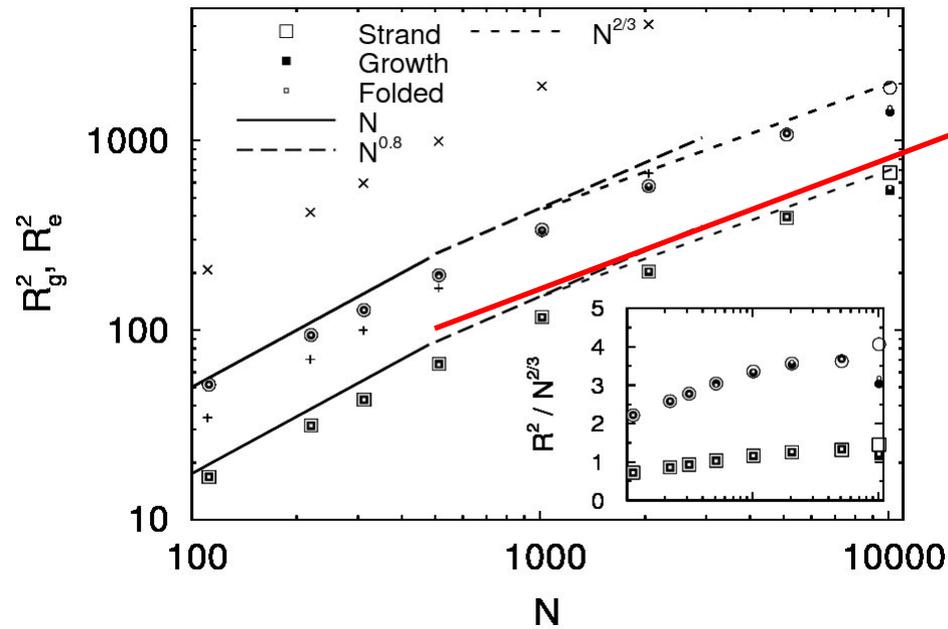
$$N \geq O(\approx 10N_e)$$

No effect, random walk

Begin “shrinking”,

Segregation regime

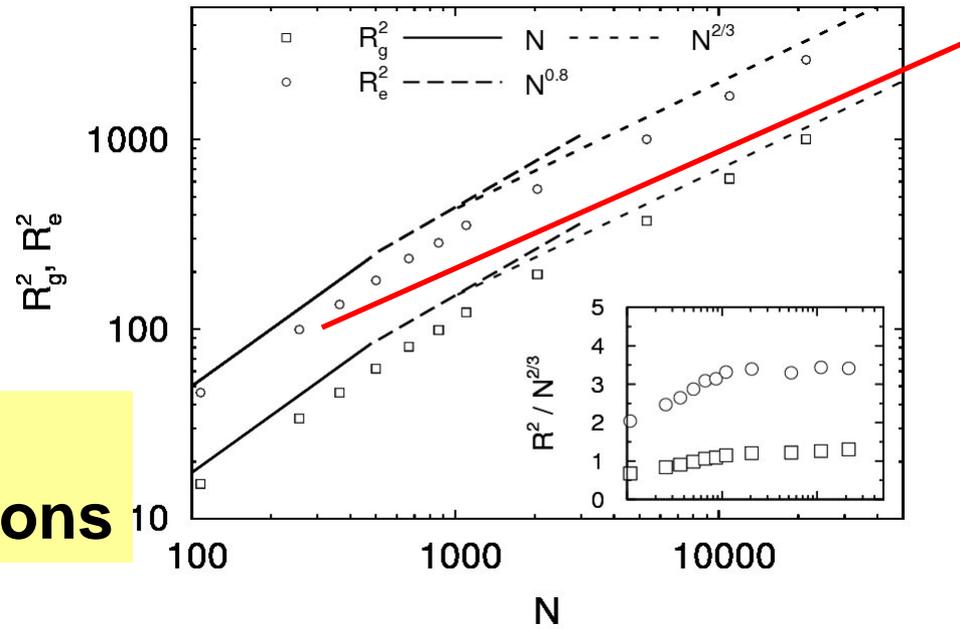
Ring extensions: MC and MD



← = melt of many rings

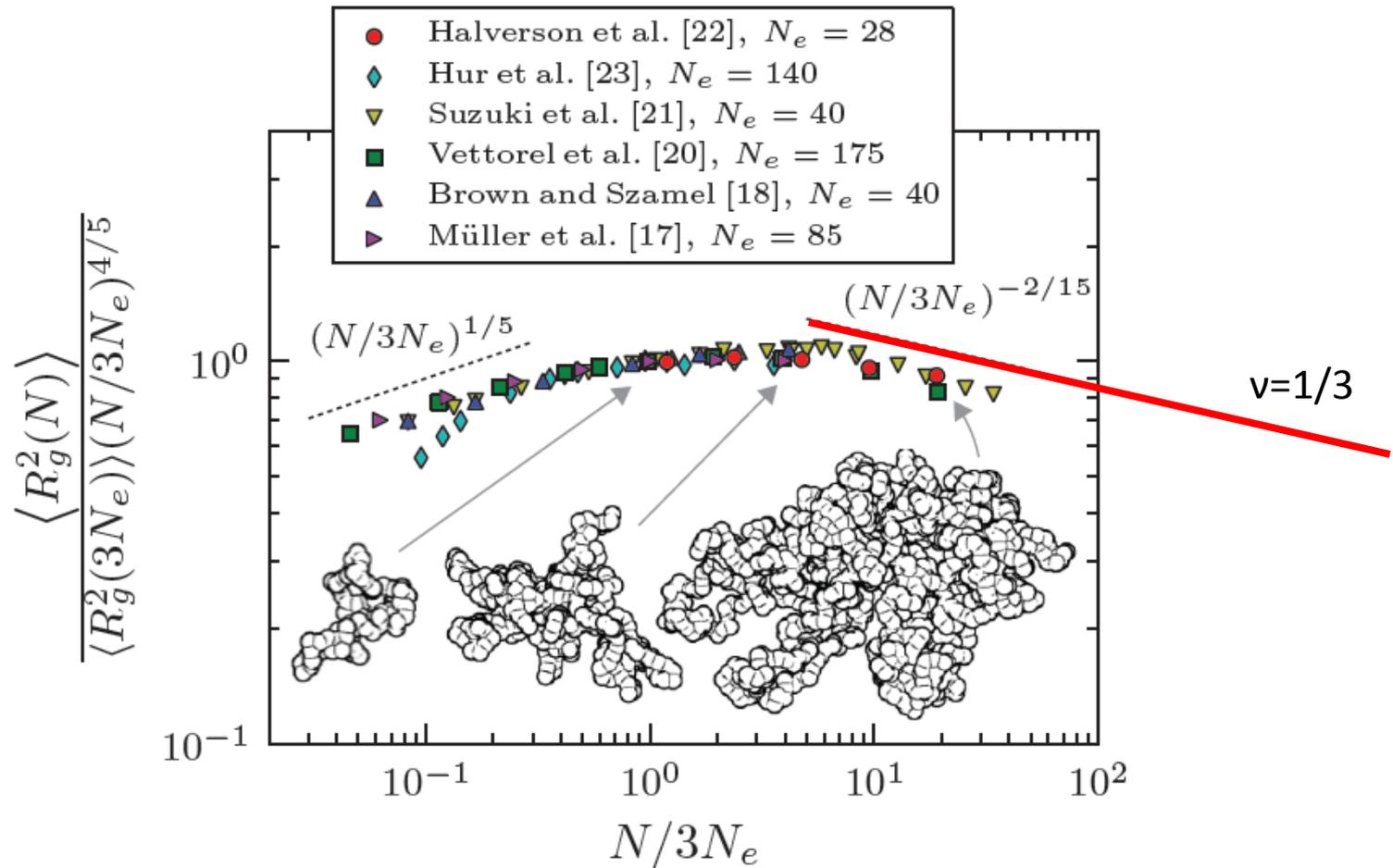
$\square = 1/3$

melt of one ring, seeing itself via periodic boundary conditions



Ring extensions: Collection of data

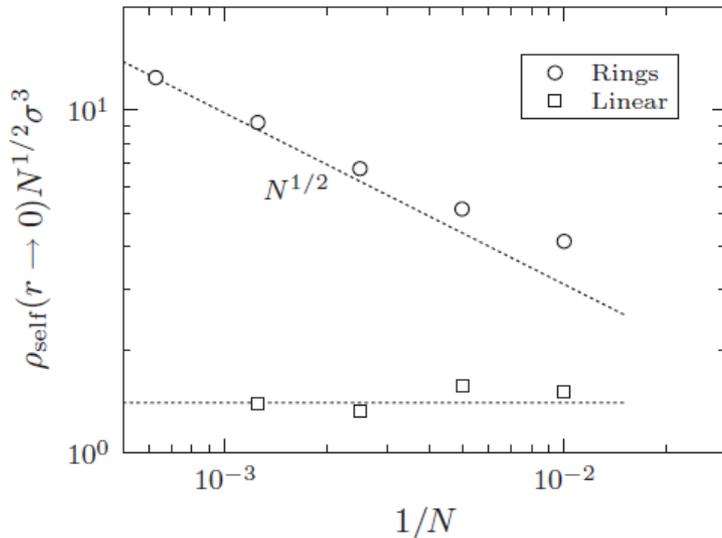
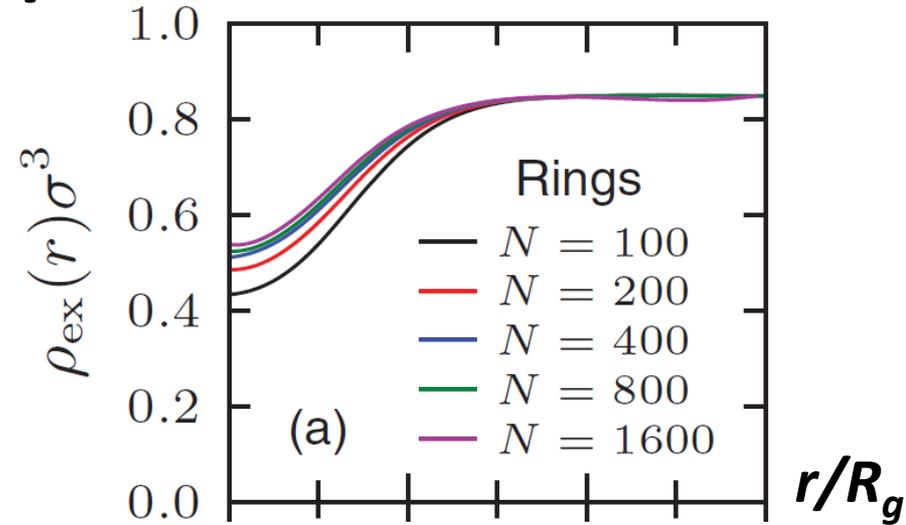
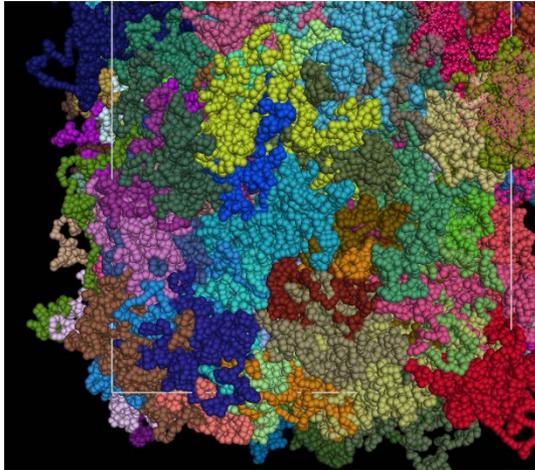
$$R_g^2(N)/R_g^2(3N_e)(N/N_e)^{4/5} \text{ vs } N/N_e$$



Melt of Ring Polymers: Core Density



$$\rho_{\text{self}}(r)$$



Self density

$$\rho_{\text{intra}}(N \rightarrow \infty) \approx 1/3 \rho_{\text{melt}} (!)$$

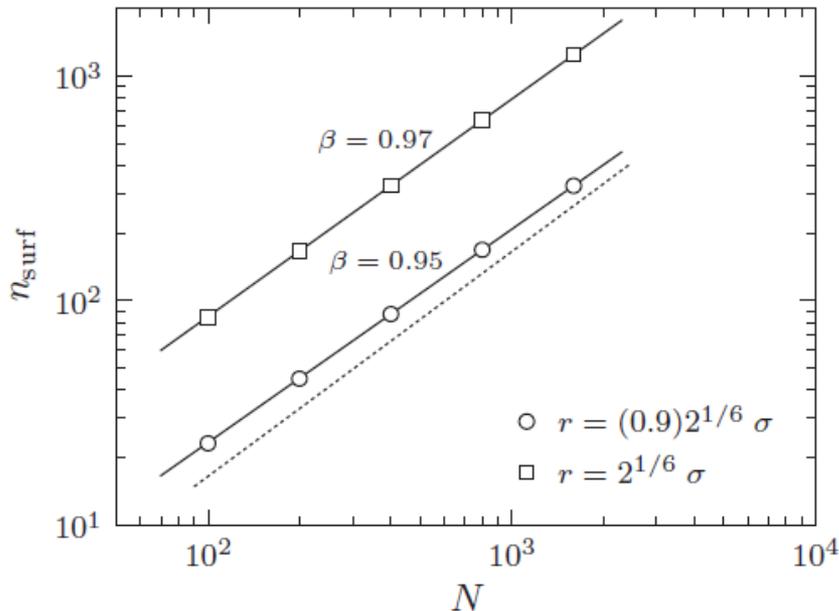
$$\approx \rho_{\text{intra}}(N_e)$$

Volume shared by several rings

sharing rings $\propto N^0$,
but incomplete segregation!

Melt of Ring Polymers

#“surface” beads $\propto N^\beta$



$\beta \approx 0.95 < 1$

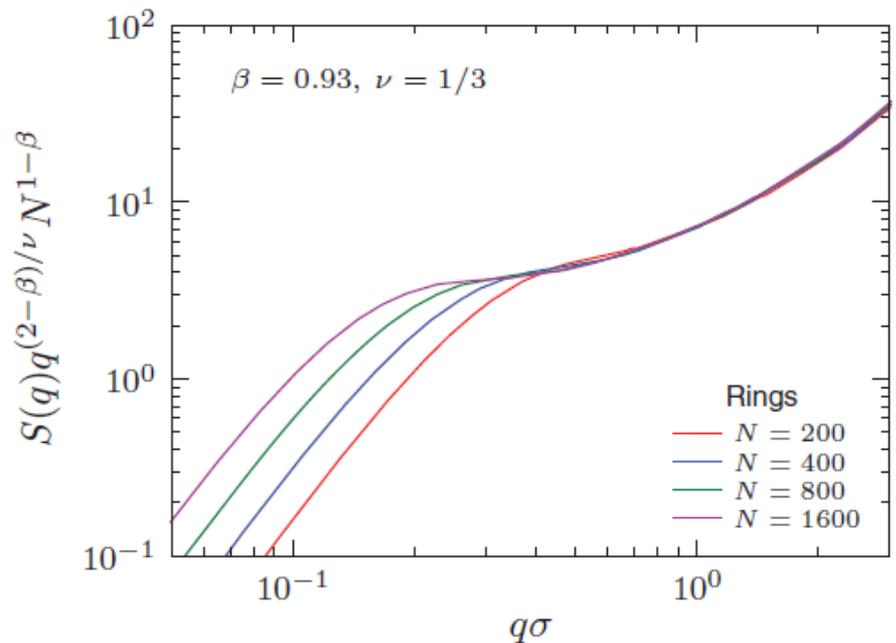
ring form factor $S(q)$

Without surface correction

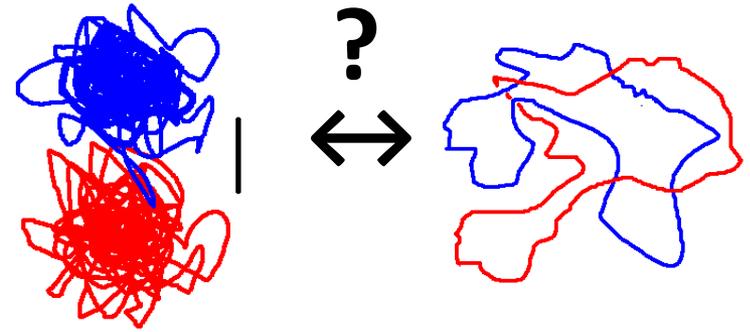
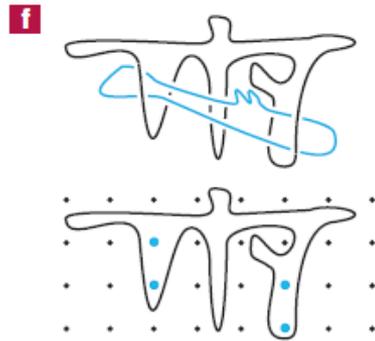
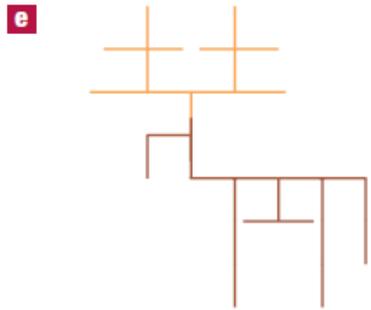
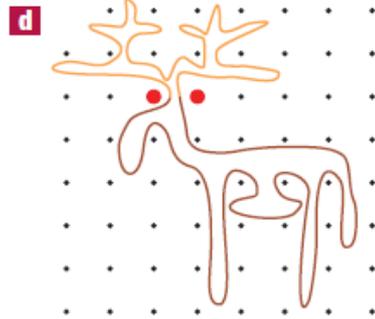
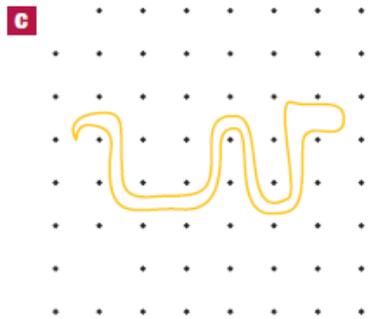
$$S(q) \propto q^{-1/\nu}$$

With surface correction

$$S(q) \propto q^{-(2-\beta)/\nu} N^{-(\beta-1)}$$



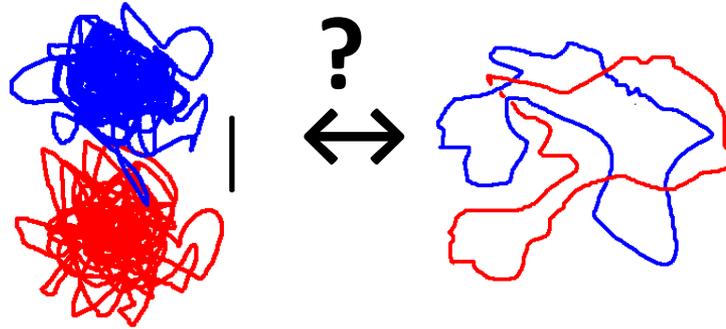
Melt of nonconcatenated rings: $N=100, 200, 400, 800, 1600 (\approx 58N_e)$



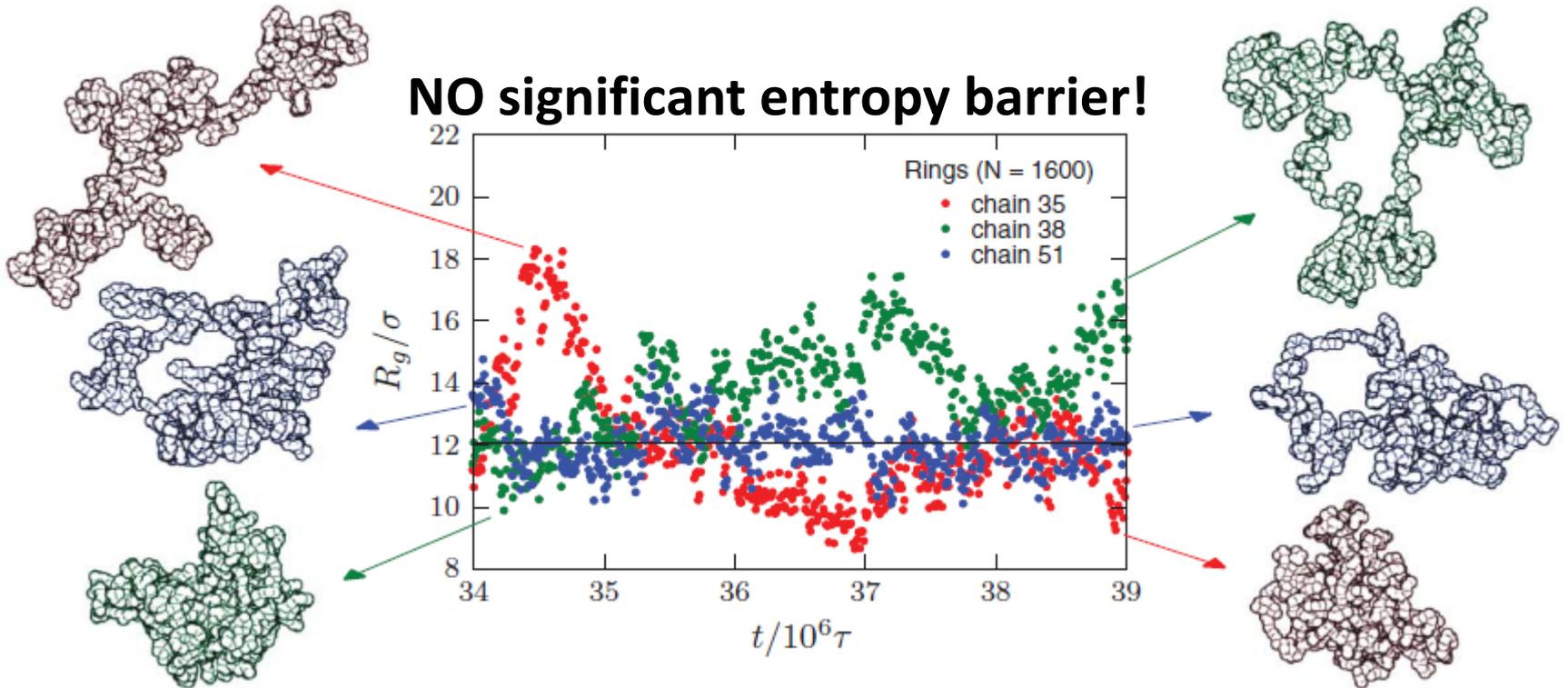
**significant
entropy barrier?**

**significant
entropy barrier!**

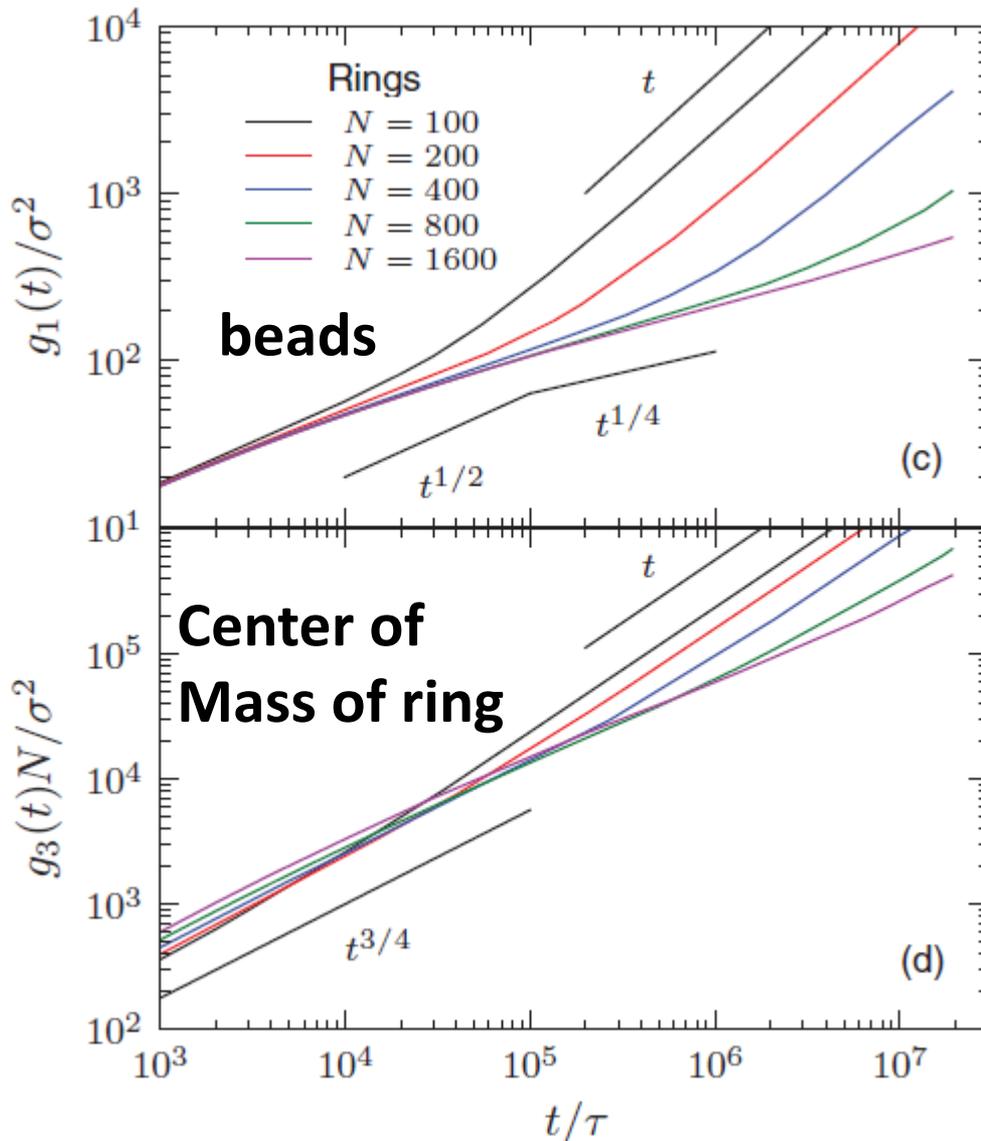
Melt of nonconcatenated rings: $N=100, 200, 400, 800, 1600 (\approx 58N_e)$



NO significant entropy barrier!



Melt of nonconcatenated rings: $N=100, 200, 400, 800, 1600$ ($\approx 58N_e$)



Mean square displacements

Slowing down

$$g_1 \propto t^{1/4}$$

and slower for large N

Subdiffusive

$$g_3 \propto t^x, x=0.75 \dots 0.65 (N=1600)$$

For distances
 larger than RG!

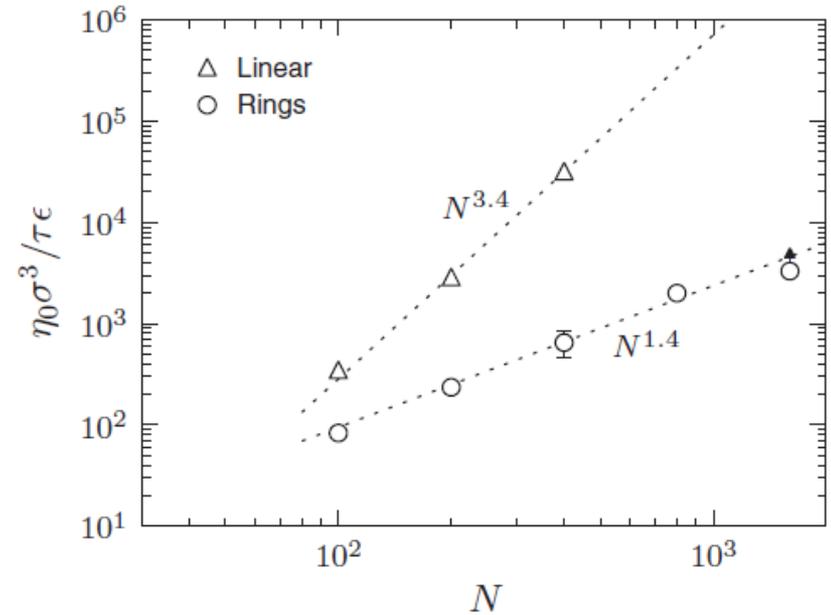
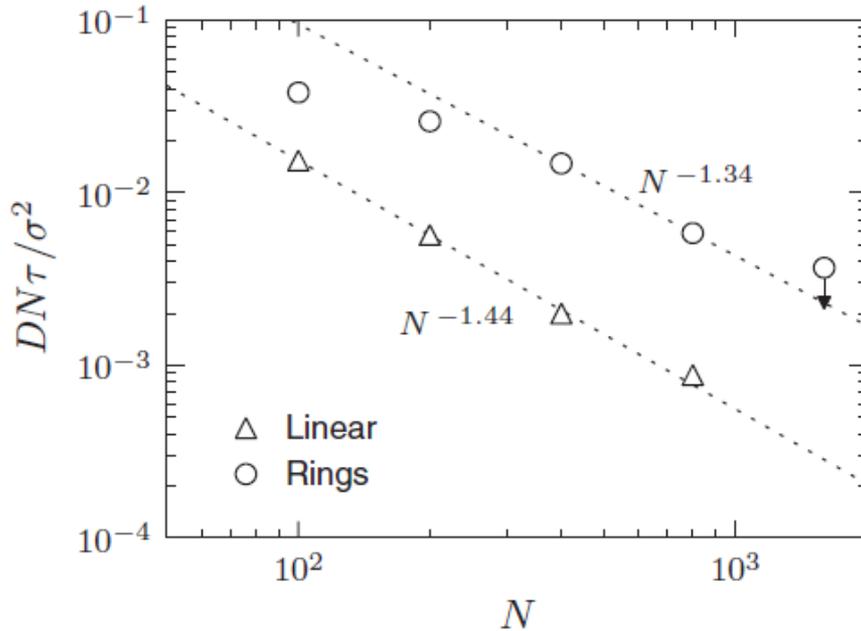
Melt of nonconcatenated rings: $N=100, 200, 400, 800, 1600 (\approx 58N_e)$



Diffusion

vs

viscosity



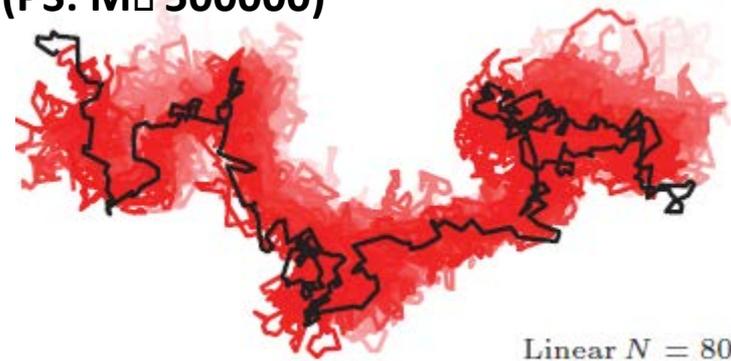
Mean square displacement of center of mass of ring and viscosity seem to decouple!

NO well established theoretical model
Problem of few linear polymers in rings and few rings in linears

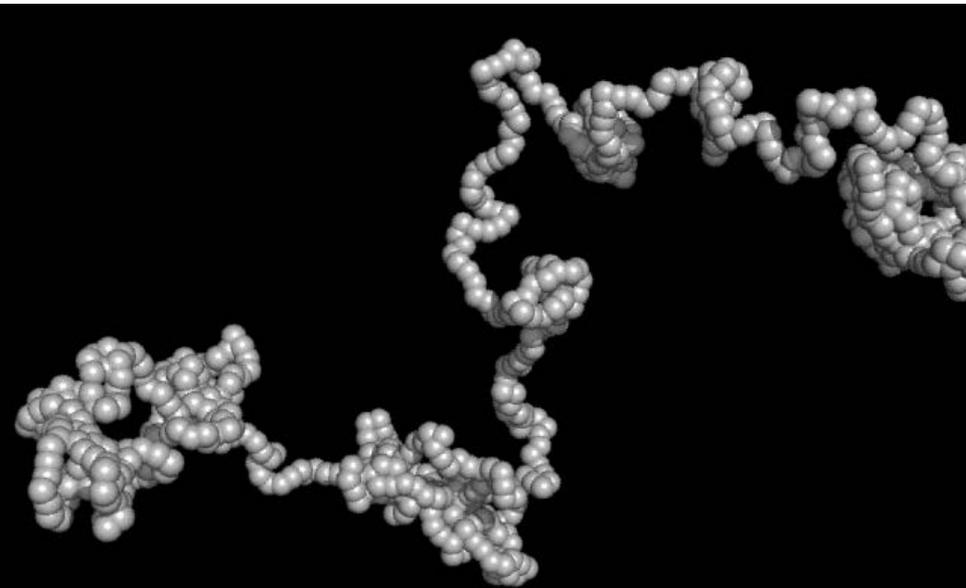
Halverson, Lee, Grest, Grosberg, KK
JCP, 2011



Linear, $N=800 \times 29 N_e$
(PS: $M \approx 500000$)

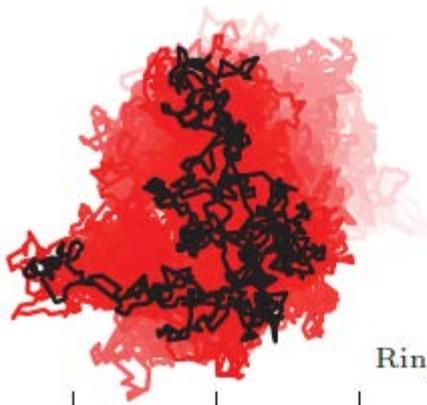


Linear $N = 800$



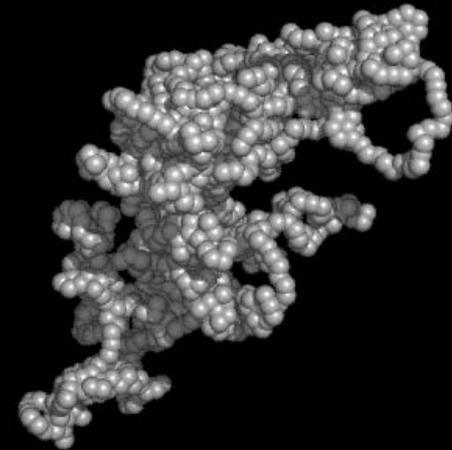
chain id = 78, monomers per chain = 800

time = 5.664e+06 tau $R_g(t)/\langle R_g \rangle = 1.22$ $\text{sqrt}[g_3(t)]/\langle R_g \rangle = 0.00$



Ring $N = 1600$

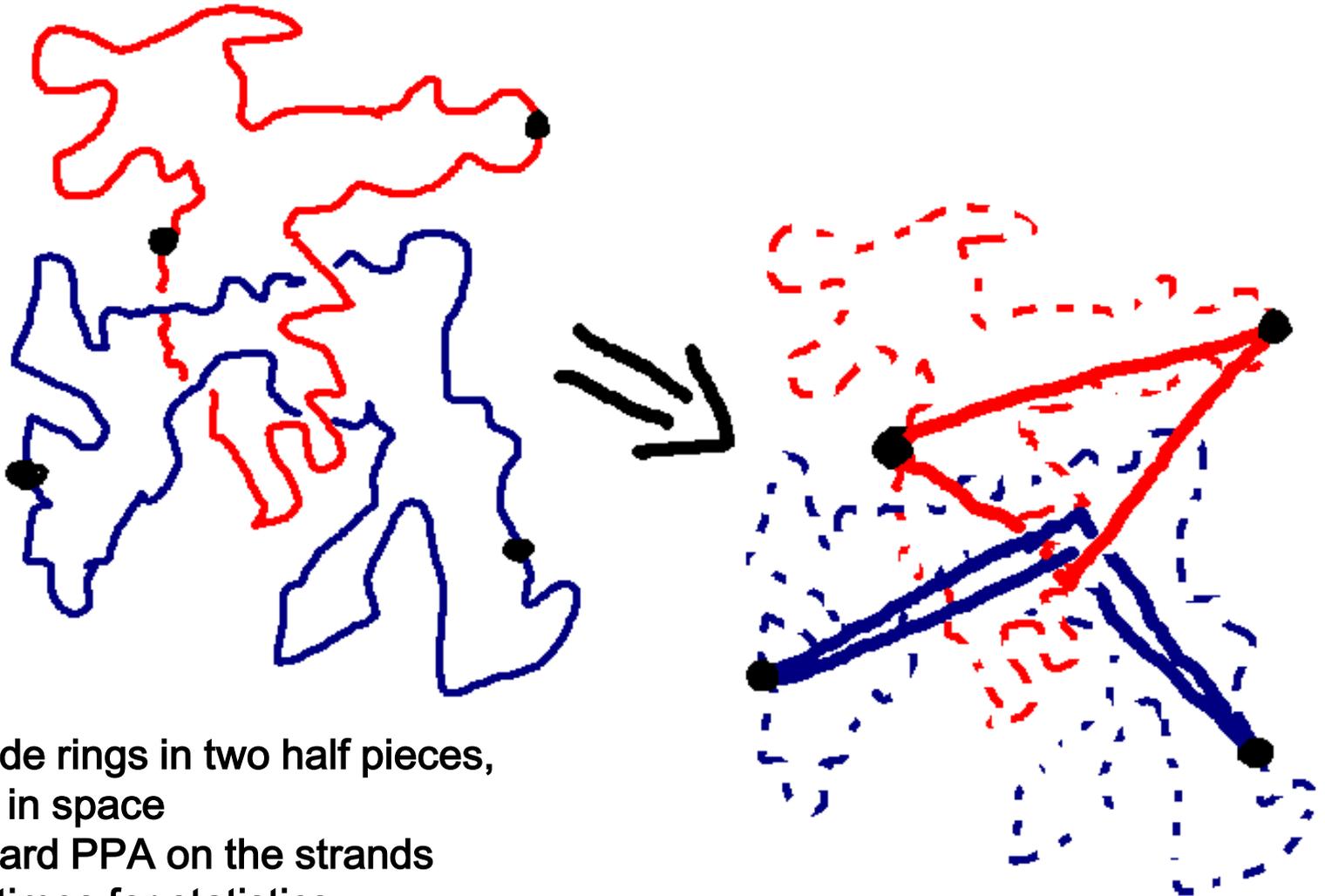
Ring, $N=1600$
(PS: $M \approx 1000000$)



chain id = 29, monomers per ring = 1600

time = 1.00e+04 tau $R_g(t)/\langle R_g \rangle = 1.13$ $\text{sqrt}[g_3(t)]/\langle R_g \rangle = 0.00$

Melt of Ring Polymers: Dynamics, Primitive Path Analysis



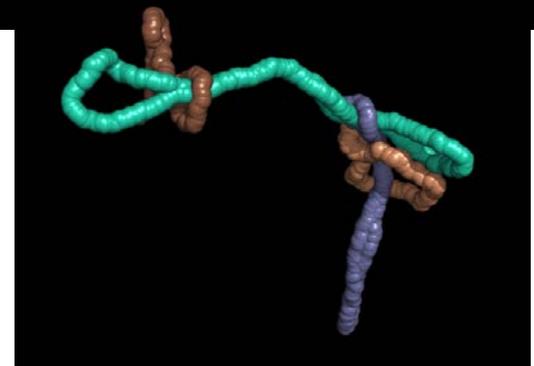
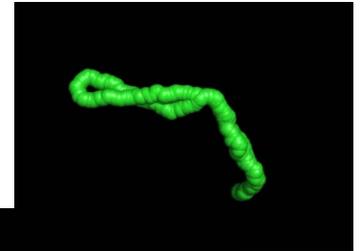
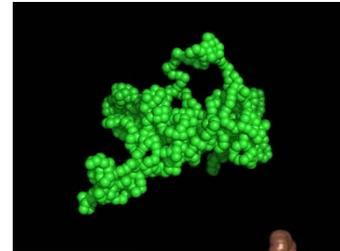
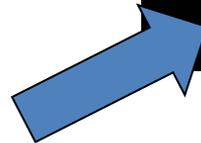
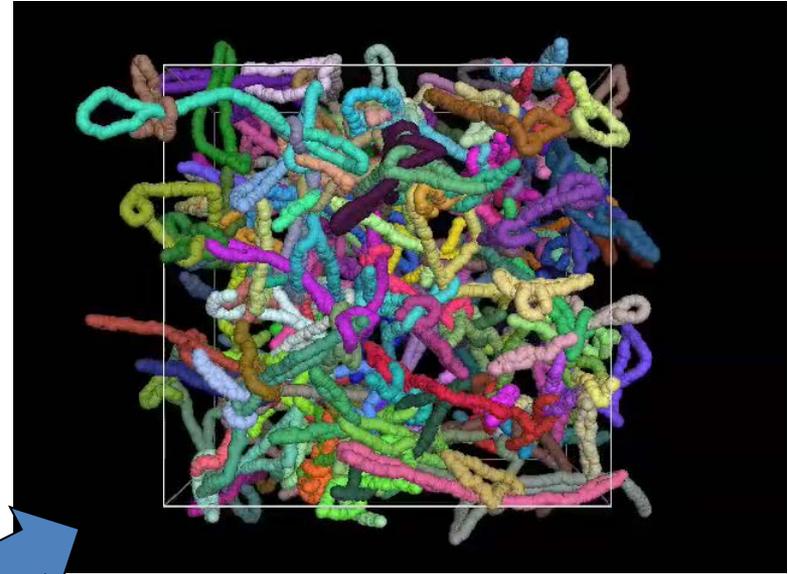
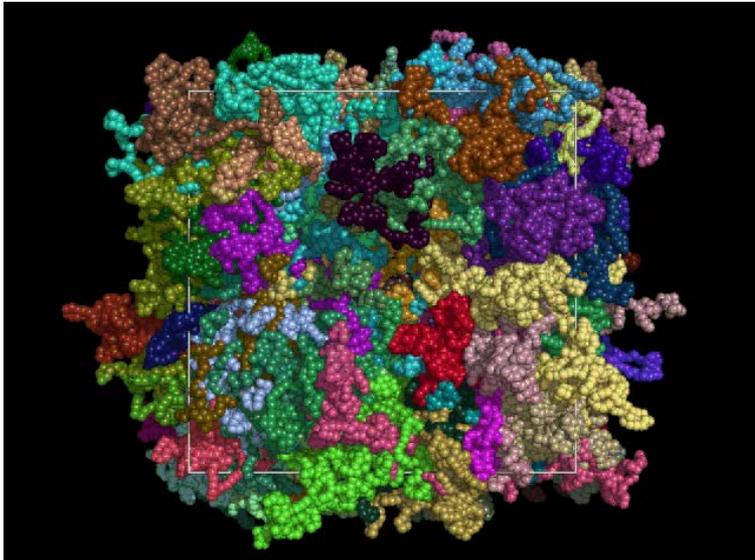
Randomly divide rings in two half pieces,
Fix end points in space
Perform standard PPA on the strands
Repeat many times for statistics



Melt of Ring Polymers: Dynamics

Primitive Path analysis
(similar to linear polymers)

Rings are somewhat
entangled



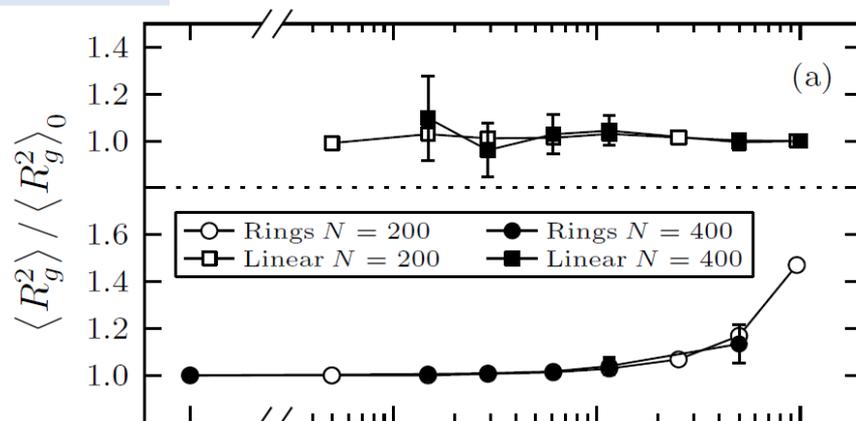
$$N_{e,\text{ring}} \approx 3 N_{e,\text{linear}}$$

N=800, # of rings 200

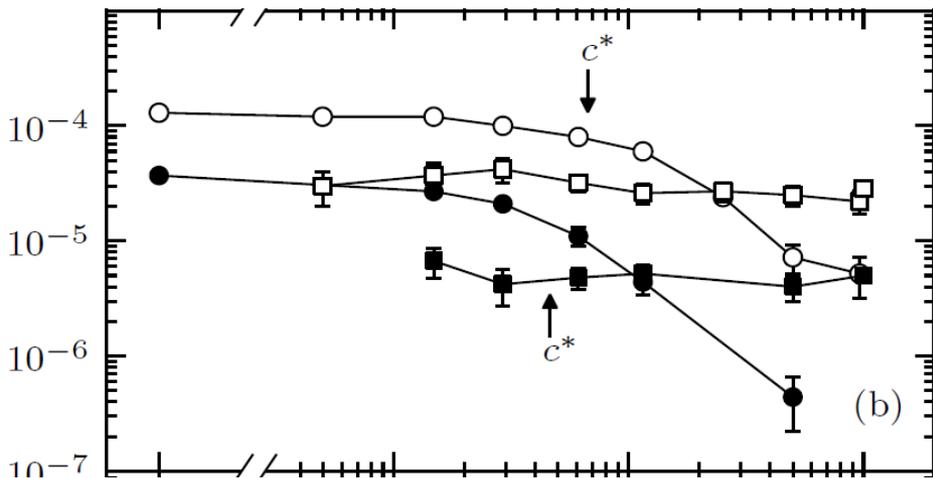
Another Puzzle: Rings plus linear Polymers

J. Halverson, et al, 2011

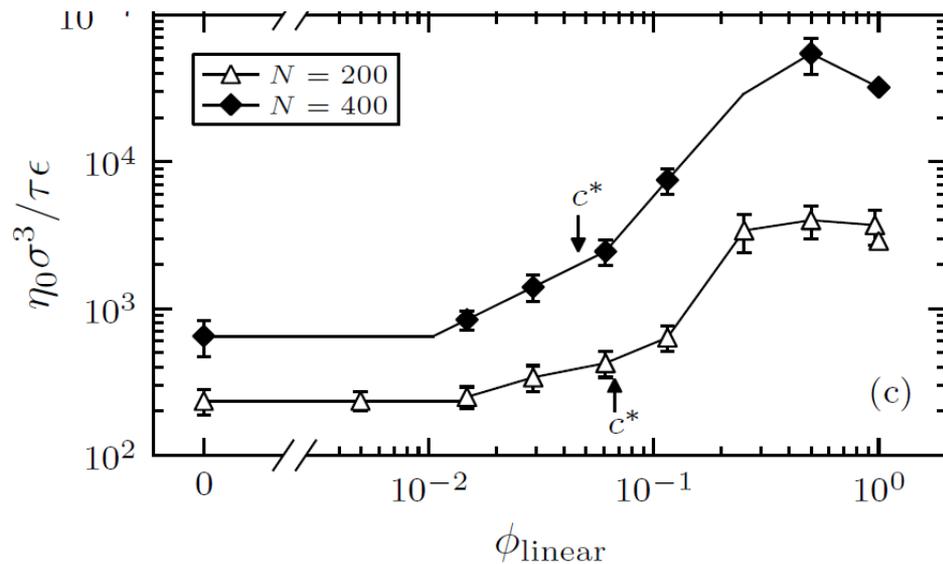
Radius of Gyration



Diffusion

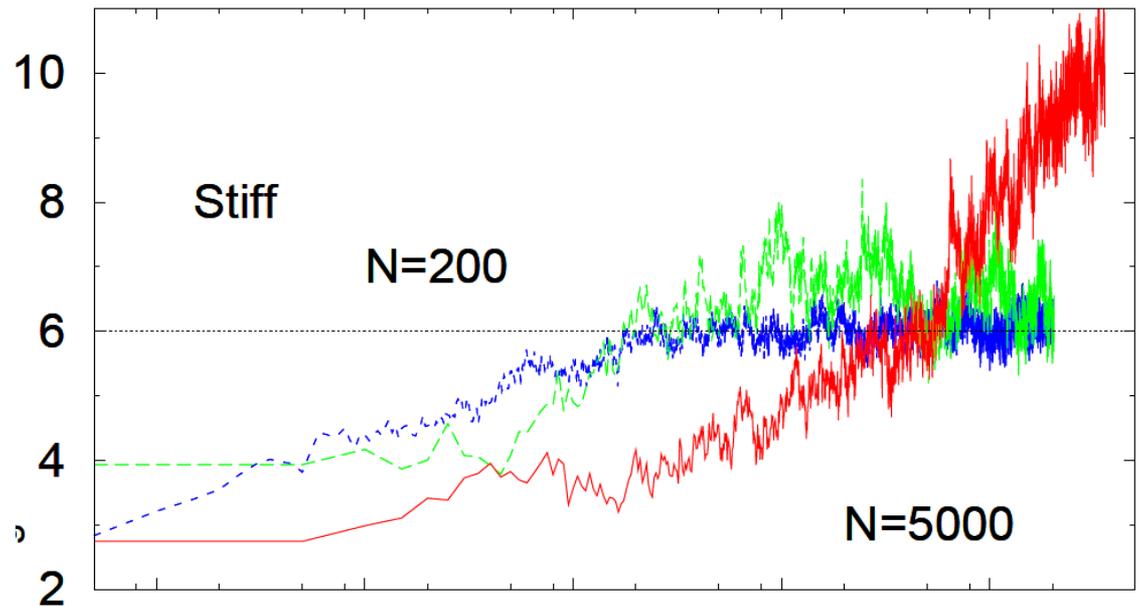
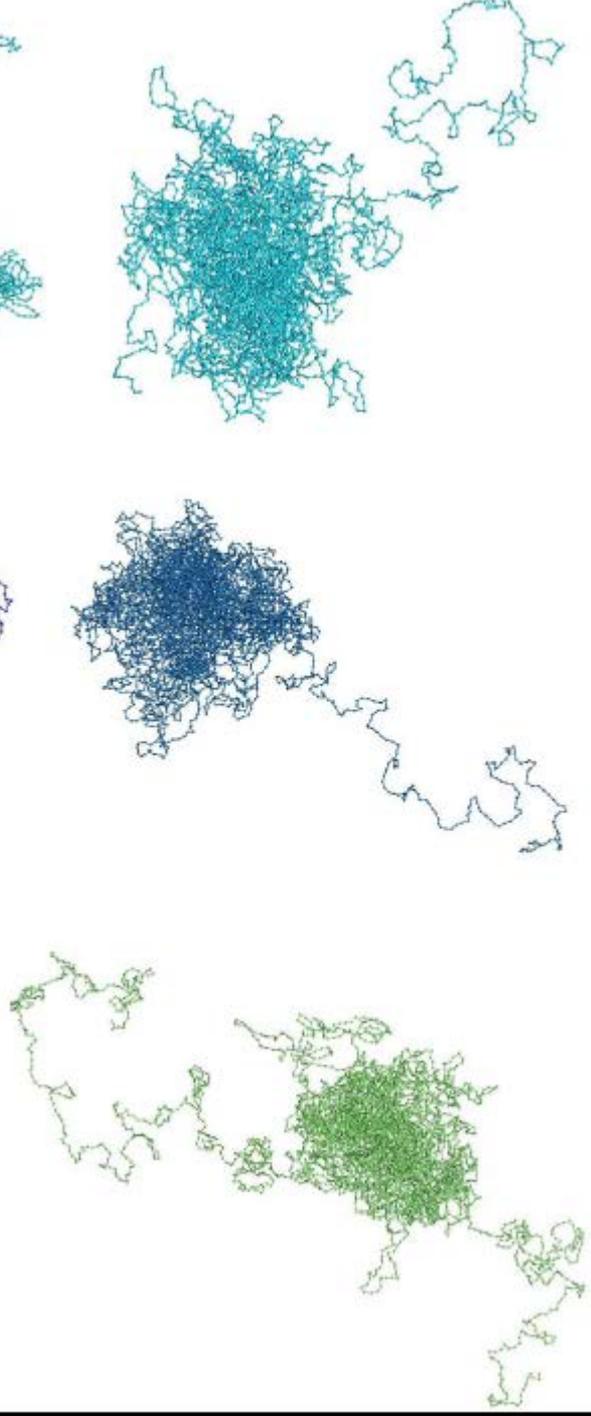


Viscosity





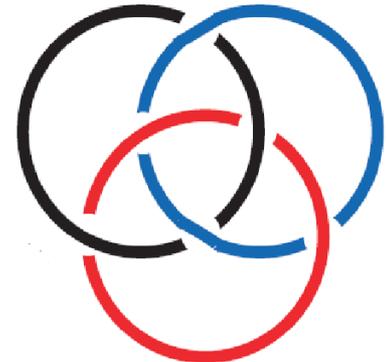
“Use” topological constraints: Melt of collapsed linear polymers



Topological Constraints Matter



- **Tube concept**
 - N_e derived from conformational statistics
 - Good agreement between simulations and experiment, both for melts and networks
- **Melt of non concatenated rings, melt of globules**
 - Still many open questions about conformations
 - Dynamics poorly understood so far
 - Mixtures, composites, applications beyond model system for chromosome territories...



Thank you for your attention



http://news.cnet.com/8301-31114_3-20030591-258.html