

Polyhedra as Ground States

LC Shells

2D systems described by field $\phi: \Sigma \rightarrow \left(\begin{smallmatrix} \text{Order} \\ \text{Field} \\ \text{Space} \end{smallmatrix} \right)$

Origin of shape $\equiv$ morphology

Shape change for $\Sigma \hookrightarrow \mathbb{R}^3$ is bend

$$H_{\text{bend}} = \frac{\kappa}{2} \int_{\Sigma} d^2x (\vec{H} - \vec{H}_0)^2 \quad \text{Will more functional}$$

$$\vec{H} = \frac{1}{2} \text{Tr} \vec{K}_{\alpha\beta} \quad \vee \quad \vec{K}_{\alpha\beta} = D_{\alpha} \vec{t}_{\beta}$$

2nd fundamental form

$$\vec{H} = \frac{1}{2} \text{Tr} \vec{K}_{\alpha\beta} = \frac{1}{2} g^{\alpha\gamma} \vec{K}_{\gamma\alpha} = \frac{1}{2} \frac{1}{\sqrt{\|g\|}} g_{\alpha\gamma} \vec{K}_{\gamma\alpha}$$

$$\vec{K} = K_{\alpha\beta} \hat{n}$$

$$\vec{H} = H \hat{n}$$

$$H_{\text{bend}} = \kappa \int d^2x (H - H_0)^2 \quad (1)$$

where $H = \frac{1}{2} K_{\alpha}^{\alpha}$

Further structure on surface:

orientational order $\equiv$ director field
 $\equiv$ line field

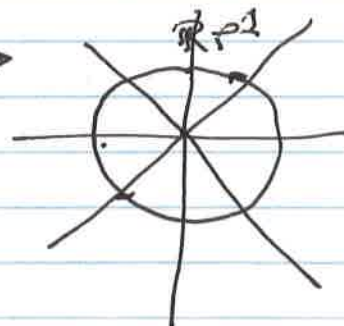
$$\vec{n}(x_1, x_2) : \underline{\Sigma} \rightarrow \underline{\mathbb{R}P^1}$$

$$\vec{n} \cdot \vec{n} = 1$$

$$\begin{aligned} \text{3D} / & \quad \vec{n} \in \mathbb{R}P^2 \\ & \equiv S^2 / \mathbb{Z}_2 \end{aligned}$$



2D $\rightarrow$



$$H_{\text{lc}} = \frac{1}{2} \int_{\Sigma} \sqrt{g} d^2x [K(\vec{D}\hat{n})^2] \quad (2)$$

$$\hat{n} = \vec{n}$$

$$H_{\text{total}} = H_{\text{lc}} + H_{\text{bend}}$$

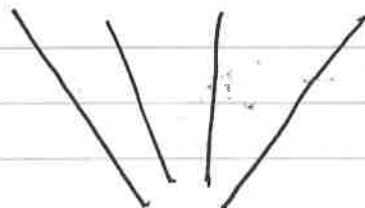
Origin of (2)

More generally invariants are

$$H_{\text{lc}} = \frac{1}{2} \int_{\Sigma} \sqrt{g} d^2x [K_1 \underbrace{(\vec{D} \cdot \hat{n})^2}_{\text{splay}} + K_3 \underbrace{(\vec{D} \times \hat{n})^2}_{\text{bend}}]$$

$$\vec{D} \cdot \hat{n} \neq 0$$

$$\vec{D} \times \hat{n} = 0$$

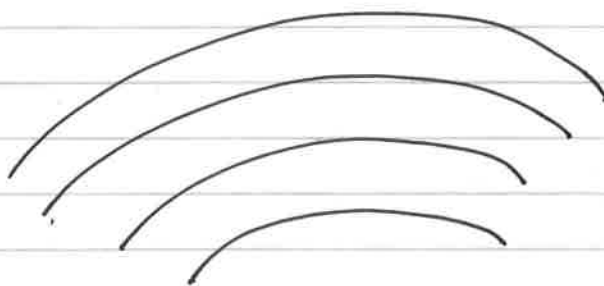


$$\vec{D} = \partial + \Gamma$$

N.B $\vec{D} \times \hat{n}$ is a scalar ϕ in 2D

$$\vec{D} \times \hat{n} \neq 0$$

$$\vec{D} \cdot \hat{n} = 0$$



$$\text{In 2D } (\hat{n} \times \vec{D} \times \hat{n})^2 = (\vec{D} \times \hat{n})^2$$

$$\text{Twist } \hat{n} \cdot (\vec{D} \times \hat{n}) = 0$$

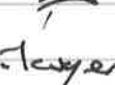
Case 1 $K_1 = K_3 = K \quad (\vec{D} \cdot \hat{n})^2 + (\vec{D} \times \hat{n})^2 = (\vec{D}\hat{n})^2$

Non-linear sigma model for projective field

Case 2 $K_3 \rightarrow \infty \Rightarrow \mathbb{D} \times \hat{n} = 0$ for those H_m

$\Rightarrow (\hat{n} \cdot \mathbb{D}) \hat{n} = 0$ ($\Rightarrow \hat{n}$ integral curve or equiv. geodesic)

Then $H_m = \frac{K_1}{2} \int (\mathbb{D} \cdot \hat{n})^2$ and call K_1 K
 Pure sp. lay. is a smectic  parallel layers

Now set $H_0 = 0$ (e.g. symmetric bilayer)  layer spacing

$$H = \frac{K}{2} \int_{\Sigma} (\mathbb{D} \hat{n})^2 + \frac{\kappa}{2} \int_{\Sigma} H^2$$

Limit 1 Rigid $\frac{\kappa}{K} \rightarrow \infty \quad \kappa \gg K$

$\Rightarrow H = H_{\text{willmore}}$ extremized by round sphere $\Sigma = S^2$ (round metric)
 $H = 1/R$ const.

$$H = \frac{\kappa}{2} \int_{S^2} d^2x \left(\frac{1}{R^2} \right)$$

$$= (2\pi\kappa) \stackrel{4\pi}{=} \text{const.}$$

$$H = \frac{1}{2} \left(\frac{1}{R} + \frac{1}{R} \right) = \frac{1}{R}$$

$\Rightarrow$ spherical metric or spherical smectic

Ribby defect structure (later)
 novel

Limit 2 Floppy w/ hard in space
 $K \gg \kappa \quad \left(\frac{\kappa}{K} \rightarrow 0 \right)$
 thin-shell $\kappa \sim (\text{thickness})^3$ ✓
 $K \sim (\text{thickness})$

Extreme case $\kappa = 0$

$$H = \frac{\kappa}{2} \int_{\Sigma} (\mathbb{D} \hat{n})^2$$

Σ top. homeomorphic to S^2

Extremize for $\{\hat{n}(x_1, x_2)\} = \{\hat{n}(\vec{x})\}$
and non-fold Σ (gap)

$\mathbb{D}^2 \hat{n} \rightarrow 0$ at extremum (harmonic projective fields)

$\hat{n} \neq \text{const.}$ costs energy - spectral inhomogeneity (texture)

$\mathbb{D} \neq 0$ costs energy - geometry not flat any Gaussian curvature costs energy

$\Rightarrow \Sigma$ should be developable ($K_{\text{Gauss}} = 0$)
but

Impossible since $\frac{1}{2\pi} \int_{\Sigma} K_{\text{Gauss}} = 2 = \chi_{S^2}$

$$K = \det K_{ij} = \frac{1}{R^2}$$

Solⁿ Σ should be facewise flat

i.e. Σ should be the boundary ∂ of a 3D polyhedron (facets of order)

i.e. $K_{\text{Gauss}}(\Sigma) = 0$ almost everywhere

$K_{\text{Gauss}}(\Sigma) = 0$ except at a discrete # of points $\vec{x}_i$
@ which $K_{\text{Gauss}}(\vec{x}_i) = c \delta(\vec{x} - \vec{x}_i)$ singular

Fortunately LC (2-fold) order on S^2 possesses zeros of $\hat{n}$ or singularities of $D\hat{n}$.
 What are the line fields w/ zeros $\equiv$ topological defects

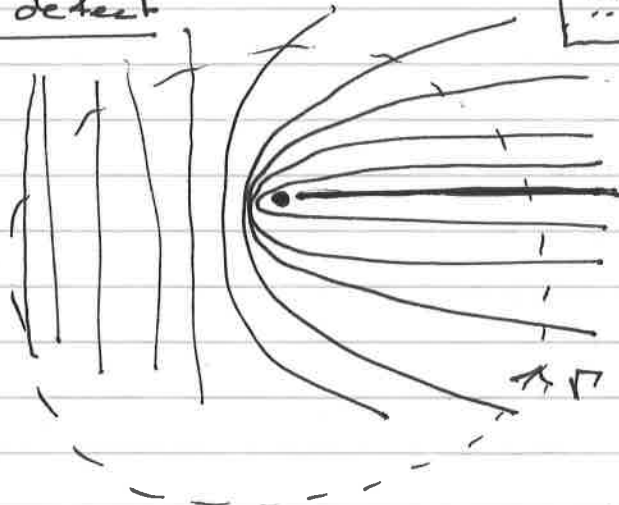
$$\pi_1(\mathbb{RP}^1) = \pi_1(S^1/\mathbb{Z}_2) = \mathbb{Z}$$

$$(3D) \quad \pi_1(\mathbb{RP}^2) = \pi_1(S^2/\mathbb{Z}_2) = \pi_0(\mathbb{Z}_2) = \mathbb{Z}_2$$

Simplest defect

NB. Point defects

$\equiv$ distinguished points



$\hat{n}$ winds by π as you encircle the defect
 $\Rightarrow \hat{n}$ traces closed loop on $\mathbb{RP}^1$

$$\Gamma: S^1 \rightarrow \mathbb{RP}^1 \quad \text{and } \Gamma \text{ is not contractible}$$

$$\Gamma \in \text{non-trivial equivalence class in } \pi_1(\mathbb{RP}^1)$$

$$S = \frac{1}{2\pi} \int_{\Gamma} \hat{n} \cdot d\mathbf{r} = \left(\frac{1}{2} \right)$$

disclination

$$\sum_i S_i = \chi_{S^2} = 2$$

$S_i =$ strength of defective line field

LC generalization of Poincaré-Hopf theorem
 $\equiv$ hairy ball theorem

$$\sum_i \text{ind}_i = \chi = 2$$

$i \in$ zero of a vector field $\mathbf{v}$ on S^2

⑥

Vector case $\vec{V}$ $\vec{V}: S^2 \rightarrow \mathbb{R}^3 \Rightarrow \vec{p} \cdot \vec{V}(\vec{p}) = 0$
vector-valued map $\forall \vec{p} \in S^2$

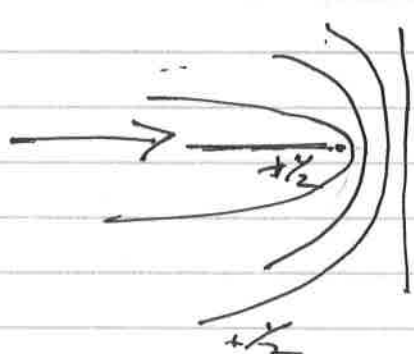
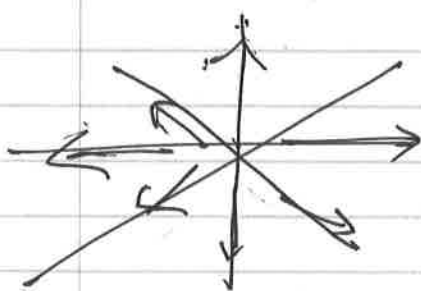
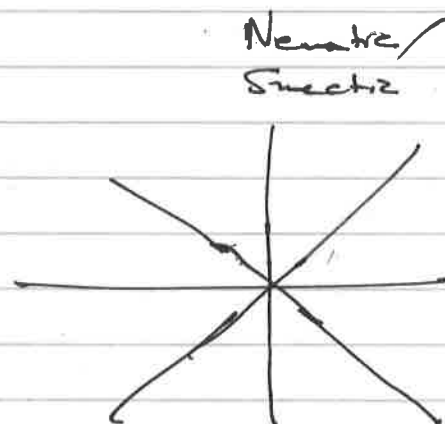
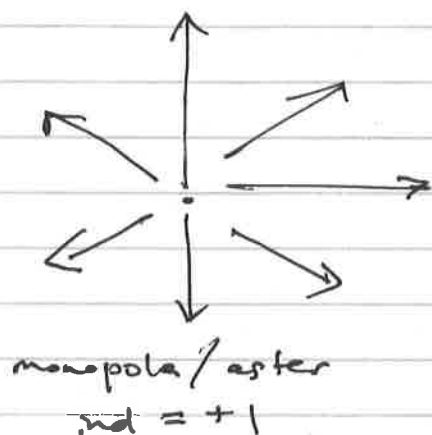
Thm Any continuous tangent vector field on S^2 must vanish somewhere.

Poincaré (1885)

Brouwer (1912)

2D / Poincaré-Hopf $\sum_{i \in \text{zeros}} \text{ind}_i = \chi$

Vector field $\pi_1(S^1) = \mathbb{Z}$



$$E_{(+1)} = (+1)^2 E_0$$

$$E_{\frac{1}{2}} = \frac{1}{4} E_0$$

$$E_{+1} = 2(E_{\frac{1}{2}} + E_{\frac{1}{2}})$$

$$= 2 E_{\frac{1}{2}} \text{ for } \frac{1}{2}$$

for the field

LC case — nematic / smectic

$$\sum_{i \in \{\text{defects}\}} s_i = 2 = \mathcal{K}(s^2)$$

$$s_i = \pm \frac{1}{2}, \pm 1, \pm \frac{3}{2}, \pm 2, \dots = n \left(\frac{1}{2} \right) \quad n \in \mathbb{Z}$$

Solns

1	strength +2	defect
2	+1	defects
1 (+1)	2 (+1/2)	
4 (1/2)		

Set of defects

Nematic

Energy

$$E_i = s_i^2 \frac{E_0}{K \ln(R/a)}$$

+2

$4E_0$

+1, +1

$2E_0$

+1, +1/2, +1/2

$\frac{3}{2} E_0$

$4 \left(+\frac{1}{2} \right)_s$

E_0

Smectic case

$$E_{+1} = 2E_{1/2} \quad (K_3 \rightarrow \infty)$$

*

$$E_{+2}, 2E_1 = 4E_{1/2}$$

+2

$4E_0$

+1, +1

$4E_0$

+1, $2 \left(+\frac{1}{2} \right)_s$

$4E_0$

$4 \left(\frac{1}{2} \right)_s$

$4E_0$

Degenerate

Now one needs at least four points to span a non-degenerate polyhedral boundary.
Minimal case $\Rightarrow$ tetrahedron

Vertices $\pm \frac{1}{2}$ defects

Thus $K \gg \kappa$, the ground state morphology of a LC vesicle w/ spherical topology is a faceted tetrahedral shell $\approx 2(\text{solid tetrahedron})$

N.B. Faceting has a completely different origin to the well-known buckling of elastic shells where buckling occurs above a critical size $R \sim \sqrt{\frac{\kappa}{Y}}$ $Y = \text{Young's modulus}$

Elastic YR^2 Bend κ

Buckling $\kappa < YR^2$
 $\Rightarrow R > R_c = \sqrt{\frac{\kappa}{Y}}$

Here it is scale free. The faceting is not size dependent. Depends only on $\frac{\kappa}{Y}$

Which tetrahedra support a suitable nematic defect configuration?

At each vertex i

$$\sum \theta_{int} = \pi$$

3 constraints

4 vertices but 3 independent constraints

$$\# \text{d.o.f} \quad 4 \text{ vertices} \times 3 \text{ coords} = 12$$

$$- (3 \text{ translations} + 3 \text{ rotations}) \quad \underline{6}$$

— scale invariance

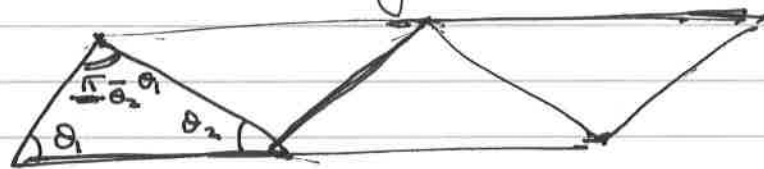
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3 constraints

2-dimensional moduli space of ground states

2

Moduli Space of grand states $\equiv$ scalene triangle (θ_1, θ_2)



Two parameter family of grand states for $\kappa = 0$

$$\kappa = 0^+$$

$H_{Le} \equiv 0$ for these structures

$$H_{bend} \propto \sum \text{edge lengths}$$

$H = H_{bend}$ minimised by regular tetrahedron (NB/ A fixed)

Grand state morphology is a regular tetrahedron when the bending rigidity $\rightarrow$ vanishingly small but non-zero

Other structures

Cylinder $\equiv$ Nanotube

$$H_{Le}^c = 0$$

$$H_{bend}^c \approx \frac{\kappa A}{8 R_c^2}$$

$$H_{tetra} = 4\kappa \frac{L}{b}$$

$$= \frac{4}{3^{3/4}} \kappa \frac{\sqrt{A}}{b}$$

$L = \text{edge length}$
 $b = \text{radius of curvature of ridge}$

$$A = \sqrt{3} L^2$$

$$L = \left(\frac{A}{\sqrt{3}}\right)^{1/2}$$

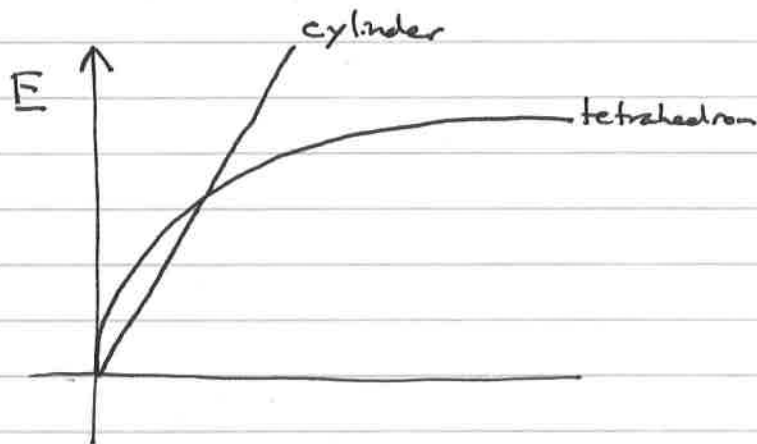


$$H = \frac{1}{2} \left(\frac{1}{R_c^2} \right)$$

$$H_{bend}^c = \frac{\kappa}{2} \int dx \left(\frac{1}{4 R_c^2} \right)$$

$$= \frac{\kappa}{8 R_c^2} \cdot (2\pi R_c \cdot L)$$

$$= \frac{\kappa A}{8 R_c^2}$$



Tetrahedron
i.e., ~~nanotube~~ wins for

$$\frac{4}{(3)^{1/4}} \frac{\kappa \sqrt{A}}{b} < \frac{\kappa}{8} \frac{A}{R_c^2}$$

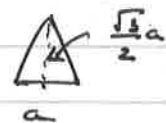
$$\frac{16}{\sqrt{3}} \frac{A}{b^2} < \frac{A^2}{64 R_c^4}$$

$$\Rightarrow A > \frac{2^{10}}{\sqrt{3}} \left(\frac{R_c^2}{b} \right)^2$$

$$\frac{A}{R_c^2} > \frac{1024}{\sqrt{3}} \left(\frac{R_c}{b} \right)^2$$

Cylinder wins
"Nanotube"

$$\frac{A}{R_c^2} < \frac{1024}{\sqrt{3}} \left(\frac{R_c}{b} \right)^2$$



$$A_{\text{face}} = \frac{\sqrt{3}}{4} a^2$$

$$A_{\text{tetra}} = \sqrt{3} a^2$$

Other topologies

Disk Planar contains $H_{\text{edge}} \propto \pi L$ $L = \text{length of perimeter}$
w/ edges

$$H = H_{\text{c}} + H_{\text{bend}} + H_{\text{edge}}$$

For large system sizes H_{edge} is costly and vesicles spontaneously close holes (surgery) to form closed S^2 rather than disk
(Helfrich 1974)

Nanofibres

Nano-size cylindrical micelles

Ellipsoidal polymersomes/vesicles: $\kappa/\kappa \sim 1$

Tori

Without topological constraint vanishing Gaussian curvature $\Rightarrow$ developable surfaces

Limit $K \gg K$

Geometry - round sphere

$K = K_1 = K_3$: Defects $(4(\pm\frac{1}{2})s)$ repel :

• Spherical Nematic

$$H_{\text{def}} = \frac{-\pi K}{8} \sum_{i \neq j} n_i n_j \ln(1 - \cos \beta_{ij}) + E(R) \sum_j n_j^2$$

$$\beta_{ij} = d_{ij}/R$$

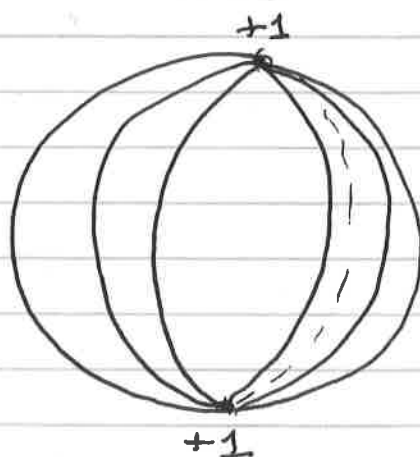
Extremum

4 defects at the vertices of an inscribed tetrahedron

$K_3 = 0$

$K = K_1$

Spherical Nematic



$$E_{+1} = 2E_{\frac{1}{2}}$$

Cut along the great circle corresponding to any given meridian and slide along director by an arbitrary angle α $\Rightarrow$ zero mode
 $\Rightarrow$ One parameter family of grounded states w/
 4 $\frac{1}{2}$ defects on a single great circle w/ arbitrary separation $\frac{1}{2}\pi$ 2 sets w/ angle α

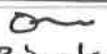


Functionalization of defects

$\pm$ defects/defective regions
 $\equiv$ valence

Geometrical distribution of defects $\equiv$ type of directional bonding

Types of Order

Order/ Symmetry Z_p	Field	Physical Model	Defects	Valence	Directional Bonding
1	vector	X-Y ferromagnet on S^2 Tilt Order	Dipolar Vortices Monopoles Index $+2, +1$	1 or 2	 Bipolar
2	RP^1	Nematic or Smectic LCs	strength $Z_k = \{ \frac{1}{2}, n \}$ strength $+2,$ $+1, +\frac{1}{2}$ disclination strength $\frac{1}{2}$	1, 2, 3 or 4	 Bipolar sp^2, sp^3 $K_3 \rightarrow \infty$ planar valence 4
3	trigonal	Triskelion proteins	$\frac{2\pi}{3}$ defects strength $\frac{1}{3}$	6	Octahedral
4	tetrahedral	Cubic colloidal made w/ cubic NPs	$\frac{\pi}{2}$ defects strength $\frac{1}{4}$	8	cube
6	crystalline/ hexatic	2d crystal	$\frac{\pi}{3}$ disclination strength $\frac{1}{6}$	12	icosahedral

Appendix

$$D \times \hat{n} = 0 \quad (\text{pure splay})$$

$$\Rightarrow \epsilon_{ab} D_a \hat{n}_b = 0$$

$$\Rightarrow D_a \hat{n}_b = D_b \hat{n}_a$$

$$\text{Thus } (\hat{n}_a D_a) \hat{n}_b = \hat{n}_a D_b \hat{n}_a = 0 \quad \text{since } \hat{n} \cdot \hat{n} = 1$$

$$\hat{n} \cdot D \hat{n} = 0$$

$$\text{Thus } (\hat{n} \cdot D) \hat{n} = 0$$

$\Rightarrow \hat{n}$ Integral curve or geodesic

$$\xrightarrow{\hat{n}} D \hat{n}$$