#	possible	score
1	25	
2	25	
3	25	
4	25	
5	25	
Total	100	

Preliminary Exam Partial Differential Equations 1:30 - 4:30 PM, Fri. Jan. 10, 2019 Room: Newton Lab (ECCR 257)

Student ID:_____

There are five problems. Solve four of the five problems. Each problem is worth 25 points.

A sheet of convenient formulae is provided.

1. Quasilinear first order equations.

Consider the Cauchy problem

$$u_t + (u + u^2)u_x = 0, \qquad x \in \mathbb{R}, \quad t > 0, u(x, 0) = f(x), \quad x \in \mathbb{R}.$$
 (1)

- (a) Suppose $f \in C^1(\mathbb{R})$ and f, f' are bounded functions. Prove that a continuously differentiable solution u(x,t) to Eq. (1) exists and is unique for $x \in \mathbb{R}$, $t \in [0, t_*)$ for some $t_* > 0$.
- (b) Provide an additional, necessary condition on f for the solution to Eq. (1) to exist for all t > 0, i.e., for u(x, t) to remain continuously differentiable for all t > 0.

Solution:

(a) Use the method of characteristics. Since

$$\frac{\mathrm{d}u}{\mathrm{d}t} = 0$$
 along the characteristic curves $x(t)$ satisfying $\frac{\mathrm{d}x}{\mathrm{d}t} = u + u^2$,

we can integrate each of these ODEs to obtain

$$u(x,t) = f(x_0)$$
, along $x = (u+u^2)t + x_0$.

This is the implicit solution for u

$$u(x,t) = f(x - (u + u^{2})t).$$
(2)

By the boundedness of f and f', there exist M, M' > 0 such that

$$|f(x)| < M, \quad |f'(x)| < M', \quad x \in \mathbb{R}.$$

In order to guarantee existence of a solution to this implicit equation, we define the mapping

$$F(u, x, t) = u - f(x - (u + u^2)t), \quad F : \mathbb{R} \times \mathbb{R}^2 \to \mathbb{R}$$

and apply the implicit function theorem. By virtue of the continuous differentiability of $f, F \in C^1(\mathbb{R} \times \mathbb{R}^2)$. Applying the initial data, $F(f(x_0), x_0, 0) = 0$. We compute

$$\frac{\partial F}{\partial u} = 1 + f'(x_0)(1+2u)t.$$

Then, for $t < T = \frac{1}{(1+2M)M'}$, $F_u \neq 0$ and the implicit function theorem implies that there exists a unique $u \in C^1(\mathbb{R} \times [0, t_*))$ such that F(u(x, t), x, t) = 0 for some $0 < t_* < T$ as desired.

(b) The Cauchy problem (1) suffers from the possibility of finite time singularity formation in which $|u_x| \to \infty$ as $(x,t) \to (x_b,t_b)$. To check when this occurs, we use the implicit form of the solution (2) to compute

$$u_x = f'(x_0)(1 - t(1 + 2u)u_x)$$

$$\Rightarrow u_x = \frac{f'(x_0)}{1 + t(1 + 2f(x_0))f'(x_0)}.$$

Consequently, a necessary condition to ensure that u_x is finite for all t > 0 is the requirement that

$$(1+2f(x_0))f'(x_0) \ge 0, \ x_0 \in \mathbb{R} \quad \iff \quad \frac{\mathrm{d}}{\mathrm{d}x}(f(x)+f(x)^2) \ge 0, \ x \in \mathbb{R}.$$

In other words, $f(x) + f(x)^2$ must be a monotone increasing function of x.

2. Heat Equation.

Let $D = (0, L) \times (0, T]$ and assume that $u \in C(\overline{D}) \cap C^2(D)$ is a solution to

$$u_t(x,t) = g(x)u_{xx}(x,t) + F(x,t), \qquad 0 < x < L, \quad 0 < t \le T.$$
(3)

$$u(x,0) = f(x), \qquad 0 < x < L, \qquad 0 < t \le T.$$
(3)

$$u(0,t) = r(t), \qquad 0 < t \le T, \qquad 0 < t \le T, \qquad 0 < t \le T, \qquad 0 < t \le T,$$

where g(x) > 0 for all $x \in (0, L)$.

(a) Let $B = \overline{D} \setminus D$. If $F \leq 0$, prove that

$$\max_{\bar{D}} u(x,t) = \max_{B} u(x,t).$$

(b) Prove that the solutions to Eq. (3) are unique.

Solution:

(a) First, let's prove the case F < 0. Suppose that the statement is not true. Then, the maximum of u would be attained at a point (x_0, t_0) with $0 < x_0 < L$,

 $0 < t \leq T$. By basic calculus, we would have $u_t(x_0, t_0) \geq 0$, $u_x(x_0, t_0) = 0$, $u_{xx}(x_0, t_0) \leq 0$. Evaluating (3) at this point, we would have

$$F(x_0, t_0) = u_t(x_0, t_0) - g(x_0)u_{xx}(x_0, t_0) \ge 0,$$

which contradicts F < 0. When $F \leq 0$, we can define $u^{\epsilon} = u - \epsilon t$. Then $u_t^{\epsilon} = g u_{xx}^{\epsilon} - \epsilon$ and we can apply the previous case with $F - \epsilon < 0$ to get $\max_{\bar{D}} u^{\epsilon}(x,t) = \max_B u^{\epsilon}(x,t)$. Letting $\epsilon \to 0$, we obtain the desired result.

(b) Suppose there are two solutions to (3), u_1 and u_2 . Then $v = u_1 - u_2$ would satisfy

$v_t(x,t) = g(x)v_{xx}(x,t),$	0 < x < L,	$0 < t \le T.$
v(x,0) = 0,	0 < x < L,	
v(0,t) = 0,	$0 < t \le T,$	
v(L,t) = 0,	$0 < t \le T.$	

For this system $v \equiv 0$ vanishes on B by applying part (a) to the IBVP for v and -v implying $u_1 \equiv u_2$ proving uniqueness.

Continue to next page

3. Wave Equation. Consider the initial boundary value problem (IBVP):

$$u_{tt} = c^2 u_{xx} \qquad x > 0, \ t > 0, u(x,0) = 0 \qquad x > 0, u_t(x,0) = \psi(x) \qquad x > 0, u_x(0,t) = 0 \qquad t > 0.$$

(a) Use an energy argument to prove the solutions to the above IBVP are unique, applying minimal assumptions on u(x,t). State these minimal assumptions. Solution: First define $w = u_1 - u_2$, which has corresponding IBVP

$$w_{tt} = c^2 w_{xx}, \ x, t > 0; \ w(x, 0) \equiv w_t(x, 0) \equiv w_x(0, t) \equiv 0,$$

and define the energy $E(t) = \frac{1}{2} \int_0^\infty w_t^2(x,t) + c^2 w_x^2(x,t) dx$, then

$$\begin{aligned} E'(t) &= \int_0^\infty w_t(x,t) w_{tt}(x,t) + c^2 w_{xt}(x,t) w_x(x,t) dx \\ &= \left[c^2 w_t(x,t) w_x(x,t) \right]_0^\infty + \int_0^\infty w_t(x,t) \left[w_{tt}(x,t) - c^2 w_{xx}(x,t) \right] dx \\ &= \int_0^\infty w_t(x,t) \cdot 0 dx = 0, \end{aligned}$$

assuming $u_t, u_x \to 0$ as $x \to \infty$ and using the Neumann BC. We also know $E(0) = \frac{1}{2} \int_0^\infty 0^2 + 0^2 dx = 0$, so $E(t) \equiv 0$. Assuming u_t and u_x are C^1 and L^2 integrable, this implies $w_t \equiv w_x \equiv 0$ implying $w \equiv 0$ (using $w(x, 0) \equiv 0$), so $u_1 \equiv u_2$.

(b) Solve for u(x,t). What assumptions on ψ are needed for a classical solution? **Solution:** The PDE implies $(\partial_t - c\partial_x)(\partial_t + c\partial_x)u = 0$ so u = F(x - ct) + G(x + ct), so when x > ct, u is not influenced by the boundary and d'Alembert's solution implies

$$u(x,t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(y) dy,$$

but when x < ct, the boundary condition influences the solution. To see how, note

$$F(x) + G(x) = 0$$

$$-cF'(x) + cG'(x) = \psi(x) \quad \Rightarrow \quad -F(x) + G(x) + A = \frac{1}{c} \int_0^x \psi(y) dy$$

$$F'(-ct) + G'(ct) = 0 \quad \Rightarrow \quad -F(-x) + G(x) + B = 0.$$

Subtracting the last two integrated equations at x = 0 implies A = B, whereas summing the first two equations implies

$$G(x) = \frac{1}{2} \left[\frac{1}{c} \int_0^x \psi(y) dy - A \right],$$

so that when x < ct, this along with the third equation imply

$$\begin{split} u(x,t) &= F(x-ct) + G(x+ct) = G(ct-x) + A + G(x+ct) \\ &= \frac{1}{2c} \left[\int_0^{ct-x} \psi(y) dy + \int_0^{x+ct} \psi(y) dy \right]. \end{split}$$

We need that $\psi \in C^1$; and $\lim_{x\to 0^+} \psi(x) = 0$.

4. Poisson's Equation/Green's Functions.

Consider the problem

$$\Delta u(\mathbf{x}) = f(\mathbf{x}), \qquad \mathbf{x} \in \Omega,$$

$$u(\mathbf{x}) = g(\mathbf{x}), \qquad \mathbf{x} \in \partial\Omega,$$
(4)

where $\Omega = \{ \mathbf{x} = (x, y, z) \subseteq \mathbb{R}^3 : z > 0, \| \mathbf{x} \|_2 < R \}.$

(a) Construct an appropriate Green's function for this problem.

Solution: Leveraging the fundamental solution $\Phi(\mathbf{x} - \mathbf{y}) = 1/(4\pi |\mathbf{x} - \mathbf{y}|)$ to $-\Delta u = \delta(\mathbf{x} - \mathbf{y})$ and defining $\hat{\mathbf{x}} = (x, y, -z)$, $\bar{\mathbf{x}} = R^2 \mathbf{x}/|\mathbf{x}|^2$, and $\tilde{\mathbf{x}} = R^2 \hat{x}/|\mathbf{x}|^2$, we write

$$G(\mathbf{x}, \mathbf{y}) = \Phi(\mathbf{x} - \mathbf{y}) - \Phi(\hat{\mathbf{x}} - \mathbf{y}) - \Phi(\frac{|\mathbf{x}|}{R} |\bar{\mathbf{x}} - \mathbf{y}|) + \Phi(\frac{|\mathbf{x}|}{R} |\tilde{\mathbf{x}} - \mathbf{y}|),$$

which vanishes along $\partial \Omega$ as along the sphere surface $\mathbf{x} = \bar{\mathbf{x}}$ and $\hat{\mathbf{x}} = \tilde{\mathbf{x}}$ and along z = 0 we have $\mathbf{x} = \hat{\mathbf{x}}$ and $\bar{\mathbf{x}} = \tilde{\mathbf{x}}$.

(b) Using the Green's function found in (a), construct an explicit formula for the solution in terms of the functions f and g.

Solution: Using Green's theorem, we find

$$\begin{split} \int_{\Omega} \left(u(\mathbf{y}) \Delta G(\mathbf{x}, \mathbf{y}) - G(\mathbf{x}, \mathbf{y}) \Delta u(\mathbf{y}) \right) d\mathbf{y} &= \int_{\partial \Omega} \left(u(\mathbf{y}) \frac{\partial G}{\partial n}(\mathbf{x}, \mathbf{y}) - G(\mathbf{x}, \mathbf{y}) \frac{\partial u}{\partial n}(\mathbf{y}) \right) dS_{\mathbf{y}} \\ - u(\mathbf{x}) - \int_{\Omega} G(\mathbf{x}, \mathbf{y}) f(\mathbf{y}) d\mathbf{y} &= \int_{\partial \Omega} \frac{\partial G}{\partial n}(\mathbf{x}, \mathbf{y}) g(\mathbf{y}) d\mathbf{y} \\ u(\mathbf{x}) &= -\int_{\Omega} G(\mathbf{x}, \mathbf{y}) f(\mathbf{y}) d\mathbf{y} - \int_{\partial \Omega} \frac{\partial G}{\partial n}(\mathbf{x}, \mathbf{y}) g(\mathbf{y}) d\mathbf{y} \end{split}$$

5. Separation of Variables.

Consider the following IBVP for the heat equation

$$u_t = k u_{xx}, \qquad x \in (0, 1), \quad t > 0, \quad k > 0,$$

$$u(x, 0) = f(x), \qquad x \in (0, 1),$$

$$u_x(0, t) = u(0, t), \qquad t > 0,$$

$$u_x(1, t) = -u(1, t), \qquad t > 0.$$

(a) Assuming separated solutions u(x,t) = X(x)T(t), derive the boundary value problem for X(x), and show it is a symmetric Sturm-Liouville problem.

Solution: Plugging the ansatz into the PDE, we have XT' = kX''T implying

$$\frac{T'}{kT} = \frac{X''}{X} = -\lambda \quad \Rightarrow \quad T' = -\lambda kT; \quad X'' = -\lambda X$$

with boundary conditions X'(0) = X(0) and X'(1) = -X(1), so by assuming $X, Y \in L^2(0,1)$ with Y'(0) = Y(0) and Y'(1) = -Y(1):

$$\begin{split} \langle X'', Y \rangle &= \int_0^1 X''(x) \bar{Y}(x) dx = X'(1) \bar{Y}(1) - X'(0) \bar{Y}(0) - \int_0^1 X'(x) \bar{Y}'(x) dx \\ &= X'(1) \bar{Y}(1) - X'(0) \bar{Y}(0) - X(1) \bar{Y}'(1) + X(0) \bar{Y}'(0) + \int_0^1 X(x) \bar{Y}''(x) dx \\ &= -X(1) \bar{Y}(1) - X(0) \bar{Y}(0) + X(1) \bar{Y}(1) + X(0) \bar{Y}(0) + \int_0^1 X(x) \bar{Y}''(x) dx \\ &= \int_0^1 X(x) \bar{Y}''(x) dx = \langle X, Y'' \rangle \end{split}$$

(b) Solve for the general form of the solution u(x,t). Make sure you show all eigensolutions must decay in time, and do not blow up as $t \to 0$.

Solution: Starting with the BVP for X(x), we note when $\lambda = -\nu^2 < 0$, solutions are of the form $X(x) = A_{\nu}e^{\nu x} + B_{\nu}e^{-\nu x}$, so Robin BCs imply

$$X(0) = A_{\nu} + B_{\nu} = \nu(A_{\nu} - B_{\nu}) = X'(0) \quad \rightarrow \quad B_{\nu} = \frac{\nu - 1}{\nu + 1}A_{\nu}$$

and

$$X(1) = A_{\nu} \left[e^{\nu} + \frac{\nu - 1}{\nu + 1} e^{-\nu} \right] = -A_{\nu} \nu \left[e^{\nu} - \frac{\nu - 1}{\nu + 1} e^{-\nu} \right] = -X'(1)$$

$$\rightarrow \quad (\nu + 1) e^{\nu} = \frac{(\nu - 1)^2}{\nu + 1} e^{-\nu} \quad \rightarrow \quad e^{2\nu} = \frac{(\nu - 1)^2}{(\nu + 1)^2} \quad \rightarrow \quad \nu = 0,$$

and for $\lambda = 0$ and X(x) = A + Bx, we have X(0) = A = B = X'(0) and X(1) = 2A = -A = -X'(1) implies A = 0, so the only nontrivial solutions are of form $X(x) = A_{\nu} \sin(\nu x) + B_{\nu} \cos(\nu x)$ with $\lambda = \nu^2 > 0$, so the Robin BCs imply

$$X(0) = B_{\nu} = \nu A_{\nu} = X'(0)$$

and

$$X(1) = A_{\nu}(\sin\nu + \nu\cos\nu) = -\nu A_{\nu}(\cos\nu - \nu\sin\nu) = -X'(1) \quad \to \quad \tan\nu = \frac{2\nu}{\nu^2 - 1}$$

which has an infinite set of solutions $\nu_{1,2,3,\ldots}$, since the RHS is decreasing and positive for $\nu > 3$ while $\tan \nu$ is periodic and ranges from $(-\infty, \infty)$. Moreover, this is a symmetric SL problem, so it is guaranteed to have a complete orthonormal set of eigenfunctions.

Clearly $\nu = 0$ is a solution, but this yields the trivial $X(x) \equiv 0$. Defining $X_j(x) = A_j(\sin(\nu_j x) + \nu_j \cos(j\pi x)) = A_j\phi_j(x)$ and $\lambda_j = \nu_j^2$, so $T_j(t) = e^{-\nu_j^2 k t}$, and

$$u(x,t) = \sum_{j=1}^{\infty} A_j e^{-\nu_j^2 k t} \phi_j(x), \qquad A_j = \frac{\langle \phi_j(x), f(x) \rangle}{\langle \phi_j(x), \phi_j(x) \rangle}.$$

(c) State minimal assumptions on f(x) needed so $u(x,t) \in C^2$ on $x \in (0,1)$ and t > 0. Solution: We need only that f(x) be L^2 -integrable, so A_j can be defined and the series converges for all $x \in (0,1)$ due to the exponential term.